

# Master's thesis

Kristian Tveitstøl

## Using EELS to measure the local conductivity in Aluminium

Master's thesis in Nanotechnology

Supervisor: Randi Holmestad

Co-supervisor: Ragnvald Mathiesen and Emil Frang Christiansen

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# Abstract

The plasmonic properties of pure aluminium have been investigated using electron energy loss spectroscopy in order to reproduce its macroscopic electrical conductivity. Temperatures varying from room temperature to 500°C have been investigated with an energy resolution of  $\lesssim 100$  meV. The plasmon energy was found to be  $\approx 15$  eV at room temperature, decreasing with approximately  $0.5$  meV K $^{-1}$  and had an asymmetric linewidth of  $\approx 500$  meV and 700 meV resulting in an underestimate of the conductivity with one order of magnitude. The findings suggest that the full width of half maximum broadens with increasing temperature, though no conclusive relation has been outlined.

# Sammendrag

De plasmoniske egenskapene til ren aluminium har blitt undersøkt ved hjelp av elektronenergitapsspektroskopi med hensikt til å reproduusere den makroskopiske elektriske ledningsevnen til materialet. Dette har blitt gjennomført for temperaturer fra romtemperatur til 500 °C med en energioppløsning på  $\lesssim 100$  meV. Plasmontoppen ble målt til å være sentrert ved omtrent 15 eV ved romtemperatur, sank med  $0.5 \text{ meV K}^{-1}$ , og hadde en asymmetrisk halvverdibredde på 500 og 700 meV, som medførte at metoden underestimerte ledningsevnen med en størrelsesorden. Funnene antyder at det er en temperatureavhengighet for halvverdibredden, men ingen konkluderene relasjon har blitt foreslått. Drude-modellen og den dielektriske formalismen innen mange-partikkel systemer har blitt brukt som utgangspunkt for det teoretiske grunnlaget for at elektronenergitapsspektroskopi kan bli brukt for å finne ledningsevnen, dog har det blitt antatt at de eksperimentelle resultatene kan forklares ut ifra plasmonspreddning i langbølgegrensen og kan direkte knyttes opp mot Drude-modellen. Det har videre blitt spekulert i at dette ikke har vært tilfellet for forsøkene som har blitt gjennomført, og dermed at bidrag fra plasmonspreddning som ikke er i langbølgegrensen i tillegg til interbåndsoverganger har hatt en utbredende effekt. For å begrense den påfølgende asymmetrien fra bevegelsesmengdeoverganger med kortere bølgelengde i eventuelle fremtidige forsøk har det blitt foreslått at akselerasjonsspenningen i transmissjonselektron-mikroskopskolonnen og oppsamlinksinkelen inn til spektrometeret reduseres.

# Preface

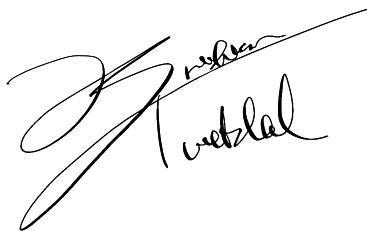
This master thesis completes a five year master program in Nanotechnology at the Norwegian University of Science and Technology (NTNU), and is a continuation of a preliminary project thesis during the autumn 2022. The low-loss regime of aluminium and platinum was investigated during that project. The work concluded that the data quality was inadequate for the purpose due to lack of energy resolution and carbon contamination, respectively.

I would like to thank my supervisors Randi Holmestad, Ragnvald Mathiesen and Emil Frang Christiansen for their guidance and patience during both this work and the project thesis. Not many people can say that they have had the pleasure to visit both England and Japan during their master thesis. Thank you Randi for fixing this and all the additional thoughts and experts you have gained valuable considerations from. Emil, thank you for all your support with sample preparation, TEM training, EELS acquisition, and tips and tricks for especially plotting (and Hyperspy last semester). I still remember Friday 16.12.2022 when we spent 12 h to acquire the EELS data for my preliminary project thesis and apologized for “stealing” my Friday afternoon. *I* am sorry for stealing *your* Friday afternoon. Thank you Ragnvald for all your insights in the physical aspects of the plasmon peak. I don’t know how this would have been without your knowledge and ability and patience to explain the concepts to enhance my understanding. Additionally, I would like to thank Per Erik Vullum for his contribution on the FIB, making sure that everything was in order prior to EELS analysis.

I would also like to thank Simon Fairclough and John Walmsley from the University of Cambridge for conducting the experiments leading to all the EELS data presented in this work. Unless otherwise specified, all data presented in this thesis has been acquired by Simon Fairclough.

When referring to a *plasmon* it will throughout the entire text be implicit that we are talking about the *bulk*-plasmon and not the *surface*-plasmon as the latter has not been investigated.

Til slutt vil jeg rette en takk til alle som har bidratt til 5 uforglemmelig år i tigerstaden. Takk for dere har gjort studenthverdagen til å ikke bare omhandle studier. En spesiell takk rettes til medlemmene av Timini kull-18 (wwwiiiiiiiiiiiiii) og NTNUI Samba (Samba-toget ruller videre!). Jeg er glad i dere < 3

A handwritten signature in black ink, appearing to read "Kristian Tveitstøl". The signature is fluid and cursive, with some parts written in a slightly different style.

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Kristian Tveitstøl  
June 23, 2023

# List of Aberrations

**ADF** angular dark field. 20, 21

**EELS** electron energy loss spectroscopy. 2, 9

**FIB** focused ion beam. 15, 20

**FWHM** full width at half maximum. 8, 9, 18

**GIF** Gatan imaging filter. 15

**SNR** signal-to-noise-ratio. 12, 17, 20

**STEM** scanning transmission electron microscope/microscopy. 2

**TEM** transmission electron microscope/microscopy. 4, 5, 15

**ZLP** zero-loss peak. 12

# Selected Experimental Values for Aluminium

Table 1: selected physical properties for aluminium. The conductivity is taken at room temperature. The electron density is taken from Ref. [1], while the other properties are taken from Ref. [2].

Quantity	Symbol	Value
Conductivity	$\sigma$	$37.7 \cdot 10^6 \text{ S m}^{-1}$
Temperature coefficient	$\alpha$	0.0053
Coefficient of linear expansion	$\beta$	$23 \cdot 10^{-6} \text{ K}^{-1}$
Electron density	$n$	$18.1 \cdot 10^{22} \text{ cm}^{-3}$

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# Chapter 1

## Introduction

Nanotechnology and nanomaterials is an emerging field of research attempting to tailor the (macroscopic) behaviour of materials and devices to meet certain needs for physical properties. Mechanical strength, electrical conductivity, thermal properties and also functional properties such as energy conversion between photons and electrical energy (solar cells) and electrochemical properties (carbon capture) are all examples of material properties one wants to control, and in order to control these aspects, both understanding and manipulation at the micro-/nano scale are required [3]. Being able to manipulate materials down to the nano-scale would contribute to the different challenges the modern society are facing. Nanomaterials and -technology may offer more efficient and economically sustainable materials in the search for environmentally sustainable energy sources, more efficient transistors, nanoparticles that can help removing contaminants in water where potable water is lacking, as well as both medical treatment and -imaging [4, 5].

Aluminium is one of the largest exports in Norway, only behind oil/gas and fish, and due to its versatility and manipulable mechanical properties, price and weight, aluminium and aluminium alloys are a big field of research. With an increasing interest in using aluminium for electrification purposes in e.g. automotive industry and wires, there are great incentives of finding alloys with a high electrical conductivity and good mechanical performance. Whilst pure and undeformed, aluminium has a high electrical conductivity and can suitably serve as an electrical conductor. However, the mechanical properties of pure aluminium are less suited for practical use. It is easily bendable and fails to reach the performance required to meet the needs in e.g. wires and chassis (cars, airplanes etc.). Aluminium alloys however, are often used in all the aforementioned [6]. Here the strength is improved by

alloying with e.g. magnesium, lithium, silicon and zinc, depending on the required properties [7]. This even holds at elevated temperatures, making the alloys suitable as an all-round material for many applications. One of the drawbacks of such alloying is that the electrical conductivity is lowered. Today, copper is most widespread as an electrical conductor. It is fairly cheap, environmentally sustainable and has a high electrical conductivity relative to other candidates. On a world basis however, there is a supply shortage resulting in an increasing price [8]. It is not enough copper, taking the massive electrification of todays society into account. This creates a further incentive of finding new materials in the future to meet the climate goals regarding electrification. Tailoring aluminium in such a way that it can replace copper as an electrical conductor would then be beneficial regarding price, recyclability and weight. With an industrial demand of both high conductivity and improved mechanical properties, a suitable trade off between the two must then be found. As alloying form microconstituents with a size typically ranging down to the nano-scale, improved understanding of the electrical effects they have at precisely the nano-scale is then required to perform this tailoring [9].

Scanning transmission electron microscope/microscopy (STEM) is a technique offering sub-nanometer resolution, and utilizing this resolution for the aforementioned analysis would be a great technique to measure a *local*- or a (sub-)nanometer electrical conductivity. This would require additional spectroscopy methods and in this project, it will be investigated whether electron energy loss spectroscopy (EELS) can be utilized for this purpose. As the electrical properties of pure aluminium are well documented, the applicability of this method will be tested for pure aluminium. A temperature range spanning from room temperature to 500 °C has been considered, affecting both the macroscopic conductivity as well as the expected signal. It has been investigated whether or not the technique can reproduce the macroscopic electrical conductivity as function of temperature. The Drude model and the dielectric formulation within solid state physics have been utilized as the physical foundation relating the conductivity to the EELS signal [10, 11].

Similar experiments have been carried out by J.-L. Verger-Gaugry and P.Guyot on crystals and quasicrystals in AlMn and AlMnSi in the mid. 1980's [12]. However, their results had uncertainties at the same order of magnitude as their calculated mean values, and as far as the author can see from the literature, there has not been conducted many studies since then. It is thus assumed that advances in TEM instrumentation and EEL spectroscopy could allow further investigation of such measurements.

The report will firstly introduce the physical and instrumental aspects of the project as well as a short introduction to some deconvolution techniques. Secondly, the methods that has been used during sample preparation, EELS acquisition and post-processing will be presented. Lastly, the results will be presented and discussed.

# Chapter 2

## Theory

This chapter will firstly introduce the underlying physics justifying the attempt to use EELS as means for calculating the conductivity, being a short introduction of the differential cross section within quantum mechanics, the Drude model followed by the dielectric formulation of many-particle systems. This will be followed by the theory of a plasma and *plasmons*, being a quantum of plasma oscillation. Secondly, the chapter introduces shortly some theory behind the working principle of EELS, and lastly, some deconvolution techniques will be presented. The reader is expected to be acquainted with transmission electron microscope/microscopy (TEM) and to have basic knowledge in solid state physics. If so is not the case, several textbooks are available covering these topics, such as Refs. [1, 11, 13] and Refs. [10, 14] for the physics and instrumentation, respectively. Unless otherwise specified, the theory presented in this chapter rests on these textbooks.

### 2.1 Scattering

As the electron beam hits the sample, multiple scattering events may occur. These are mainly divided into *elastic*- and *inelastic* scattering. Elastic scattering originates from interactions between the incoming electron and the electrostatic field from the nuclei and bound electrons in the sample. Although *some* energy gets lost in the process in terms of excitation of phonons, the energy loss in these interactions are typically not resolved in a TEM [10]. Inelastic scattering on the other hand, results in a detectable and measurable energy loss and originates from electron-electron interactions. These scattering events include excitation of the electrons present in the sample, i.e. absorption. Additionally, and more relevant for our purposes, the beam

electrons in a TEM may interact with the free electron gas as a collective. This will be further explained in section 2.4 when discussing plasmas and plasmons.

Central to the scattering mechanisms is the differential cross section

$$\frac{d\sigma}{d\Omega} \quad (2.1)$$

representing the probability of an incident electron being scattered per unit solid angle by an atom. Here,  $\sigma$  is the cross section of scattered electrons at a solid angle  $\Omega$ . Finding this cross section is itself a large field of study within quantum mechanics and will in the following be limited to the work of Ritchie [15] and Lindhart [16].

## 2.2 Drude Model

The Drude model was developed by Paul Drude in 1900 and is an application of the kinetic theory of gases. Originally, the Drude theory coarsely assumed that the electrons behaved in a classical manner (i.e. billiard balls), colliding with positively charged and immobile spheres [1, 17]. In later work, the theory has been refined due to new discoveries such as the atomic model, phonons and other scattering events. The Drude model is now considering free electrons in a lattice of positive ions. The only scattering event that is considered is the interactions between the free electrons and the ion lattice, thus neglecting electron-electron interactions. If the free electrons have a momentum per electron  $\mathbf{p}(t)$  at a time  $t$  and a momentum  $\mathbf{p}(t + dt)$  at some later time  $dt$ , the fraction of electrons having “collided” with the ion lattice will be  $dt/\tau$ , where  $\tau$  is the *relaxation time* of the system. With no external forces, the electron gas will (if not at rest) come to rest, and the net momentum of the free electrons will become zero. If an external force is present however, the electrons will gain net momentum between the collisions. The resulting equation of motion for the electron gas then becomes

$$\frac{d\mathbf{p}}{dt} = \frac{-\mathbf{p}}{\tau} - e\mathcal{E}, \quad (2.2)$$

where the external force has been taken to originate from an external electrical field  $\mathcal{E}$ , and  $e$  is the elementary charge. The effect of single-electron-lattice interactions therefore acts as a dampening force to the free electron gas as a collective. For a many particle electron gas, the conductivity,  $\sigma$  (not to be

confused with the cross section in section 2.1), is defined through

$$\mathbf{j} = -ne\langle \mathbf{v} \rangle = \sigma \mathcal{E} \quad (2.3)$$

where  $\mathbf{j}$  is the current density,  $n$  is the electron density and  $\langle \mathbf{v} \rangle$  is the mean velocity of the electrons. Taking  $x, \mathcal{E} \propto \exp(-i\omega t)$  where  $\omega$  is the frequency, gives

$$\sigma(\omega) = \sigma_0 \frac{1 + i\tau\omega}{1 + \omega^2\tau^2} \quad (2.4)$$

where  $\sigma_0$  is defined through the relationship

$$\sigma_0 = \frac{ne^2\tau}{m}, \quad (2.5)$$

whereas  $m$  is the electron mass. The current  $\sigma_0$  is also known as the *Drude conductivity*, and represents the DC conductivity of the material.

## 2.3 Dielectric formulation

The dielectric function,  $\epsilon_r(\mathbf{k}, \omega)$ , describes the dielectric response of the material to an external electromagnetic field and is therefore also known as the *dielectric response function*. Its dependency on both frequency and wave vector,  $\mathbf{k}$ , is strong, though the long wavelength limit  $\epsilon_r(\mathbf{k} \rightarrow 0, \omega)$  will be assumed sufficient for our purposes describing the collective excitation of the electron gas [11].

By representing a transmitting electron as a point charge with coordinate  $\mathbf{r}$  and velocity  $\nu$ , satisfying the Poisson equation,

$$\epsilon_0 \epsilon_r(\mathbf{k}, \omega) \nabla^2 \phi(\mathbf{r}, t) = e\delta(\mathbf{r}, t), \quad (2.6)$$

where  $\epsilon_0$  is the vacuum permittivity,  $\phi(\mathbf{r}, t)$  is the electrostatic potential and  $\delta(\mathbf{r}, t)$  is the Dirac delta function, R. H. Ritchie showed that the differential cross section in Equation 2.1 for small angles can be written as

$$\frac{d^2\sigma}{d\Omega dE} \propto \frac{1}{\nu^2} \Im \left( \frac{-1}{\epsilon_r(k, E)} \right) \left( \frac{1}{\Theta^2 + \Theta_E^2} \right), \quad (2.7)$$

where  $\Theta_E = E/(\gamma m\nu^2)$  is a characteristic angle,  $\gamma$  is a relativistic factor,  $\nu$  is the velocity of the incident electrons,  $E$  is the energy,  $\Theta$  is the scattering angle and  $\Im(-1/\epsilon(k, E))$  is the *energy-loss function*, providing a complete description of the medium the transversing electron is going through [10, 15].

Two of Maxwell's equations read

$$\nabla \cdot \mathbf{D} = \rho_{\text{ext}} \quad ; \quad \nabla \cdot \mathcal{E} = \rho_{\text{tot}}/\epsilon_0 \quad (2.8\text{a,b})$$

where  $\mathbf{D}$  and  $\mathcal{E}$  are the electric displacement and electric field, respectively, whilst  $\rho_{\text{ext}}$  and  $\rho_{\text{tot}}$  is the external charge density and the total charge density, respectively. The dielectric function is defined through

$$\mathbf{D} = \epsilon_0 \mathcal{E} + \mathbf{P} = \epsilon_r \epsilon_0 \mathcal{E}, \quad (2.9)$$

where  $\mathbf{P}$  is the polarization density, and will in the long wavelength limit take the form

$$\epsilon_r(\omega) \approx \epsilon_{\omega \rightarrow \infty} \left[ 1 - \frac{\bar{\omega}_p^2}{\omega^2} \right]. \quad (2.10)$$

Here  $\bar{\omega}_p$  is the *plasma frequency* accounting for the ion core distribution at high frequency:

$$\bar{\omega}_p^2 = \omega_p^2 / \epsilon_{\omega \rightarrow \infty} = \frac{1}{\epsilon_{\omega \rightarrow \infty}} \frac{n e^2}{\epsilon_0 m}, \quad (2.11)$$

$n$  and  $m$  are the electron density and the electron mass, respectively, whilst  $e$  is the elementary charge and  $\epsilon_{\omega \rightarrow \infty}$  is the dielectric contribution from the ion background at high frequency. Inserting plane waves into the wave equation for a non-magnetic medium

$$\mu_0 \ddot{\mathbf{D}} = \nabla^2 \mathcal{E}, \quad (2.12)$$

$\mu_0$  being the vacuum permeability, and using Equation 2.9 provides the dispersion relation, being a function of  $\mathbf{k}$  and  $\omega$

$$\epsilon_r(\mathbf{k}, \omega) \frac{\omega^2}{c^2} = \|\mathbf{k}\|^2. \quad (2.13)$$

From here, it is clear that for  $\epsilon_r < 0$ ,  $\|\mathbf{k}\|$  becomes imaginary and the waves are attenuated exponentially. The plasma frequency therefore acts as a lower cutoff for wave propagation in a plasma, in which  $\mathbf{P} = -\epsilon_0 \mathcal{E}$ . At this frequency, the dielectric function equates to 0 and a collective longitudinally polarized wave mode of the system gets excited, whereas transversal modes are excited at higher frequencies. This is schematically illustrated in Figure 2.1 showing both the dispersion relation in Equation 2.13, the optical limit as well as a shaded region indicating that waves cannot propagate at frequencies lower than  $\omega_p$ . Not considering spin, the equation of motion of

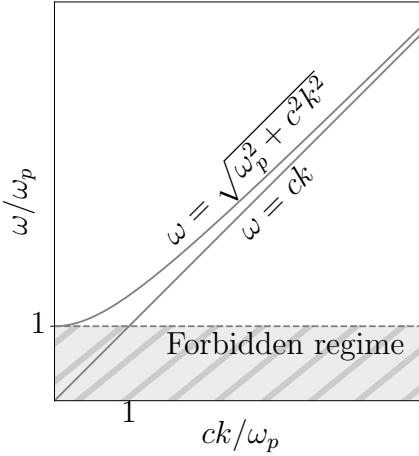


Figure 2.1: the dispersion relation for the plasma waves as well as the optical limit. No propagation occurs in the shaded region.

the plasma is given by

$$\begin{aligned} \mathbf{P} = -ne\langle \mathbf{x} \rangle &= \frac{-ne^2\tau}{m(\omega^2\tau - i\omega)} [\boldsymbol{\mathcal{E}} + \langle \mathbf{v} \rangle \times \mu \mathbf{H}] \\ &\approx \frac{-ne^2\tau}{m(\omega^2\tau - i\omega)} \boldsymbol{\mathcal{E}} \end{aligned} \quad (2.14)$$

where  $\mathbf{H}$  is the magnetic field strength, and the latter holds true in the long wavelength limit  $\langle \mathbf{v} \rangle = \hbar\langle \mathbf{k} \rangle/m \rightarrow 0$ , being the same form as for the Drude theory. Conclusively, the dielectric formalism in the long wavelength limit predicts that the Drude conductivity defined in Equation 2.4 can be investigated thorough the properties of the collective electron motion at the plasma frequency.

Jumping back to the differential cross section derived by Ritchie and introduced in Equation 2.7, the energy-loss function in this long wavelength limit is given by

$$-\Im \left( \frac{1}{\epsilon_r(k=0, E)} \right) = \frac{E(\Delta E) E_p^2}{(E^2 - E_p^2)^2 + (E\Delta E)^2}, \quad (2.15)$$

known as the *inverted Drude-Sellmeier function* [13]. In this equation,  $E_p = \hbar\omega_p$  is the plasma energy,  $\Delta E = \hbar/\tau$  is the full width at half maximum

(FWHM) and  $\hbar$  is Planck's reduced constant, thus connecting the EELS measurements to the conductivity given in Equation 2.18, and will for  $E_p/\Delta E \gg 1$  take the shape of a Lorentzian function [18] (see Equation 2.17).

## 2.4 Plasmons

As briefly mentioned, at the plasma frequency, the material sets up a depolarizing field to the electric field  $\mathbf{P} = -\epsilon_0 \mathcal{E}$ . The total dielectric displacement then equates to 0, and Equation 2.2 and Equation 2.14 takes the form of a damped harmonic oscillator, and a longitudinal collective travelling wave of the conduction electron gas is excited. Similarly to a phonon being a quantum of lattice oscillation, a *plasmon* is a quantum of electron density oscillation with energy  $E_p = \hbar\omega_p$ . A plasmon can be excited in a metallic thin film by a transversing electron as presented in Equation 2.6. The negative charge of the electron couples to the electron gas, and the transversing electron will then have lost energy equal to an integer multiple of the plasmon energy  $\hbar\omega_p$ . Equation 2.14 then takes the form

$$n\ddot{\mathbf{x}} = -\frac{n}{\tau}\dot{\mathbf{x}} - \frac{n^2 e^2}{\epsilon_0 m} \mathbf{x}. \quad (2.16)$$

For  $\omega_p\tau \gg 1$ , the normalized solution to the differential equation takes the form of a Lorentzian in reciprocal space:

$$L(\omega) = \frac{1}{2\pi} \frac{\Gamma}{(\omega - \omega_p)^2 + \Gamma^2/4}, \quad (2.17)$$

where  $\Gamma = 1/\tau$  is the FWHM [18]. The quantities in Equation 2.5 can then be found in an EELS experiment through Equation 2.11, such that

$$\sigma_0 = \frac{\epsilon_0}{\hbar} \frac{E_p^2}{\Delta E}. \quad (2.18)$$

## 2.5 Metals and thermal dependencies

With increasing temperature, the electron-phonon interactions increase, thus decreasing the conductivity in a metal. According to Matthiesen's rule, the relationship between the total relaxation time in the system is related to the

individual scattering events through

$$\begin{aligned}\frac{1}{\tau_{\text{tot}}} &= \sum_i \frac{1}{\tau_i} \\ &= \frac{1}{\tau_{\text{im}}} + \frac{1}{\tau_{\text{ph}}}\end{aligned}\quad (2.19)$$

where the sum is taken over all scattering events [11]. In this equation,  $\tau_{\text{im}}$  and  $\tau_{\text{ph}}$  are the relaxation times related to impurities and phonon excitation, respectively. For a metal at room temperature,  $\tau_{\text{im}} \ll \tau_{\text{ph}}$ , and thus Equation 2.19 can be approximated as

$$\tau_{\text{tot}} \approx \tau_{\text{ph}}. \quad (2.20)$$

The collision rate with phonons is proportional to the concentration of phonons, which increases linearly above the Debye temperature  $\Theta_D$ . With this temperature dependency, the Drude conductivity in Equation 2.5 is expected to decrease as  $1/T$  for  $T > \Theta_D$  [11], commonly written on the form:

$$\sigma_0(T)^{-1} = [\sigma_0(T_0)(1 - \alpha\Delta T)]^{-1}, \quad (2.21)$$

where  $\alpha$  is the temperature coefficient,  $T_0$  is a reference temperature with known conductivity and  $\Delta T = T - T_0$  is the temperature difference between the measured temperature and the reference temperature.

A general property of metals is that they expand when the temperature increases. This happens in a linear manner for our temperature ranges, and it can readily be shown from linear expansion of solids and Equation 2.11 that the plasmon energy is expected to decrease linearly with increasing temperature [19]:

$$E_p(T) = E_p(T_0) \left[ 1 - \frac{3}{2}\beta\Delta T \right]. \quad (2.22)$$

In this equation,  $\beta$  is the coefficient of linear expansion and is taken here as a material constant as enlisted in Table 1.

## 2.6 Electron Energy Loss Spectroscopy

Electron energy loss spectroscopy measures the kinetic energy of initially monoenergetic electrons after specimen interaction. The beam electrons hits the sample at what is known as the *convergence angle*,  $\alpha$ . After the beam

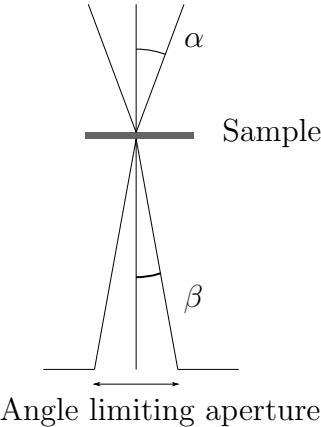


Figure 2.2: schematic of convergence angle,  $\alpha$ , collection angle,  $\beta$ , and the angle limiting aperture in EELS experiments.

electrons have transmitted and interacted with the sample, they travel further into the EEL spectrometer through an angle limiting aperture. As the word suggests, this aperture filters out all electrons having scattered more than the *collection angle*,  $\beta$ , see Figure 2.2. The electrons are then bent by magnetic forces in a magnetic prism. The curvature of the electrons are then depending on their velocity as they enter the magnetic field:

$$R = \frac{\gamma m \nu}{e B}. \quad (2.23)$$

In Equation 2.23,  $m$  is the electron mass,  $e$  the electric charge,  $\nu$  the electron velocity after specimen interaction,  $B$  the strength of the magnetic field, and  $\gamma$  a relativistic factor. The EELS detector therefore detects the energy loss of the electrons after the magnetic prism based this curvature. For our purposes the energy loss regime around the plasmon peak in aluminium is of interest, and will be the main focus.

## 2.7 Deconvolution techniques

As for all experimental techniques, measurable features are limited by the resolution of the instrument. For our purposes, the term *resolution* can be two different things. The *spatial* resolution is how small features in real space that is detectable. That is, if microconstituents are present in the material, it is the spatial resolution that determines how small clusters/clouds/precipitates one is able to distinguish. The *energy* resolution is a measure of

how well the instrument can distinguish different energy loss peaks. If two physically distinct peaks are closer to each other than the energy resolution, the two peaks will appear as one, wider peak. In general, the resolution affects the detected signal by a convolution between the actual signal and the *instrumental response function*,  $R(E)$ , determining the resolution. Further on, the response function will be referred to as the energy resolution and how the instrumental setup responds to no inelastic scattering, i.e. the zero-loss peak (ZLP). Setting the response function to have unit area, the ZLP will take the form

$$Z(E) = I_0 R(E), \quad (2.24)$$

where  $I_0$  is the incoming intensity,  $Z(E)$  is the acquired signal of the ZLP and  $E$  is the energy loss of an electron. The energy resolution of the system is then given by the FWHM of  $Z(E)$ , in which for our purposes will be approximated by curve-fitting with a Gaussian and a SplitGaussian (see section 3.4). All other features, such as the plasmon peak, will also be affected by the response function through a convolution [10]:

$$J(E) = Z(E) * S(E). \quad (2.25)$$

Here  $J(E)$  is the signal recorded by the instrument and  $S(E)$  is the signal from the physical processes. Directly solving this equation for  $S(E)$  by Fourier transforms would greatly amplify noise, and a *reconvolution function* is required to maintain a sufficiently high signal-to-noise-ratio (SNR). The *Fourier-ratio* method solves this by reconvolving Equation 2.25 with a zero-centered unit area Gaussian,  $G(E)$ , as a reconvolution function:

$$G(E) * S(E) = \mathcal{F}^{-1} \{ g(\nu) j(\nu) / z(\nu) \}. \quad (2.26)$$

Here,  $g(\nu)$ ,  $r(\nu)$  and  $j(\nu)$  are the Fourier transforms of  $G(E)$ ,  $R(E)$  and  $Z(E)$ , respectively. This gives an additional free parameter, as the FWHM of the Gaussian reconvolution function can be chosen.

*Fourier-log deconvolution* attempts not only to remove the effect of a finite resolution, but also to eliminate multiple scattering. Assuming that the scattering events follow Poisson statistics, the Fourier transform of the single scattering signal,  $s$ , takes the form [10]

$$s(\nu) = I_0 \log(j(\nu) / z(\nu)). \quad (2.27)$$

Once again, solving this by taking the inverse Fourier transform is very prone to noise. The Fourier-log deconvolution technique handles this issue by using  $Z(E)$  or  $G(E)$  as reconvolution functions, termed *zero-loss modifier* and

*Gaussian modifier*, respectively, and Equation 2.27 then takes the form:

$$z(\nu)s(\nu) = I_0 z(\nu) \log\left(\frac{j(\nu)}{z(\nu)}\right) \quad ; \quad g(\nu)s(\nu) = I_0 g(\nu) \log\left(\frac{j(\nu)}{z(\nu)}\right). \quad (2.28a,b)$$

The zero-loss modifier will not alter the energy resolution, but remove effects from multiple scattering, whilst the Gaussian modifier will remove multiple scattering effect and can both compensate for a potential asymmetry of  $Z(E)$  and improve the resolution by setting the FWHM lower than that of  $Z(E)$ . This can also introduce other artefacts, especially for noisy data.

## 2.8 Lindhart Extension to higher $k$ -values

This section is heavily based on Ref [13]. The derivations can be found in this reference and will thus not be repeated here.

Lindhart extended the inverted Drude-Sellmeier function (Equation 2.15) to not only yield in the optical limit, but also for small scattering vectors. He found that the plasmon energy  $E_p$  is itself depending on the momentum transfer through

$$E_p(k) = E_p(0) + \frac{\alpha' \hbar^2}{m} k^2 \quad (2.29)$$

where

$$\alpha' = \frac{3}{5} \frac{E_F}{E_p(0)} \quad (2.30)$$

and  $E_F$  is the Fermi energy [10]. He further considered both the real and imaginary part of the dielectric function:  $\epsilon = \epsilon_1 + i\epsilon_2$ . For the real part, a similar expression as in Equation 2.10 is found, but the imaginary part divides the dielectric function into four parts. For our purposes, one of these parts are of interest, and restricts collective excitations of plasmons to frequencies lower than the parabola

$$\hbar\omega = \frac{\hbar^2}{2m}(k^2 + 2kk_F), \quad (2.31)$$

where  $k_F$  is the Fermi wave vector. This is shown in Figure 2.3 along with Equation 2.29. The parabola marks the boundary of where  $\epsilon_2$  equates to 0, whereas for  $\hbar\omega$  greater than this parabola, no electron-hole excitations of wave vector  $k$  is possible. This creates a curve in the  $(k, \omega)$ -plane where  $\epsilon_1$  and  $\epsilon_2$  are simultaneously zero, up to a cut-off wave vector  $k_c$  where they intersect. This curve is shown in blue in the aforementioned figure, showing

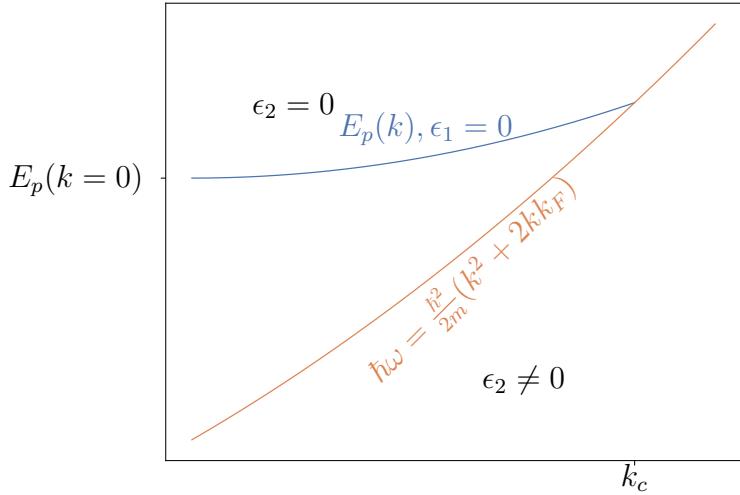


Figure 2.3: schematic of the  $(k, \omega)$ -plane for the dielectric function as derived in the Lindhart model. The orange curve separates the regions where  $\epsilon_2 = 0$  and where  $\epsilon_2 \neq 0$

the relevant parts (for our purposes) of the  $(k, \omega)$ -plane of the dielectric function.

# Chapter 3

## Method

### 3.1 Sample preparation

For this project, a 25 µm thick aluminium foil from Goodfellow with a purity of 99.999%<sup>1</sup> was disk punched to have a diameter of approximately 3 mm. The disk was then ion milled in a GATAN PIPS II Model 695, where the beam energy and time are summarized in Table 3.1. The sample was then investigated in a JEOL JEM 2100 TEM operating at 200 kV before extracting two suitable regions in a Helios G5 Plasma Focused ion beam (FIB) and placing them on a P.T.H.TS.1 DENS chip allowing in-situ heating of the sample. During transfer to the DENS chip, the regions of interest were illuminated as little as possible after precursor deposition steps to minimize contamination and damage. An image taken during FIB preparation is shown in Figure 3.1. The two samples were then plasma cleaned in a Fischione Model 1020 plasma cleaner for 1 min 30 s immediately before EELS analysis.

### 3.2 EELS acquisition

EELS data was acquired using a Gatan Continuum Gatan imaging filter (GIF) with a CMOS detector on a X-FEG monochromated Thermo Fisher Scientific Spectra 300 operated with an acceleration voltage of 300 kV. In order to narrow the ZLP and thus improve the energy resolution, the exposure time was set to 10 ms and the scans were recorded

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<sup>1</sup>Commercially available at <https://www.goodfellow.com/p-a100-f1-000131/aluminium-foil> (04.06.2023)

Table 3.1: beam energy and duration used in the ion milling step.

Beam energy [eV]	Time [min]
5	140
4	60
3.5	30
3	30
2.5	30
2	30
1.5	30
1	30
0.5	60

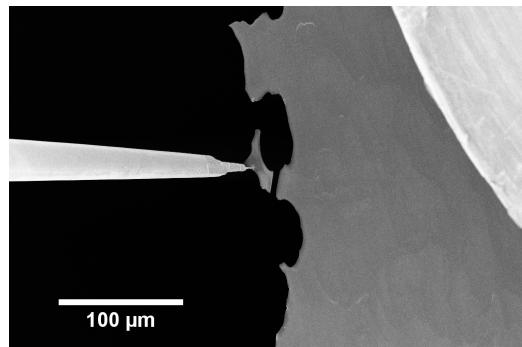


Figure 3.1: sample during transfer to DENS chip in a FIB. Image acquired by Per Erik Vullum

10 times in DualEELS™ mode with a collection aperture of 100 mrad. The energy ranges for the two simultaneous EEL recordings were set to -2 to 14 eV and 8 to 24 eV with a dispersion of 10 meV per channel. The temperature of the sample was raised sequentially, starting off at room temperature and up to 500 °C as is plotted in Figure A.1. The imaging system was unstable and the resolution was prone to errors as the temperature of the sample was increased. Slight adjustments of the alignment were therefore performed before the first scan at each temperature.

The monochromator is an essential part of the EELS system when energy resolutions of  $\lesssim 100$  meV are required. Instead of being limited by the uncertainty in the electron energy from the electron gun, the electrons are filtered out before entering the rest of the TEM column. The improved energy resolution comes at the expense of a lowered intensity (i.e. fewer electrons).

DualEELS™ makes it possible to acquire two different energy loss intervals simultaneously. It is often used for recording both the ZLP (or low-loss) and a core-loss, where the latter is typically at energies in the keV range. For our purposes, it allows recording both the ZLP and the rest of the low-loss regime, thus avoiding overexposure from the ZLP, affecting the signal from the plasmon peak. Additionally, the energy dispersion can be lowered, in which for our purposes was 10 meV.

### 3.3 Data handling

Prior to all further data handling, the EEL signals were aligned such that the ZLP was placed at 0 eV using HyperSpy [20]. The ZLP used in Equation 2.28a,b was taken to be defined for all electrons having lost  $\leq 2$  eV. The different signal lines that were curve fitted to the signal was a Lorentzian (Equation 2.17), SplitLorentzian (Equation 2.17, see also section 3.4), and the inverted Drude-Sellmeier function<sup>2</sup> (Equation 2.15) using SciPy [21]. Both weighted and unweighted fits were performed, and the energy loss interval for optimization was between 13.5 and 16.5 eV. All pixels were curve fitted in this manner, followed by filtering out the pixels differing from the mean value by one standard deviation for all free parameters<sup>3</sup>. The remaining pixels were segmented using scikit-image [22] and added together to improve the SNR. Segmenting the pixels was based on their measured

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<sup>2</sup>Results not shown as it was indistinguishable from the Lorentzian fit.

<sup>3</sup>Obvious outliers affecting the mean value to a sufficiently great extent was filtered out manually.

plasmon energy. Background subtraction was performed by subtracting the median value of the signal between 10-12 eV, before deconvolving the signal with the methods presented in section 2.7 using NumPy [23]. As means to obtain information regarding  $Z(E)$  in Equation 2.24, the ZLP was curve fitted to both a Gaussian and a SplitGaussian function (see section 3.4). A linear regression was made for all different approaches, selecting the  $R^2$  values as means to estimate how well the signal lines fit the signal [24].

All code can be found in Appendix D.

### 3.4 Note

As briefly mentioned, several signal lines have been curve fitted to the spectra. Due to asymmetry of the peak(s), it is has been necessary to define distributions with two different widths: one width for energy losses lower than the center, and one width for energy losses greater than the center. A signal line  $S(E, E_0, \Delta E)$  will be tagged with “Split” if such a procedure has been performed. For the Split signal line, the function will take the form

$$\text{Split}S(E, E_0, \Delta E_1, \Delta E_2) = I_0 \begin{cases} \frac{1}{S_{E_0^-}} S(E, E_0, \Delta E_1) & , E \leq E_0 \\ \frac{1}{S_{E_0^+}} S(E, E_0, \Delta E_2) & , E > E_0, \end{cases} \quad (3.1)$$

where  $E_0$  is the center,  $\Delta E_{1,2}$  are FWHM and  $1/S_{E_0^-}$  and  $1/S_{E_0^+}$  are defined as

$$S_{E_0^-} = S(E = E_0, E_0, \Delta E_1) \quad ; \quad S_{E_0^+} = S(E = E_0, E_0, \Delta E_2)$$

making sure that the signal line is continuous at  $E = E_0$ .

Additionally, other attempts has been made in order to acquire the most information regarding the plasmon peak. This includes:

1. Curve fit to the derivative of a SplitLorentzian. This was done as it was hypothesized that the curve fit optimizer would prioritize the energy intervals that was rapidly increasing/decreasing and could thus lead to additional information regarding the FWHM. Additionally, the background level has been assumed to be a constant, and any potential misjudgement of this level would for this procedure be eliminated. It turned out to be equivalent to a weighted fit, though more prone to noise.

2. Altering the energy interval for optimization based on the estimated center. This was done to isolate the energy interval closer to the center, and in particular for the energies between the center and 500 meV lower than the center. This did not make any particular difference.
3. Adjusting the expression for the (Split)Lorentzian such that a change in FWHM would alter the integral and not the maximum value. This function would be mathematically equivalent to the original signal lines, but one could theoretically find the FWHM based on the ratio between the integral and the maximum value. This gave estimates that had a very high variance, and therefore considered inadequate for this purpose.
4. Finding the numerical inflection point for  $E < E_p$  as well as the center value. Their difference can be used to find the FWHM, but this was (unsurprisingly) very prone to noise.
5. Inserting a peak with a FWHM derived from the macroscopic conductivity and setting hard restrictions for how far away its center value could be compared to previous estimates. The measured area of this peak was dropped to 0 unless restricted by the optimizer. If restricted, the optimizer would choose the lowest possible area for the peak, and the results from this procedure was therefore considered unphysical.
6. Centering the signal on the plasmon peak instead of the ZLP. This gave a bigger FWHM, likely due to that the SNR for the ZLP was far better, resulting in a less accurate estimate for the plasmon peak compared to the ZLP (also being evident that DualEELS™ is working properly).

All these attempts did not provide any additional information regarding the peak, and their results will thus not be presented.

# Chapter 4

## Results and Discussion

As a brief overview, an angular dark field (ADF) image of the sample is shown in Figure 4.1 along with an image taken in the FIB during sample preparation. It shows that the sample is mostly pure aluminium, though contaminated on one of its edges<sup>1</sup>. Examples of how the EEL signal looked like is plotted in Figure 4.2, and shows that summation over multiple pixels vastly enhances the SNR. A figure showing the energy resolution defined in section 2.7 is shown in Figure B.1 and shows that the FWHM of the ZLP was around  $90 - 100$  meV, with the exception of scans taken at room temperature and at  $120^\circ\text{C}$ . Although some deviations of the results depending on the data handling methods were present, the overall tendencies and order of magnitudes were rather similar. The main differences were if the Split-Lorentzian method was used compared to the Lorentzian and if the signal was weighted or not during curve fitting. Unless otherwise specified, the data handling method presented has 10 pixels per segment, was Fourier-log deconvolved with a Gaussian modifier with a reconvolution function having a FWHM of 90 meV, and was not weighted in the curve fitting procedure.

The overall conductivity was measured to be far lower than the macroscopic conductivity, measuring a conductivity of  $(4.3 \pm 0.2) \cdot 10^6 \text{ S m}^{-1}$  at room temperature with a temperature dependency as is shown in Figure 4.3 along with a linear regression of Equation 2.21. It is already here clear that the conductivity at room temperature has not been reproduced, as the measured value and the macroscopic value differ with a factor of 9. Further, the temperature dependency for the plasmon energy is shown in Figure 4.4 and was measured to be  $14.94 \pm 0.01$  eV at room temperature, declining with approximately  $0.5 \text{ meV K}^{-1}$ . Solving Equation 2.11 for the electron density

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<sup>1</sup>This region was not included in the further data handling

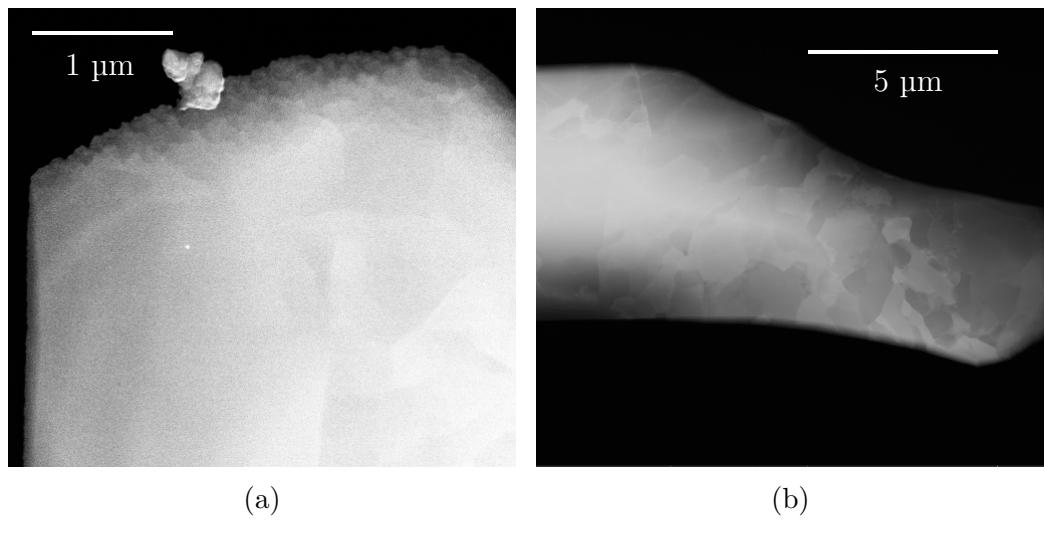


Figure 4.1: a) ADF image of the sample, and b) an image from the FIB during FIB-milling. The images a,b) were acquired by Simon Fairclough and Per Erik Vullum, respectively.

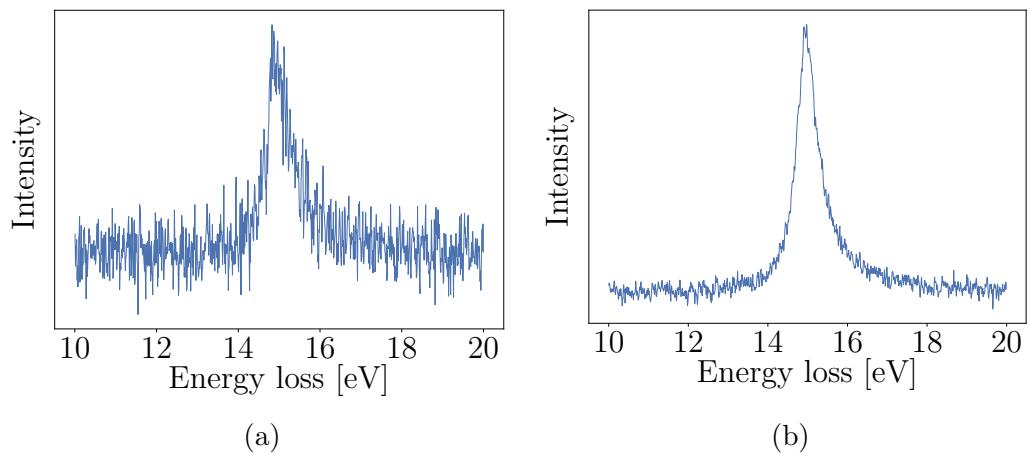


Figure 4.2: examples of unprocessed EELS signals from a) one pixel and b) segment with 100 pixels per segment.

gives  $n = 18 \cdot 10^{22} \text{ cm}^{-1}$ . The decline in plasmon energy for increasing temperature can be explained through Equation 2.22, measuring a temperature coefficient of approximately  $23 \cdot 10^{-6} \text{ K}^{-1}$ . Comparing both this decline and the electron density to the values enlisted in Table 1 reveals that these values are well within an acceptable accuracy, and the plasmon energy can therefore be considered to be found. This is also congruent with existing literature, showing both the same plasmon energy [25] and temperature dependency [19]. It must here be stressed that the plasmon energy at room temperature was slightly different for the SplitLorentzian and the Lorentzian curve fit, whereas the latter is more similar to previous EELS experiments.

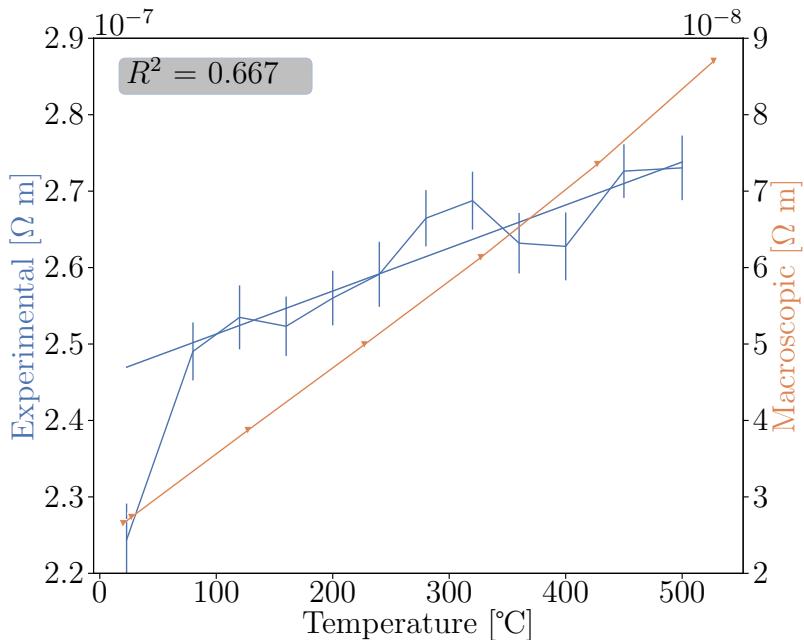


Figure 4.3: measured resistivity (inverse conductivity) of aluminium as function of temperature compared to the macroscopic resistivity. Note that the offset for the two colors are different, but the scale is similar. For the regression analysis, the data at room temperature were left out, see also Figure 4.5b. The macroscopic data points are taken from Ref. [2]

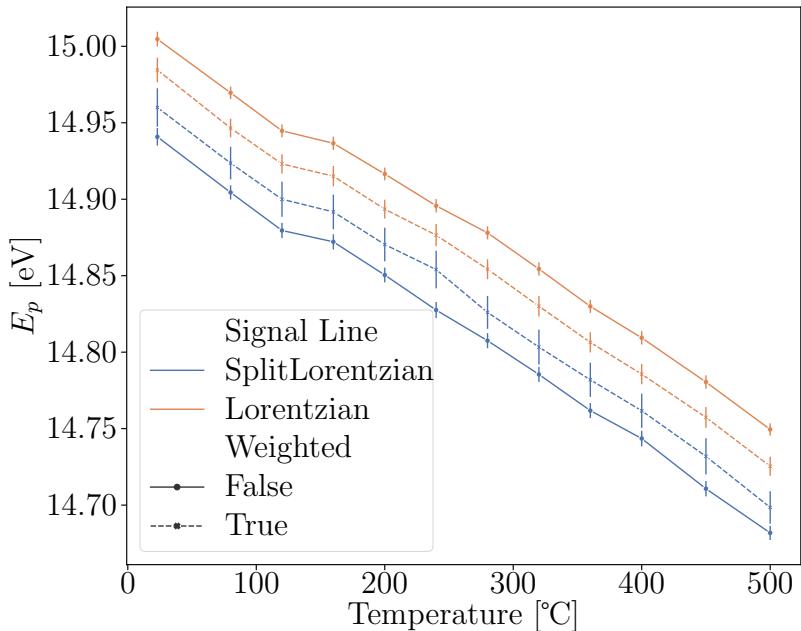


Figure 4.4: temperature dependency for the measured plasmon energy for the SplitLorentzian and the Lorentzian function.

It is therefore evident that it is the FWHM of the plasmon peak that is causing such a big discrepancy between our measurements and the macroscopic conductivity. Predicted at room temperature, the FWHM should have been around 80 meV, while our results gave a FWHM greater than 450 meV, and should have increased to about 240 meV at 500 °C. Although the instrumental resolution of  $\lesssim 100$  meV would have failed to find the true FWHM at room temperature, there is a strong discrepancy between the measured FWHM and a convolution between the theoretical Lorentzian peak and the ZLP. This is shown in Figure 4.5a, showing the FWHM for  $E < E_p$  and  $E > E_p$  for the SplitLorentzian. In addition, Figure 4.5b shows how well the SplitLorentzian fit is compared to the Lorentzian fit, making it clear that the asymmetry of the SplitLorentzian function makes a better fit than the Lorentzian function. With this in mind, it is clearly evident that the plasmon has been broadened towards higher energies. One of the premises for the method to work, was that the EEL spectra was due to plasmon excitations in the long wavelength limit,  $k \rightarrow 0$ . As the collection aperture was set to 100 mrad and the acceleration voltage was 300 kV, the incoming electrons had an initial wave vector of  $k_0 = 3191 \text{ nm}^{-1}$ , allowing all momentum

transfers up to  $k \approx k_0\theta_{\text{collection aperture}} = 319.1 \text{ nm}^{-1}$  into the spectrometer. The dielectric scattering cross section in Equation 2.7 would predict a rapid decline in intensity for higher values of  $k$  as  $\Theta_E$  is of the order 0.1 mrad [10] and the dispersion relation in Figure 2.1 predicts that the plasma oscillations quickly becomes transversal rather than longitudinal. However, experimental  $k$ -resolved evidence show that there are great contributions from higher  $k$ -values [26]. Batson and Silcox mapped the  $(k, \omega)$  plane of the plasmon response, mapping out both the plasmon energy, linewidth and intensity [27]. They measured the coefficient,  $\alpha'$ , in Equation 2.29 to be  $\alpha' = 0.38 \pm 0.02$ , the FWHM to increase for increasing  $k$ , and that the contribution from  $k \not\rightarrow 0$  is non-negligible even for  $k > k_c = 1.3 \text{ \AA}^{-1}$  (see Figure 2.3, numerical value also given in Ref. [27]) with an incident beam of 75 keV. They found that the differential cross section in Equation 2.7 is insufficient to describe the observed EEL spectra as it only considers single scattering and that multiple scattering inflicts the measurements. This was for both their and our purposes eliminated through the Fourier-log deconvolution techniques, and the corrections to the FWHM this yielded in this experiment is shown in Figure C.1 for the SplitLorentzian. The Fourier-log deconvolution technique does not however, take the contribution for higher  $k$ -values into account, and a broadening to higher energy losses will still be present. Although Batson and Silcox indeed obtained an intensity map of the momentum transfer- and energy loss-dependency for the plasmon sufficient to explain our broadening and asymmetry in a qualitative way, the accuracy of their measurements were limited as their equipment was not up to modern standards<sup>2</sup>. It is therefore evident that more sophisticated theoretical framework is needed to explain the contribution to our measured energy loss. Such calculations have been made by Ferrell, based on Hartree-Fock wave function analysis [28] and is shown in Figure 4.6 along with a sharp cut-off approximation taken from Ref. [10]. His calculations showed that despite the dispersion relation in Figure 2.1, there are indeed non-negligible contributions for  $k < 0.74k_F = 1.3 \text{ \AA}^{-1}$ , where  $k_F$  is the Fermi wave vector and the numerical value is taken from Ref. [2]. The cutoff wave vector will in the Lindhart extension to higher  $k$ -values gives a plasmon energy at approximately 18 eV. This contribution explains both the asymmetry and the broadening of the peak for  $E > E_p$ . However, as the contributions from scattering events not being in the long-wavelength limit strictly contributes to  $E > E_p(k = 0)$  the energy losses from  $E < E_p$  is not

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<sup>2</sup>or in their words: “(...) for the most part, the accuracy of the intensity measurements reported here surpasses the ability of the draftsman and printer to reproduce them on plots of this nature” [27].

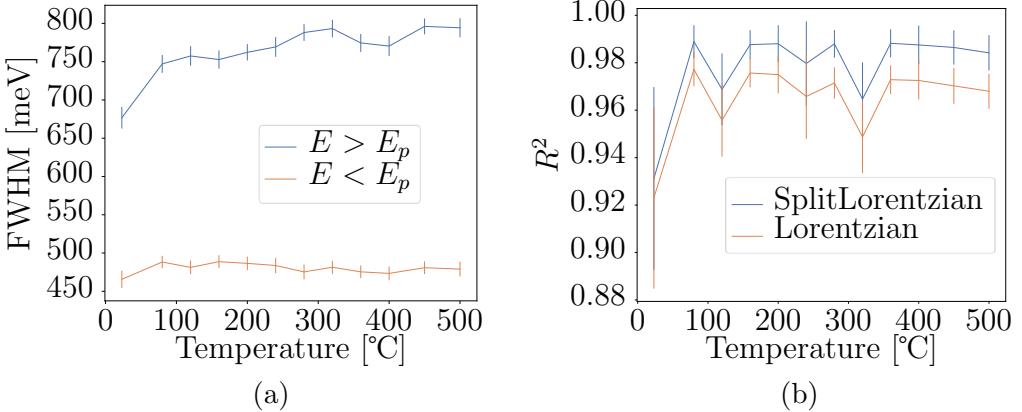


Figure 4.5: temperature dependency for a) the FWHM for the Split-Lorentzian, and b) the  $R^2$  value for the SplitLorentzian and Lorentzian function.

completely explained. The additional intensity from higher values of  $k$  close to- but not *sufficiently* close to 0 still contribute with a rather high intensity. Their contribution may be enough to shift the maximum of the acquired energy loss peak to a higher energy than the plasmon energy at  $k = 0$ . As the spectrometer used in these experiments are not  $k$ -resolved, finding the intensity distribution as function of  $k$  is unavailable, and our resulting energy-loss spectra are effectively integrated over all  $k$ -values not filtered out by the collection aperture. Figure 4.7a shows a resulting prediction of how a signal would appear based on Ferrell's calculations along with experimental data taken at 500 °C for a segment of 100 pixels. This figure makes it clear that the scattering cross section from Ferrell is indeed promising to explain the asymmetry of the plasmon peak for  $E > E_p$ , but fails for  $E < E_p$ . Close to the cutoff frequency, Ferrell's calculations have been experimentally verified by Schmuser [26] for an electron beam with an energy of 40 keV, though the aforementioned paper from Batson and Silcox suggested that the intensity was extending to even higher values of  $k$  with a beam energy of 75 keV. As far as the author has seen through literature, Ferrell's calculations have not been experimentally verified for sufficiently low  $k$ . If the scattering cross section predicts an increasing intensity up to a wave vector  $k_E$  as suggested by the sharp cut-off approximation as seen as the dashed line in Figure 4.6, the intensity in the experimental peak would have a peak value shifted to even higher energies that would lead to a measured broadening also for  $E < E_p$ .

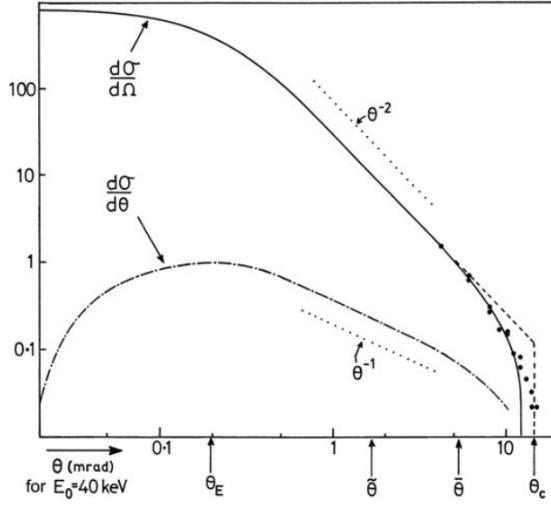


Figure 4.6: the differential cross section as calculated by Ferrell [28] (solid) and using a sharp cutoff approximation [10]. The figure is taken from Ref. [10].

This is illustrated in Figure 4.7b, and shows more promising features, also for  $E < E_p$ . More mathematically:

$$S_{\text{measured}}(E) = \int_{k>0} S(k, E) dk \neq S(k=0, E), \quad (4.1)$$

where  $S(E)$  is the measured peak from *Equation 2.27*, even for  $E < E_{p,\text{measured}}$ . If so is the case, this would explain why the plasmon energy estimated from an unweighted fit does not coincide with the peak values obtained from a weighted fit on a quantitative level<sup>3</sup>: the intensity for  $E < E_{p,\text{measured}}$  would simply not be dominated by scattering events from the optical limit.

Another interpretation has been suggested by Smith and Segall, arguing that the frequency  $\omega_p$  of the bulk plasmon emerges as consequence of both intraband and interband transitions [29], and that the observed plasmon energy in the acquired signal is a consequence of that both contributions show the same asymptotic behavior at this frequency. The interband contributions are not a part of the analysis and their contribution would prevent the analysis as it has been an underlying assumption from *Equation 2.21* that the observed scattering events are due to increased amount of phonon

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<sup>3</sup>Quantitative analysis showed a noisy and unpredictable values of the energy shift, though of the order 20 – 30 meV, being only 2-3 times the energy dispersion.

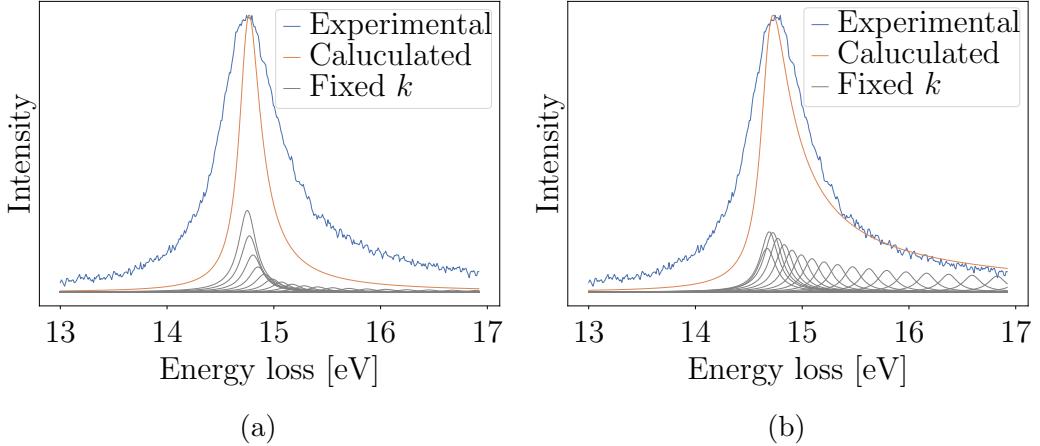


Figure 4.7: estimated shape of plasmon given a Lorentzian plasmon shape as calculated by a) Ferrell and b) the sharp cut-off approximation [10, 28]. The orange line is the sum of all the grey lines (intensity of grey lines not to scale with the other two).

excitations. This interpretation is supported by K. Sturm [30, 31], arguing that the Drude addition to the linewidth is drowning in the interband transitions due to Umklapp scattering. This would explain the shortcomings of both Ferrells calculations and the sharp cut-off approximation as is apparent in Figure 4.7. Further, Smith and Segall found the intraband component to have a frequency of  $\Omega_p = 12.5 \pm 0.3$  eV with a relaxation time of  $\tau = (1.06 \pm 1.2) \cdot 10^{-14}$  s, resulting in a conductivity at room temperature differing from the macroscopic electrical conductivity with approximately 10% [29].

Regardless of what has caused the broadening of the peak, it is clear that the FWHM indeed is too big to reproduce the scattering time in the Drude model. Equation 2.22 was derived on the premise that the dominating contribution to the overall relaxation was the Drude contribution. Although this is not congruent with Smith and Segall as well as Sterms work, the increase in relaxation time for the Drude contribution should increase the overall FWHM. As the reference conductivity in Equation 2.22 has been measured to be incorrect, it would be more sensible to choose the form  $aE + b$ , where Figure 4.3 suggests that there are similarities to macroscopic measurements. The slope in the aforementioned form from our experiments is a factor of 1.8 lower than its macroscopic counterpart. In terms of the Lindhart model, this can be explained through that the contribution is drowning in the effects of

higher  $k$ -values, whereas the same argument holds for the interband explanation. The fact that there is a temperature dependency however, is likely due to a decreasing conductivity. That the FWHM for  $E > E_p$  is the side showing a more prominent temperature dependency can be explained through Figure 4.7, where the extension to higher  $k$ -values makes more impact for  $E > E_p$ , as well as Batson and Silcox findings of an increasing FWHM for increasing  $k$ .

# Chapter 5

## Conclusion

### 5.1 Conclusion

This project has attempted to find the DC-conductivity of pure aluminium using EELS. The signal line from the plasmon peak over a temperature range from room temperature to 500 °C has been investigated for this purpose. Overall, the plasmon peak has been measured to have a FWHM far too big for this purpose, and a thorough explanation has not been found. It has been speculated that this can be due to contributions from plasmon excitation involving sufficiently high momentum transfer and interband transitions. It has also been shown that the plasmon peak is too asymmetric to assert a Lorentzian shape, likely originating from the contributions from higher  $k$ -values. A temperature dependency of the width of the plasmon peak has been found, though more research is required to outline an explicit relation. Furthermore, it has been found that the plasmon energy is around 15 eV at room temperature, decreases linearly with increasing temperature. This tendency is well explained through linear thermal expansion of solids.

### 5.2 Future work

As has become clear throughout chapter 4, broadening effects due to  $k$ -dependency of the plasmon peak is likely a big reason of why the conductivity has not been retrieved. A more thorough understanding of especially the differential cross section as function of both energy loss and momentum transfer is then required for further analysis. This can be done similarly to what Batson and Silcox did [27], but now with instruments up to modern

standards, such as in Refs. [32, 33]. If the experiments are to be repeated without  $k$ -resolved EELS, again attempting to find the plasmon behaviour in the long wavelength limit, it would be necessary to lower the acceleration voltage and have a smaller collection aperture in order to prevent contributions from higher  $k$ -values. Other papers, as referenced in chapter 4 suggest that this would not be sufficient as they argue that the plasmon is dominated by interband transitions.

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# Bibliography

- [1] Neil W. Ashcroft and N. David Mermin. *Solid state physics*. Saunders College Publishing, 1976.
- [2] David R Lide. *CRC handbook of chemistry and physics*, volume 85. CRC press, 2004.
- [3] Ahmed Kadhim Hussein. Applications of nanotechnology in renewable energies-a comprehensive overview and understanding. *RENEWABLE & SUSTAINABLE ENERGY REVIEWS*, 42:460–476, FEB 2015.
- [4] Parshant Kumar Sharma, Shraddha Dorlikar, Pooja Rawat, Vidhu Malik, Nishant Vats, Manu Sharma, Jong Soo Rhyee, and Ajeet Kumar Kaushik. 1 - nanotechnology and its application: a review. In Kamil Reza Khondakar and Ajeet Kumar Kaushik, editors, *Nanotechnology in Cancer Management*, pages 1–33. Elsevier, 2021.
- [5] Lina Wang, Mavd P.R. Teles, Ahmad Arabkoohsar, Haoshui Yu, Kamal A.R. Ismail, Omid Mahian, and Somchai Wongwises. A holistic and state-of-the-art review of nanotechnology in solar cells. *Sustainable Energy Technologies and Assessments*, 54:102864, 2022.
- [6] Tolga Dursun and Costas Soutis. Recent developments in advanced aircraft aluminium alloys. *MATERIALS & DESIGN*, 56:862–871, APR 2014.
- [7] Ankitkumar K. Shriwas and Vidyadhar C. Kale. “ impact of aluminium alloys and microstructures on engineering properties-review ”. 2016.
- [8] Lee Ying Shan. There isn't enough copper in the world — and the shortage could last till 2030. <https://www.cnbc.com/2023/02/07/there-isnt-enough-copper-in-the-world-shortage-could-last-.html>, February 2023.

- [9] Tusr Ranjan Soren, Ramanuj Kumar, Isham Panigrahi, Ashok Kumar Sahoo, Amlana Panda, and Rabin Kumar Das. Machinability behavior of aluminium alloys: A brief study. *Materials Today: Proceedings*, 18:5069–5075, 2019. 9th International Conference of Materials Processing and Characterization, ICMPC-2019.
- [10] R.F. Egerton. *Energy-Loss Instrumentation*. Springer US, Boston, MA, 2011.
- [11] Charles Kittel. *Introduction to solid state physics*. John Wiley & Sons, 8 edition, 2004.
- [12] J.-L. Verger-Gaugry and P. Guyot. PLASMON ELECTRON LOSS SPECTROSCOPY AND ELECTRICAL CONDUCTIVITY AT 300 K OF CRYSTALS AND QUASICRYSTALS IN AlMn AND AlMnSi. *Journal de Physique Colloques*, 47(C3):C3–477–C3–483, 1986.
- [13] Giuseppe Grosso and Giuseppe Pastori Parravicini. Chapter 7 - excitons, plasmons, and dielectric screening in crystals. In Giuseppe Grosso and Giuseppe Pastori Parravicini, editors, *Solid State Physics (Second Edition)*, pages 287–331. Academic Press, Amsterdam, second edition edition, 2014.
- [14] David B. Williams and C. Barry Carter. *The Transmission Electron Microscope*. Springer US, Boston, MA, 2009.
- [15] R. H. Ritchie. Plasma losses by fast electrons in thin films. *Phys. Rev.*, 106:874–881, Jun 1957.
- [16] Jens Lindhard. On the properties of a gas of charged particles. *Dan. Vid. Selsk Mat.-Fys. Medd.*, 28:8, 1954.
- [17] P. Drude. Zur elektronentheorie der metalle. *Annalen der Physik*, 306(3):566–613, 1900.
- [18] Lukas Novotny and Bert Hecht. *Principles of Nano-Optics*. Cambridge University Press, 2 edition, 2012.
- [19] Hiroyuki Abe, Masami Terauchi, Ryuichi Kuzuo, and Michiyoshi Tanaka. Temperature Dependence of the Volume-Plasmon Energy in Aluminum. *Journal of Electron Microscopy*, 41(6):465–468, 12 1992.

- [20] Francisco de la Peña, Eric Prestat, Vidar Tonaas Fauske, Pierre Burdet, Jonas Lähnemann, Petras Jokubauskas, Tom Furnival, Magnus Nord, Tomas Ostasevicius, Katherine E. MacArthur, Duncan N. Johnstone, Mike Sarahan, Joshua Taillon, Thomas Aarholt, pquinn dls, Vadim Mi-gunov, Alberto Eljarrat, Jan Caron, Carter Francis, T. Nemoto, Timothy Poon, Stefano Mazzucco, actions user, Nicolas Tappy, Niels Cau-taerts, Suhas Somnath, Tom Slater, Michael Walls, Florian Winkler, and Håkon Wiik Ånes. hyperspy/hyperspy: Release v1.7.3, October 2022.
- [21] Pauli Virtanen, Ralf Gommers, Travis E. Oliphant, Matt Haberland, Tyler Reddy, David Cournapeau, Evgeni Burovski, Pearu Peterson, Warren Weckesser, Jonathan Bright, Stéfan J. van der Walt, Matthew Brett, Joshua Wilson, K. Jarrod Millman, Nikolay Mayorov, Andrew R. J. Nelson, Eric Jones, Robert Kern, Eric Larson, C J Carey, İlhan Polat, Yu Feng, Eric W. Moore, Jake VanderPlas, Denis Laxalde, Josef Perktold, Robert Cimrman, Ian Henriksen, E. A. Quintero, Charles R. Harris, Anne M. Archibald, Antônio H. Ribeiro, Fabian Pedregosa, Paul van Mulbregt, and SciPy 1.0 Contributors. SciPy 1.0: Fundamental Algorithms for Scientific Computing in Python. *Nature Methods*, 17:261–272, 2020.
- [22] Stéfan van der Walt, Johannes L. Schönberger, Juan Nunez-Iglesias, François Boulogne, Joshua D. Warner, Neil Yager, Emmanuelle Gouil-lart, Tony Yu, and the scikit-image contributors. scikit-image: image processing in Python. *PeerJ*, 2:e453, 6 2014.
- [23] Charles R. Harris, K. Jarrod Millman, Stéfan J. van der Walt, Ralf Gommers, Pauli Virtanen, David Cournapeau, Eric Wieser, Julian Tay-lor, Sebastian Berg, Nathaniel J. Smith, Robert Kern, Matti Picus, Stephan Hoyer, Marten H. van Kerkwijk, Matthew Brett, Allan Hal-dane, Jaime Fernández del Río, Mark Wiebe, Pearu Peterson, Pierre Gérard-Marchant, Kevin Sheppard, Tyler Reddy, Warren Weckesser, Hameer Abbasi, Christoph Gohlke, and Travis E. Oliphant. Array pro-gramming with NumPy. *Nature*, 585(7825):357–362, September 2020.
- [24] Karl Pearson. Vii. note on regression and inheritance in the case of two parents. *proceedings of the royal society of London*, 58(347-352):240–242, 1895.

- [25] Heinz Raether. *Excitation of plasmons and interband transitions by electrons*, volume 88. Springer Tracts in Modern Physics, Springer, 1980.
- [26] Peter Schmüser. Anregung von volumen- und oberflächenplasmaschwingungen in al und mg durch mittelschnelle elektronen. *Zeitschrift für Physik*, 180:105–126, 1964.
- [27] P. E. Batson and J. Silcox. Experimental energy-loss function,  $\text{Im}[-\frac{1}{\epsilon}(q, \omega)]$ , for aluminum. *Phys. Rev. B*, 27:5224–5239, May 1983.
- [28] Richard A. Ferrell. Characteristic energy loss of electrons passing through metal foils. ii. dispersion relation and short wavelength cutoff for plasma oscillations. *Phys. Rev.*, 107:450–462, Jul 1957.
- [29] D. Y. Smith and B. Segall. Intraband and interband processes in the infrared spectrum of metallic aluminum. *Phys. Rev. B*, 34:5191–5198, Oct 1986.
- [30] Kurt Sturm. Pseudopotential theory of the width of the long wavelength plasmon in simple metals. *Zeitschrift für Physik B Condensed Matter*, 25(3):247–253, 1976.
- [31] Kurt Sturm. Pseudopotential theory of the k-dependent plasmon line width in simple metals. *Zeitschrift für Physik B Condensed Matter*, 28(1):1–7, 1977.
- [32] Ryosuke Senga, Kazu Suenaga, Paolo Barone, Shigeyuki Morishita, Francesco Mauri, and Thomas Pichler. Position and momentum mapping of vibrations in graphene nanostructures. *Nature*, 573(7773):247–250, 2019.
- [33] Hikaru Saito, Hugo Lourenço Martins, Noémie Bonnet, Xiaoyan Li, Tracy C. Lovejoy, Niklas Dellby, Odile Stéphan, Mathieu Kociak, and Luiz Henrique Galvão Tizei. Emergence of point defect states in a plasmonic crystal. *Phys. Rev. B*, 100:245402, Dec 2019.

# Appendix A

## Temperature plot

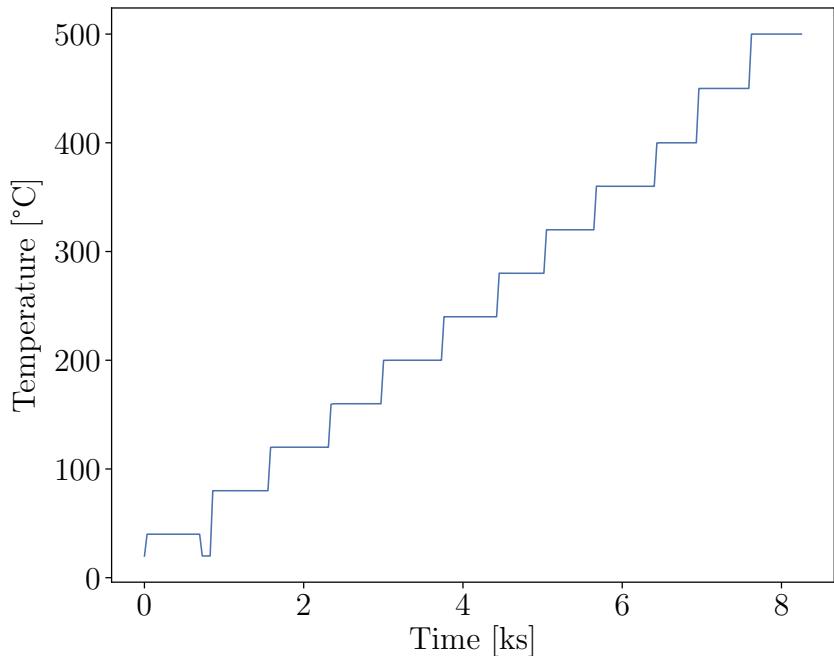


Figure A.1: measured temperature log of the sample during TEM operation.

# Appendix B

## Resolution

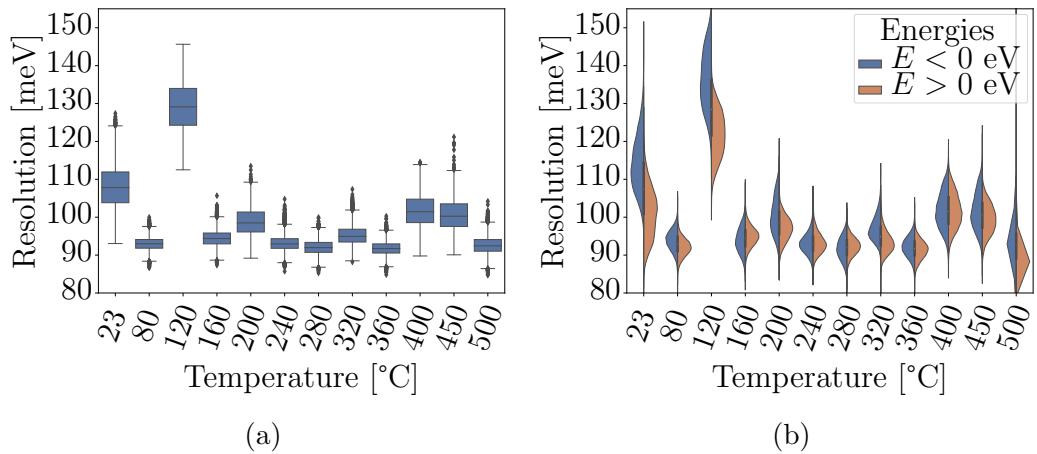


Figure B.1: FWHM of the ZLP after fitting the signal to a) Gaussian and b) SplitGaussian. Note that the temperature axis is not equidistant.

# Appendix C

## Effect of Deconvolution for SplitLorentzian

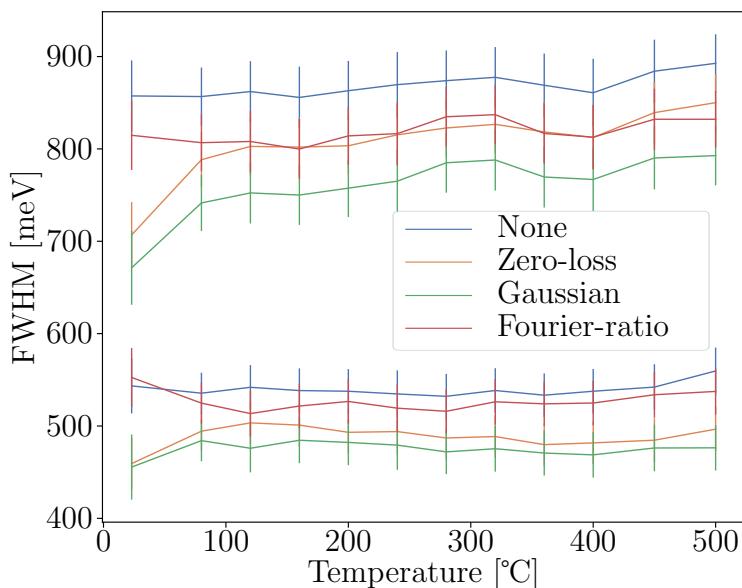


Figure C.1: measured temperature dependency for the signal for the different deconvolution techniques for a SplitLorentzian

# Appendix D

## Code

```
1 from numpy import pi
2 """This file contains the parameters that can be used for
   initialization of the curve-fitting algorithm
3
4 The order of the dictionary is essential. Only one example is
   shown.
5 """
6
7 class Param:
8     def __init__(self, value, bmin, bmax):
9         self.value = value
10        self.bmin = bmin
11        self.bmax = bmax
12
13
14 drudec      = Param(14.8,12,18)
15 drudefwhm = Param(0.5 ,0 ,1)
16 drudea      = Param(2e5 ,0 ,1e9)
17
18
19 drude = {'center' : drudec,
20           'fwhm'    : drudefwhm,
21           'area'     : drudea}
22
23
24
25 """ This file contains wrappers for the different signal lines.
   These include Gaussian, Lorentzian, Voigt, and more."""
26
27 from abc import ABC
28 import numpy as np
```

```

29 import sympy
30 import scipy.special
31 import logging
32
33
34
35
36 dx  = '+0.01'
37 dx_ = '-0.01'
38 def prod_id(x):
39     while True:
40         yield chr(ord('a')+x)
41         x+=1
42 id_ = prod_id(0)
43
44 gauss_expr          = '(area * 2 * sqrt(log(2))/(fwhm*sqrt(pi))'
45             * exp(-4*(log(2)) * (x -center)**2/fwhm**2) ) '
46 loren_expr           = '(area / pi * fwhm / 2 / ((x -center)'
47             **2 + fwhm**2/4))'
48 # gauss_expr          = 'area * exp(-4*(log(2)) * (x-center)**2/fwhm'
49 #             **2 ) '
50 # loren_expr          = 'area * ((x-center)**2 + fwhm**2/4)'
51 # voigt_expr          = 'area* real(Gfwhm / (4*sqrt(pi*log(2))) * scipy.'
52 #             special.wofz(2*sqrt(log(2)) /Gfwhm * (x-center+1j*Lfwhm/2) )'
53 #             ) '
54 voigt_expr           = '(area * real(2*sqrt(log(2))/sqrt(pi)*1/
55 #             Gfwhm * wofz(2*sqrt(log(2)) /Gfwhm * (x-center+1j*Lfwhm/2) ) )'
56 # splitvoigt_expr = 'area * (((1-frac)*'+gauss_expr.replace('
57 #             area','1')).replace('fwhm','fwhm1') + ' + frac*'+loren_expr.
58 #             replace('area','1')).replace('fwhm','fwhm1')+')* (x<=center*1)
59 #             '\
60 #             '+ ((1-frac)*'+gauss_expr.replace('
61 #             area','1')).replace('fwhm','fwhm2') + ' + frac*'+loren_expr.
62 #             replace('area','1')).replace('fwhm','fwhm2')+')* (x >center*1)
63 #             ) '
64 splitvoigt_expr      = 'area * ( 1/((1-frac)*'+gauss_expr.
65 #             replace('area','1')).replace('fwhm','fwhm1').replace(' x ',''
66 #             center'+dx_) + ' + frac*'+loren_expr.replace('area','1').
67 #             replace('fwhm','fwhm1').replace(' x ','center'+dx_)+') * ((1-
68 #             frac)*'+gauss_expr.replace('area','1')).replace('fwhm','fwhm1'
69 #             ) + ' + frac*'+loren_expr.replace('area','1')).replace('fwhm',
70 #             'fwhm1')+')* (x<=center*1) '\
71 #             '+ 1/((1-frac)*'+gauss_expr.replace('
72 #             area','1')).replace('fwhm','fwhm2').replace(' x ','center'+dx_
73 #             ) + ' + frac*'+loren_expr.replace('area','1')).replace('fwhm',
74 #             'fwhm2').replace(' x ','center'+dx_)+') * ((1-frac)*'+

```

```

gauss_expr.replace('area','1').replace('fwhm','fwhm2') + '+
frac'+loren_expr.replace('area','1').replace('fwhm','fwhm2')
+')*(x>center*1) )'
54 volume_expr      = '(area * x * fwhm * center**2 / ((x**2-
center**2)**2 + (x*fwhm)**2 ) )'
55 step_expr        = '(x>start*1)'
56 inter_expr       = '((x>start)-(x>stop)*1)'
57 power_expr       = 'area*(x-center)**(-k)'
58 # power_expr     = 'area*(x-center * ((x-0.3)>center*1))**(-k)'
#Note!
59 eps_1_expr        = 'area * (1-center**2 / (x**2 + fwhm**2))'
60 eps_2_expr        = 'area * fwhm*center**2/(x*(x**2+fwhm**2))'
,
61 experimental_expr = '(1-area* (center**2-x**2)/((x**2-center
*x**2)**2+x**2*fwhm**2))'
62 lognormal_expr   = 'area * 1/(x*sigma*sqrt(2*pi))*exp(-(log(
x)-center)**2/(2*sigma**2))'    # see e.g. wikipedia
63 dho_expr          = 'area * sigma / (pi*(1-exp(-center/gamma)
)) * (1/((x-center)**2 +sigma**2) - 1/((x+center)**2 +
sigma**2)) '
64 # experimental_expr = 'area*(1+ (1+1j*fwhm/x)*() / '
65 fano_expr          = '(area * abs(fwhm/2- background*exp(1j*
phi)*(1j*( x -center)-fwhm/2) )**2 / (( x -center)**2 +
fwhm**2/4))'
66 # derivativeL_expr = '-area * 16 * ( x - center) * fwhm / (
pi *(4*( x -center)**2+fwhm**2)**2 )'
67 # derivativeL_expr = '-area * 16 * ( x - center) * fwhm1 /
(pi *(4*( x -center)**2+fwhm1**2)**2 )**2 + -area
* 16 * ( x - center) * fwhm2 / (pi *(4*( x -center)**2+
fwhm2**2)**2 )'
68 # derivativeL_expr = '-area * 16 * ( x - center) * fwhm1 /
(pi *(4*( x -center)**2+fwhm1**2)**2 )*(( x <= center)*1)-
area * 16 * ( x - center) * fwhm2 / (pi *(4*( x -center)**2+
fwhm2**2)**2 )*(( x > center)*1)'
69 derivativeL_expr  = '(-area * 16 * ( x - center) * fwhm1
/ (pi *(4*( x -center)**2+fwhm1**2)**2 )*(( x <= center)*1)
-area * fwhm2/fwhm1 * 16 * ( x - center) * fwhm2 / (pi *(4*( x -
center)**2+fwhm2**2)**2 )*(( x > center)*1) )'
70
71 splitLorent_expr = '( 1/(2*pi) * area * fwhm1**2 / (( x -
center)**2 + fwhm1**2/4) * (x <= center ) + 1/(2*pi) *
area * fwhm2**2 / (( x -center)**2 + fwhm2**2/4) * ( x >
center ) )'
72 splitGauss_expr  = '(area *( exp(-4*(log(2)) * ( x -center)
**2/fwhm1**2 ) *( x <=center) + exp(-4*(log(2)) * ( x -
center)**2/fwhm2**2 ) *( x > center) ))'#A is max value
73

```

```

74 def step_func(x, start):
75     return ((x-start) > 0)*1
76 def define_on_interval(x, start, stop):
77     return step_func(x,start) - step_func(x,stop)
78
79
80 def gaussian(x,center,fwhm,area):
81     sigma = fwhm/(2*np.sqrt(2*np.log(2)))
82     return area * 1/(sigma*np.sqrt(2*np.pi)) * np.exp(-(x-center)
83         )**2/(2*sigma**2))
84
85
86 def lorentzian(x, center, fwhm, area):
87     return area * 1/np.pi * fwhm/2 * 1/((x-center)**2 + (fwhm/2)
88         **2)
89
90 def double_lorentzian(x,center,fwhm1,fwhm2,area):
91     return lorentzian(x,center,fwhm1,area)*(x<=center) +
92     lorentzian(x,center,fwhm2,area*fwhm2/fwhm1)*(x>center)
93
94
95 def voigt(x:np.array,c,FWHM,gamma,area):
96     gamma /= 2
97     sigma = FWHM/(2*np.sqrt(2*np.log(2)))
98     z = (x - c + 1j * gamma) / (np.sqrt(2)*sigma)
99     V = scipy.special.wofz(z)/(np.sqrt(2.0*np.pi)*sigma)
100
101
102 class SignalLine(ABC):
103     def __init__(self,parameters:dict, expr:str = None, id=None,
104         symbols=None) -> None:
105         super().__init__()
106         self.parameters = parameters
107         self.expr = expr
108         self.id = id
109         self.symbols = symbols
110     def append_step_func(self,start:float,stop:float = None):
111         """Defines the interval the signal line is defined at"""
112         if stop is None:
113             self.expr += '*' + step_expr.replace('start',str(
114                 start))
115         else:

```

```

114         self.expr += '*' + inter_expr.replace('start', str(
115             start)).replace('stop', str(stop))
116
117 class Lorentzian(SignalLine):
118     def __init__(self, parameters:dict, name:str = None) -> None
119     :
120         if name is None:
121             self.id = next(id_)
122         else:
123             self.id = name
124         expr = loren_expr.replace('area', 'area'+self.id).
125             replace('center', 'center'+self.id).replace('fwhm', 'fwhm'+
126             self.id)
127         self.symbols = {}
128         for sym, param in zip(sympy.symbols('center fwhm area'.
129             replace('area', 'area'+self.id).replace('center', 'center'+
130             self.id).replace('fwhm', 'fwhm'+self.id)), parameters.keys()):
131             :
132                 self.symbols[sym] = parameters[param]
133             super().__init__(parameters, expr, self.id, self.symbols
134         )
135
136 class Gaussian(SignalLine):
137     def __init__(self, parameters:dict, name:str = None) -> None
138     :
139         if name is None:
140             self.id = next(id_)
141         else:
142             self.id = name
143         expr = gauss_expr.replace('area', 'area'+self.id).
144             replace('center', 'center'+self.id).replace('fwhm', 'fwhm'+
145             self.id)
146         self.symbols = {}
147         for sym, param in zip(sympy.symbols('center fwhm area'.
148             replace('area', 'area'+self.id).replace('center', 'center'+
149             self.id).replace('fwhm', 'fwhm'+self.id)), parameters.keys()):
150             :
151                 self.symbols[sym] = parameters[param]
152             super().__init__(parameters, expr, self.id, self.symbols
153         )
154
155 class Voigt(SignalLine):
156     def __init__(self, parameters:dict, name:str = None) -> None
157     :

```

```

145     if name is None:
146         self.id = next(id_)
147     else:
148         self.id = name
149     expr = voigt_expr.replace('area', 'area'+self.id).
150     replace('center','center'+self.id).replace('Gfwhm', 'Gfwhm'+
151     self.id).replace('Lfwhm', 'Lfwhm'+self.id)
152     self.symbols = {}
153     for sym, param in zip(sympy.symbols('center Gfwhm Lfwhm
154     area').replace('area', 'area'+self.id).replace('center',
155     center'+self.id).replace('Gfwhm', 'Gfwhm'+self.id).replace(
156     'Lfwhm', 'Lfwhm'+self.id)), parameters.keys()):
157         self.symbols[sym] = parameters[param]
158     super().__init__(parameters, expr, self.id, self.symbols
159 )
160
161
162
163
164
165 class SplitVoigt(SignalLine):
166     """Takes in the center position, fwhm1 (<= center), fwhm2
167     (>= center), frac"""
168     def __init__(self, parameters:dict, name:str = None) -> None
169     :
170         if name is None:
171             self.id = next(id_)
172         else:
173             self.id = name
174         expr = splitvoigt_expr.replace('area', 'area'+self.id).
175         replace('center','center'+self.id).replace('fwhm1', 'fwhm1'+
176         self.id).replace('fwhm2', 'fwhm2'+self.id).replace('frac', 'frac'+
177         self.id)
178         self.symbols = {}
179         for sym, param in zip(sympy.symbols('center fwhm1 fwhm2
180         frac area').replace('area', 'area'+self.id).replace('center',
181         center'+self.id).replace('fwhm1', 'fwhm1'+self.id).replace(
182         'fwhm2', 'fwhm2'+self.id).replace('frac', 'frac'+self.id)),
183         parameters.keys()):
184             self.symbols[sym] = parameters[param]
185         super().__init__(parameters, expr, self.id, self.symbols
186 )
187
188
189
190
191
192
193
194 class VolumePlasmon(SignalLine):
195     """Inverted Drude-Sellmeier"""
196
197     def __init__(self, parameters:dict, name:str = None) -> None
198     :
199         if name is None:
200             self.id = next(id_)
```

```

175     else:
176         self.id = name
177         expr = volume_expr.replace('area', 'area'+self.id).
178             replace('center','center'+self.id).replace('fwhm', 'fwhm'+
179                 self.id)
180         self.symbols = {}
181         for sym, param in zip(sympy.symbols('center fwhm area').
182             replace('area', 'area'+self.id).replace('center','center'+
183                 self.id).replace('fwhm', 'fwhm'+self.id)), parameters.keys()):
184             :
185                 self.symbols[sym] = parameters[param]
186             super().__init__(parameters, expr, self.id, self.symbols
187             )
188
189
190 class DHO(SignalLine):
191     def __init__(self, parameters:dict, name:str = None) -> None
192     :
193         if name is None:
194             self.id = next(id_)
195         else:
196             self.id = name
197             expr = dho_expr.replace('area', 'area'+self.id).replace(
198                 'center','center'+self.id).replace('sigma', 'sigma'+self.id).
199                 replace('gamma', 'gamma'+self.id)
200             self.symbols = {}
201             for sym, param in zip(sympy.symbols('center sigma gamma
202                 area').replace('area', 'area'+self.id).replace('center',
203                     'center'+self.id).replace('sigma', 'sigma'+self.id).replace(
204                         'gamma', 'gamma'+self.id)), parameters.keys()):
205                 self.symbols[sym] = parameters[param]
206             super().__init__(parameters, expr, self.id, self.symbols
207             )
208
209
210 class Fano(SignalLine):
211     def __init__(self, parameters:dict, name:str = None) -> None
212     :
213         if name is None:
214             self.id = next(id_)
215         else:
216             self.id = name
217             expr = fano_expr.replace('area', 'area'+self.id).replace(
218                 'center','center'+self.id).replace('fwhm', 'fwhm'+self.id).
219                 replace('phi', 'phi'+self.id).replace('background', ' '
220                     'background'+self.id)
221             self.symbols = {}

```

```

204     for sym, param in zip(sympy.symbols('center fwhm area
phi background').replace('area', 'area'+self.id).replace(
'center','center'+self.id).replace('fwhm', 'fwhm'+self.id).
replace('phi', 'phi'+self.id).replace('background', 'background'+self.id)), parameters.keys()):
205         self.symbols[sym] = parameters[param]
206     super().__init__(parameters, expr, self.id, self.symbols
)
207
208
209 class DerivativeLorentzian(SignalLine):
210     def __init__(self, parameters:dict, name:str = None) -> None
:
211         if name is None:
212             self.id = next(id_)
213         else:
214             self.id = name
215         expr = derivativeL_expr.replace('area', 'area'+self.id).
replace('center','center'+self.id).replace('fwhm1', 'fwhm1'+self.id).replace('fwhm2', 'fwhm2'+self.id)
216         self.symbols = {}
217         for sym, param in zip(sympy.symbols('center fwhm1 fwhm2
area').replace('area', 'area'+self.id).replace('center','center'+self.id).replace('fwhm1', 'fwhm1'+self.id).replace('fwhm2', 'fwhm2'+self.id)), parameters.keys()):
218             self.symbols[sym] = parameters[param]
219         super().__init__(parameters, expr, self.id, self.symbols
)
220
221
222 class SplitLorentzian(SignalLine):
223     def __init__(self, parameters:dict, name:str = None) -> None
:
224         if name is None:
225             self.id = next(id_)
226         else:
227             self.id = name
228         expr = splitLorentzian_expr.replace('area', 'area'+self.id).
replace('center','center'+self.id).replace('fwhm1', 'fwhm1'+self.id).replace('fwhm2', 'fwhm2'+self.id)
229         self.symbols = {}
230         for sym, param in zip(sympy.symbols('center fwhm1 fwhm2
area').replace('area', 'area'+self.id).replace('center','center'+self.id).replace('fwhm1', 'fwhm1'+self.id).replace('fwhm2', 'fwhm2'+self.id)), parameters.keys()):
231             self.symbols[sym] = parameters[param]

```

```

232         super().__init__(parameters, expr, self.id, self.symbols
233     )
234
235 class SplitGaussian(SignalLine):
236     def __init__(self, parameters:dict, name:str = None) -> None
237     :
238         if name is None:
239             self.id = next(id_)
240         else:
241             self.id = name
242             expr = splitGauss_expr.replace('area', 'area'+self.id).
243             replace('center','center'+self.id).replace('fwhm1', 'fwhm1'+
244             self.id).replace('fwhm2', 'fwhm2'+self.id)
245             self.symbols = {}
246             for sym, param in zip(sympy.symbols('center fwhm1 fwhm2
247             area').replace('area', 'area'+self.id).replace('center',
248             'center'+self.id).replace('fwhm1', 'fwhm1'+self.id).replace(
249             'fwhm2', 'fwhm2'+self.id)), parameters.keys()):
250                 self.symbols[sym] = parameters[param]
251             super().__init__(parameters, expr, self.id, self.symbols
252     )
253
254
255 class Eps_1(SignalLine):
256     """I dont think it is actually center, but just kept the
257     name as it is the plasmon energy"""
258     def __init__(self, parameters:dict, name:str = None) ->
259     None:
260         if name is None:
261             self.id = next(id_)
262         else:
263             self.id = name
264             expr = eps_1_expr.replace('area', 'area'+self.id).
265             replace('center','center'+self.id).replace('fwhm', 'fwhm'+
266             self.id)
267             self.symbols = {}
268             for sym, param in zip(sympy.symbols('center fwhm area').
269             replace('area', 'area'+self.id).replace('center','center'+
270             self.id).replace('fwhm', 'fwhm'+self.id)), parameters.keys()):
271                 self.symbols[sym] = parameters[param]
272             super().__init__(parameters, expr, self.id, self.symbols
273     )
274
275 class Eps_2(SignalLine):
276     """I dont think it is actually center, but just kept the
277     name as it is the plasmon energy"""

```

```

261     def __init__(self, parameters:dict, name:str = None) ->
262         None:
263             if name is None:
264                 self.id = next(id_)
265             else:
266                 self.id = name
267                 expr = eps_2_expr.replace('area', 'area'+self.id).
268                     replace('center','center'+self.id).replace('fwhm', 'fwhm'+
269                         self.id)
270                 self.symbols = {}
271                 for sym, param in zip(sympy.symbols('center fwhm area'.
272                     replace('area', 'area'+self.id).replace('center','center'+
273                         self.id).replace('fwhm', 'fwhm'+self.id)), parameters.keys()):
274                     :
275                         self.symbols[sym] = parameters[param]
276                     super().__init__(parameters, expr, self.id, self.symbols
277 )
278
279
280
281
282
283
284
285
286 class Experimental(SignalLine):
287     """I dont think it is actually center, but just kept the
288     name as it is the plasmon energy"""
289     def __init__(self, parameters:dict, name:str = None) ->
290         None:
291             if name is None:
292                 self.id = next(id_)
293             else:
294                 self.id = name
295                 expr = experimental_expr.replace('area', 'area'+self.id).
296                     replace('center','center'+self.id).replace('fwhm', 'fwhm'+
297                         self.id)
298                 self.symbols = {}
299                 for sym, param in zip(sympy.symbols('center fwhm area'.
300                     replace('area', 'area'+self.id).replace('center','center'+
301                         self.id).replace('fwhm', 'fwhm'+self.id)), parameters.keys()):
302                     :
303                         self.symbols[sym] = parameters[param]
304                     super().__init__(parameters, expr, self.id, self.symbols
305 )
306
307
308
309
310
311
312
313
314
315
316 class LogNormal(SignalLine):
317     def __init__(self, parameters:dict, name:str = None) -> None
318         :
319             if name is None:
320                 self.id = next(id_)
321             else:
322                 self.id = name

```

```

292         expr = lognormal_expr.replace('area', 'area'+self.id).
293         replace('center','center'+self.id).replace('sigma', 'sigma'+
294         self.id)
295         self.symbols = {}
296         for sym, param in zip(sympy.symbols('center sigma area').
297         replace('area', 'area'+self.id).replace('center','center'+
298         self.id).replace('sigma', 'sigma'+self.id)), parameters.keys(
299         )):
300             self.symbols[sym] = parameters[param]
301         super().__init__(parameters, expr, self.id, self.symbols
302     )
303
304
305
306
307
308 class PowerLaw(SignalLine):
309
310     def __init__(self, parameters: dict, name=None) -> None:
311         if name is None:
312             self.id = next(id_)
313         else:
314             self.id = name
315         self.symbols = {}
316
317         for sym, param in zip(sympy.symbols('center k area').
318         replace('area', 'area'+self.id).replace('center','center'+
319         self.id).replace('k', 'k'+self.id)), parameters.keys()):
320             self.symbols[sym] = parameters[param]
321             expr = power_expr.replace('area','area'+self.id).replace(
322             ('k','k'+self.id)).replace('center','center'+self.id)
323
324             super().__init__(parameters, expr, self.id, self.symbols
325     )
326
327
328
329
330
331 class Constant(SignalLine):
332     """The parameter it takes in must be a dictionary with a key
333     C. All other parameters are not used.
334
335     TODO: fix the above"""
336     def __init__(self, parameters: dict, name=None) -> None:
337         if name is None:
338             self.id = next(id_)
339         else:
340             self.id = name
341         self.symbols = {}
342         # for sym, param in zip(sympy.symbols('C'.replace('C','C
343         '+self.id)), parameters.keys()):

```

```

326     sym, param = sympy.symbols('C'.replace('C','C'+self.id))
327     , parameters['C']
328     self.symbols[sym] = param
329     super().__init__(parameters, 'C'.replace('C','C'+self.id)
330     ), self.id, self.symbols)
331
332 """Model.py"""
333 import numpy as np
334 from multiprocessing.pool import Pool
335 import logging
336 from tqdm.notebook import trange, tqdm
337 import MyFunc
338 from scipy.optimize import curve_fit
339 from signal_lines import *
340 import scripting
341 from skimage.segmentation import slic
342 from time import time
343 from scipy.stats import skew, kurtosis
344 import pickle
345
346
347 hbar = 6.582119569e-16      #eV s
348 e_0  = 8.8541878128e-12    #F/m
349 e     = 1.6e-19
350
351
352
353
354
355
356 class Model:
357     """A model similar to that of hyperspy."""
358     def __init__(self, data, x:np.array=None) -> None:
359         self.data          = data
360         self.signal_lines = {}
361         self.symbols       = {}
362         self.expr          = ' '
363         self.masks         = None
364         self.N             = None
365         self.x             = x
366
367
368     def append(self, curve: SignalLine)-> None:
369         """Appends a SignalLine to the model"""

```

```

371     if curve.id in self.signal_lines.keys():
372         logging.warning('This is already in the model')
373         self.signal_lines[curve.id] = curve
374
375     for sym in curve.symbols.keys():
376         self.symbols[sym] = curve.symbols[sym]
377         self.expr += '+'*(self.expr!='') + curve.expr
378
379
380     def append_labels(self, masks, N=None):
381         """If N < number of labels the multifit algorithm will
382         only fit for the first N labels. Can also use segment_model
383         to use slic for this purpose"""
384         if N is None:
385             N = masks.max()
386         self.masks = masks
387         self.N      = N
388
389     def segment_model(self, mask=None, n_segements=32,
390                      compactness=0.015, enforce=False, **kwargs):
391         """The mask will ultimately make sure that only the TRUE
392         pixels are curve fitted."""
393         img = self.data
394         segments = slic(img, n_segments=n_segements,
395                         compactness=compactness, mask=mask, enforce_connectivity=
396                         enforce, **kwargs)
397         self.append_labels(segments)
398
399     def multifit(self, num_workers = 4, use_parallel = False, **kwargs) -> np.array:
400         """Fits the appended curves to the dataset using scipy.
401         optimize.curve_fit. Default is not parallelizing"""
402         if self.masks is None:
403             self.masks = np.ones(shape=(self.data.shape[0],self.
404 data.shape[1]))
405             self.N      = 1
406             num_params = len(self.symbols.keys())
407             for i in self.symbols.keys():
408                 if self.symbols[i].bmin == -1:
409                     num_params -=1
410             results = np.zeros(shape=(self.data.shape[0], self.data.
411 .shape[1], num_params)) #2 for skew and kurtosis
412             pcov      = np.zeros(shape=(self.data.shape[0], self.data.
413 .shape[1], num_params, num_params)) # , dtype=object)
414             iterable = [[self, results, self.masks==i] for i in range
415 (1,self.N+1)]
416             t = []

```

```

406     if not use_parallel:
407
408         for i in range(1,self.N+1):
409             tic = time()
410
411             fit(self,results, pcov, self.masks==i, maxfev
412             =2000, **kwargs)
413                 toc = time()
414                 t.append((toc-tic)/60)
415                 # print(f'Segment {i} took {(toc-tic)/60} min')
416             return results, pcov
417
418     res = []
419     # the below works, but not with a tqdm
420     with Pool(num_workers) as pool:
421         for result in tqdm(pool.starmap(fit, iterable,
422         chunkszie=1)):
423             res.append(result)
424             pass
425     res = sum([res[i][:,:] for i in range(len(res))])
426             #returns only one array
427     return res
428
429 def fit(self, result, pcov ,mask = None, maxfev:int = 2000, **
430     kwargs):
431     if 'weighted' not in kwargs.keys():
432         weighted=False
433     else:
434         weighted=kwargs['weighted']
435     if mask is None:
436         mask = np.ones(shape= (self.data.shape[0], self.data.
437         shape[1]))
438         self.N = 2
439     param_values = []
440     param_min = []
441     param_max = []
442     for i in self.symbols.keys():
443         if self.symbols[i].bmin == -1:
444             continue
445         param_values.append(self.symbols[i].value)
446         param_min.append(self.symbols[i].bmin)
447         param_max.append(self.symbols[i].bmax)
448     for i in range(self.data.shape[0]):
449         for j in range(self.data.shape[1]):
```

```

448     # print(mask[i,j]*1.0, (i,j))
449     if mask[i,j]*1.0 == 0:
450         continue
451     if np.any(np.isnan(self.data[i,j])):
452         mask[i,j] = 0
453         logging.info(f'Pixel {i,j} has been removed, as
it was not a number')
454         print(f'Pixel {i,j} has been removed, as it was
not a number')
455         continue
456     try:
457         result[i,j], _ = curve_fit(MyFunc.model_function
, self.x, self.data[i,j], p0=param_values, bounds=(param_min,
param_max), maxfev=maxfev)
458         if weighted:
459             temp_result = MyFunc.model_function(self.x,*result[i,j])
460             result[i,j], pcov[i,j] = curve_fit(MyFunc.
model_function, self.x, self.data[i,j], p0=temp_result,
bounds=(param_min,param_max), maxfev=maxfev, absolute_sigma=
True, sigma=1/temp_result**2)
461             param_values = result[i,j]      # uses previous
output as input. Improves speed by quite much
462             if np.any((np.array(param_values) == np.array(
param_min)) )or np.any((np.array(param_values) == np.array(
param_max)) ):
463                 logging.warning(f'Boundary value reached for
coordinate {(i,j)}')
464
465     except RuntimeError:
466         logging.warning('Couldnt find an appropriate fit
. Try to imporve initialization of increase maxfev')
467         continue
468         return result
469
470
471     return result, pcov
472
473 class Result:
474     def __init__(self, model:Model, result:np.array) -> None:
475         self.model = model
476         self.data = result
477     def print_averages(self, data = None, label = None, mask =
None):
478         if data is None:
479             data = self.data.copy()
480         if label is not None:

```

```

481         data = self.data.copy()
482         data[self.model.masks != label] = np.nan
483         i = 0
484         for sym in self.model.symbols.keys():
485             if self.model.symbols[sym].bmin == -1:
486                 print(f'{sym}: {self.model.symbols[sym].value}')
487             else:
488                 print(f'{sym}: {np.nanmean(data[:, :, i])} {np.
489 nanstd(data[:, :, i])}')
490                 i+=1
491
492     def save_model(self, fname:str, mask, skewness=None,
493 kurtosis=None, in_correct_folder = True):
494         """Skewness and kurtosis is appended after everything
495 else. Pass it as a tuple of (mean, std)"""
496         if in_correct_folder:
497             fname = f'Results\{fname}'
498             n_array = self.data
499             # print(n_array.shape)
500             results = {}
501             i = 0
502             for sym in self.model.symbols.keys():
503                 if self.model.symbols[sym].bmin == -1:
504                     # print(f'{sym}: is fixed')
505                     continue
506                     # print(np.nanmean(n_array[:, :, i][mask]))
507                     results[sym.name] = (np.nanmean(n_array[:, :, i][mask
508 ]), np.nanstd(n_array[:, :, i][mask]))
509                     i+=1
510                     if skewness is not None:
511                         results['skewness'] = skewness
512                     if kurtosis is not None:
513                         results['kurtosis'] = kurtosis
514                     f = open(f'{fname}.pkl', 'wb')
515                     pickle.dump(results,f)
516                     f.close()
517
518     def get_pixel_results(self,x,y):
519         r = []
520         i = 0
521         for sym in self.model.symbols.keys():
522             if self.model.symbols[sym].bmin == -1:
523                 r.append(self.model.symbols[sym].value)
524             else:
525                 r.append(self.data[x,y,i])
526                 i+=1
527
528     return r

```

```

524
525     def as_dictionary(self, label=None):
526         d = {}
527         if label is not None:
528             data = self.data.copy()
529             data[self.model.masks != label] = np.nan
530         else:
531             data = self.data.copy()
532
533         i = 0
534         for sym in self.model.symbols.keys():
535             if self.model.symbols[sym].bmin == -1:
536                 d[str(sym)] = self.model.symbols[sym].value
537             else:
538                 d[str(sym)] = data[:, :, i]
539                 i+=1
540
541         # for sym,param in zip(self.model.symbols, range(data.
542         # shape[-1])):
543             #     d[str(sym)] = data[:, :, param]
544             # print(f'{sym}: {data[:, :, param]} {data[:, :, param]}')
545
546     def prod_latex_file(self, filename=r'Tables\test_table.tex',
547     label=None, significant_digits = 2):
548         """This function is not the best"""
549
550         if label is not None:
551             data = self.data.copy()
552             data[self.model.masks != label] = np.nan
553         else:
554             data = self.data.copy()
555             file = open(filename,'w')
556             for sym, param in zip(self.model.symbols.keys(),range(
557             data.shape[-1])):
558                 line = f"""{sym} & {np.round(np.nanmean(data[:, :, param]),significant_digits)} \pm {np.round(np.nanstd(data[:, :, param]), significant_digits)}\\ \\ \n"""
559                 file.write(line)
560             file.close()
561
562
563 """scripting.py"""

```

```

564 """ Warning: running this file in the same cell (.ipynb) or in
565     the same .py file as a curve fit procedure won't work. The
566     file it writes to will only be scripted after the cell/python
567     -file has been completed."""
568
569 def compile_function(model, file_name = 'fitting\\MyFunc.py'):
570     with open(file_name, 'w') as file:
571         args = ''
572         for arg in (model.symbols.keys()):
573             if model.symbols[arg].bmin == -1:
574                 continue
575             args += str(arg) + ','
576         file.write('from numpy import pi, sqrt, log, exp, real\n')
577
578         file.write('from scipy.special import wofz\n')
579         file.write('from scipy.stats import rv_continuous\n')
580         file.write(f'def model_function(x,{args[:-1]}):\n')
581         line = model.expr
582         for sym in model.symbols.keys():
583             if model.symbols[sym].bmin == -1:
584                 line = line.replace(str(sym), str(model.symbols[sym].value))
585             file.write(f'\treturn {line}')
586             file.write('\n')
587             file.write(f'class my_func_gen(rv_continuous):\n\tdef')
588             file.write(f'\t_pdf(self, x,{args[:-1]}):\n\t\t')
589             file.write(f'\t\treturn {line}\n')
590         file.write('''my_func = my_func_gen(name='my_func')''')
591
592     print('Function created')
593
594
595
596
597
598
599 """concatenate.py
600 The function(s) in this file takes in the lowloss and highloss
601 from EELS data in the form of two Hyperspy files.
602 It thereby concatenates them into one single data file. The
603 plots are for now merged at 11eV

```

```

603
604 TODO: Fix the """
605 import numpy as np
606 import hyperspy.api as hs
607 import hyperspy
608 import logging
609 import eels_addon_hyperspy as eah
610
611 merge_energy = 11. # eV, where the high- and
612 lowloss are merged
613 rrtol = 1e-3
614
615
616 def merge(hl, ll, merge_at:float = merge_energy, rtol:float=
617 rrtol, modifier = 'zero_loss', FWHM=0.09, cut_zlp=5.):
618     """hl: Highloss, type hyperspy._signals.eels.EELSSpectrum
619     ll: Lowloss, type hyperspy._signals.eels.EELSSpectrum"""
620     # if modifier not in ['zero_loss', 'gaussian', 'fourier_ratio']:
621     #     NotImplementedError('Only zero_loss and gaussian')
622     assert hl.data.shape[0] == ll.data.shape[0] and hl.data.
623     shape[1] == ll.data.shape[1]
624     lowloss = ll.isig[:merge_at*1.0].deepcopy()
625     highloss = hl.isig[merge_at*1.0:].deepcopy()
626     start = highloss.axes_manager['Energy loss'].offset
627     stop = start + highloss.axes_manager['Energy loss'].scale*
628     highloss.axes_manager['Energy loss'].size
629     x = np.linspace(start, stop, num=highloss.axes_manager['
630     Energy loss'].size)
631     start = lowloss.axes_manager['Energy loss'].offset
632     stop = start + lowloss.axes_manager['Energy loss'].scale*
633     lowloss.axes_manager['Energy loss'].size
634     x_zlp = np.linspace(start, stop, num=lowloss.axes_manager['
635     Energy loss'].size)
636     x_tot = np.unique(np.concatenate([x_zlp, x]))
637     if not np.all(x_tot == np.sort(x_tot)):
638         logging.warning('x-values not sorted')
639     y = np.zeros(shape=(highloss.data.shape[0], highloss.data.
640     shape[1], x_tot.shape[0]))
641     y[:, :, x_tot <= x_zlp[-1]] = lowloss
642             #concatenates the counts
643     y[:, :, x_tot >= x_zlp[-1]] = highloss
644             #concatenates the counts
645     s_ = hyperspy._signals.eels.EELSSpectrum(y)
646     s_.axes_manager[2].name = 'Energy loss'

```

```

639     s_.axes_manager['Energy loss'].offset = x_tot[0]
640     s_.axes_manager['Energy loss'].scale = (x_tot[-1]-x_tot[0])/s_
641     s_.axes_manager['Energy loss'].size
642     s_.axes_manager['Energy loss'].units = 'eV'
643     if not (np.isclose(s_.axes_manager['Energy loss'].scale, hl.
644     axes_manager['Energy loss'].scale, rtol) and np.isclose(s_.
645     axes_manager['Energy loss'].scale, ll.axes_manager['Energy
646     loss'].scale , rtol)):
647         logging.warning('Scale is off!')
648     if modifier=='zero_loss':
649         return x_tot, s_, s_.fourier_log_deconvolution(lowloss,
650     add_zlp=False)
651     if modifier=='gaussian':
652         return x_tot, s_, eah.fourier_log_deconvolution(s_, s_.
653     isig[:cut_zlp*1.0], add_zlp=False, FWHM=FWHM)
654     if modifier=='fourier_ratio':# Not really a modifier, but
655     its convenient for now
656         return x_tot,s_, s_.fourier_ratio_deconvolution(s_.isig
657     [:cut_zlp*1.0], fwhm=FWHM, extrapolate_lowloss=False,
658     extrapolate_coreloss=False)
659     if modifier=='None':
660         return x_tot,s_, s_
661     NotImplementedError('modifier has to be specified. Set it to
662     str(None) if no deconvolution')
663
664
665 """Some handy functions to use in several python scripts"""
666 import numpy as np
667 temperatures = [23,80,120,160,200,240,280,320,360,400,450,500]
668
669 def create_reconvolve_func():
670     raise NotImplementedError
671
672 def fourier_ratio(y:np.array, zlp:np.array,reconvolve_func:np.
673 array=None):
674     """Returns the Fourier signal of the reconvolved signal"""
675     if reconvolve_func is None:
676         reconvolve_func = create_reconvolve_func()
677     return np.fft.fft(reconvolve_func)*np.fft.fft(y)/np.fft.fft(
678     zlp)

```

```

674
675
676 def produce_filename(T,parameters):
677     # print(parameters)
678     components = parameters['components']
679     class_type = str(type(components[0]))[21:-2]
680     # class_type = parameters['class_type']
681     modifier = parameters['modifier']
682     if not parameters['deconvolve']:
683         modifier = 'None'
684     energy_interval = parameters['energy_interval']
685     curve_FWHM = parameters['curve_FWHM']
686     cut_zlp = parameters['cut_zlp']
687     add_constant = parameters['add_constant']
688     align_on_plasmon_peak = parameters['align_on_plasmon_peak']
689     subtract_median = parameters['subtract_median']
690     add_pixels = parameters['add_pixels']
691     # T = parameters['T']
692     if 'add_to_dict_name' in parameters.keys():
693         addition = parameters['add_to_dict_name']
694         return f'{class_type}_{modifier}_energyInterval{energy_interval}_FWHM{curve_FWHM}_cutZLP{cut_zlp}_w{o}*(not add_constant) + f'constant_+alignedPP_*align_on_plasmon_peak+medianSubtracted'*subtract_median+_unweighted'* (not parameters['weighted']) +'_summed_*add_pixels+f'{T}C'+f'_{addition}'

695
696     return f'{class_type}_{modifier}_energyInterval{energy_interval}_FWHM{curve_FWHM}_cutZLP{cut_zlp}_w{o}*(not add_constant) + f'constant_+alignedPP_*align_on_plasmon_peak+medianSubtracted'*subtract_median+_unweighted'* (not parameters['weighted']) +'_summed_*add_pixels+f'{T}C'
697 def get_zlp_thermal():
698     zlp = {}
699     for T in temperatures:
700         avg = []
701         std = []
702         val = []
703         for i in range(1,10):
704             try:
705                 f = f'ZLP/{(T,i)}.npy'
706                 # print(np.any(np.isnan(np.load(f))))
707                 val.append(np.load(f))
708             except:
709                 print(f)
710                 pass

```

```

711         avg = np.mean(np.array(val))
712         std = np.std(np.array(val))
713         zlp[T] = (avg, std)
714     return zlp
715
716 def derivative(y,x, order = 2):
717     """Takes the first derivative. Available for forward
718     differences (order=1) and central differences (order=2)"""
719     dx = (x[-1]-x[0])/x.shape[0]
720     if order==1:
721         return (y[1:] - y[:-1])/dx
722     elif order==2:
723         return np.gradient(y,x)
724     else:
725         raise NotImplementedError('Implementation for higher
726         order finite differences of the first derivative is not
727         implemented')
728
729 def get_zlp_thermal_intensities():
730     zlp = {}
731     for T in temperatures:
732         avg = []
733         std = []
734         val = []
735         for i in range(1,10):
736             try:
737                 f = f'ZLP/{(T,i)}_integral.npy'
738                 # print(np.any(np.isnan(np.load(f))))
739                 val.append(np.load(f))
740             except:
741                 print(f)
742                 pass
743         avg = np.mean(np.array(val))
744         std = np.std(np.array(val))
745         zlp[T] = (avg, std)
746     return zlp
747
748 temp_coeff = 0.00429
749 R_Al       = 2.65e-8
750 def theoretical(T):
751     return R_Al * (1+temp_coeff*(T-(20)))
752
753
754 mask = np.zeros(shape=(51,60)) # Shape of the images

```

```

755 mask[10:,10:] = 1
756 mask = mask > 0.5
757
758
759 def reshape_segments(segments, approved_files, end_shape):
760     """takes the in the good files and the end-shape of the
761     experimental data as well as the segment mask and fits the
762     shape appropriately"""
763     seg = {}
764     for T in segments.keys():
765         s = np.zeros(end_shape)
766         count = 0
767         for i in range(0,11):
768             if T == 23:
769                 if ('RM',i+1) not in approved_files:
770                     print(T,i)
771                     continue
772             else:
773                 if (T,i+1) not in approved_files:
774                     continue
775                 s[:, :, i] = segments[T][:, :, count]
776                 print(count)
777                 count+=1
778         seg[T] = s
779     return seg
780
781 %matplotlib qt          # This is an .ipynnb file used for
782                         # getting the information regarding the resolution
783 import numpy as np
784 import os
785 import hyperspy.api as hs
786 import hyperspy
787 import matplotlib.pyplot as plt
788 import scipy.ndimage as nd
789 from datetime import date
790 from tqdm import trange, tqdm
791 import logging
792 from time import time
793 import sys
794 sys.path.append('c:\\\\Users\\\\krist\\\\OneDrive - NTNU\\\\Semestre\\\\10
795 - 2023 V r\\\\Master\\\\Data handling\\\\fitting')
796 import signal_lines
797 import Params
798 import Model
799 import MyFunc

```

```

797 from Model import hbar, e_0          # note that hbar is taken in
    ev
798 import rcparams
799 rcparams.main()                      # sets the rcParams for
    plotting
800
801 # For automatic loading of the scripts
802 # This is needed in order for the MyFunc.py to be updated when
    needed
803 %load_ext autoreload
804 %autoreload 2
805
806 ##### New cell
807
808 ZLP = signal_lines.SplitGaussian(Params.testing_splitG,'ZLP')
809 ZLP.append_step_func(-1,1)
810 m = Model.Model(np.array(0),np.array(0))
811 m.append(ZLP)
812 m.compile_function()
813
814 ##### New cell
815
816 ##### Loading the data
817 temperatures = ['80','120','160','200','240','280','320','360','
    400','450','500']
818 temperatures = [int(T) for T in temperatures]
819 lowloss      = {}
820 tic = time()
821 result = {}
822 for T in tqdm([23,160,200,240,280,320,360,400,450,500]):
823     result[T] = []
824     for i in range(1,10):
825         if T == 23:
826             temp = hs.load(r'Cambridge_7\S1_RMTemp _ 0.01
s_58nmpixelSTEM SI'+str(i)+'.dm4')
827         else:
828             temp = hs.load(r'Cambridge_7\S1_'+str(T)+'_CTemp _ 
0.01s_58nmpixel'+str(i)+'.dm4')
829         try:
830             l_ = temp[-2]
831             l_.align_zero_loss_peak(print_stats=False,
show_progressbar=False, signal_range=(-1.,1.))
832             start = l_.axes_manager['Energy loss'].offset
833             # stop = start + l_.axes_manager['Energy loss'].scale*l_
             .size

```

```

834     x = start+np.arange(l_.axes_manager['Energy loss'].size)*l_.axes_manager['Energy loss'].scale
835
836     ZLP = signal_lines.SplitGaussian(Params.
837     testing_splitG, 'ZLP')
838     m = Model.Model(l_.data,x)
839     m.append(ZLP)
840     # m.compile_function()
841     # m.append(ZLP)
842     res, pcov = m.multifit(show_progressbar=False,
843     iterpath='serpentine')
844     result[T].append(res)
845     # val = m.components.Gaussian.sigma.map['values']*2*
846     np.sqrt(2*np.log(2))
847     np.save('ZLP/'+'f'{T,i}_SplitG_L', res[:, :, 1])
848     np.save('ZLP/'+'f'{T,i}_SplitG_R', res[:, :, 2])
849     # break
850     except FileNotFoundError:
851         print(i, 'FileNotFoundException')
852     except TypeError:
853         print(i, 'TypeError')
854     except ValueError:
855         print(i, 'ValueError')
856
857
858
859
860 """experimental.py
861
862 The compilation of the signal line may not be optimal
863
864 It is not very efficient as it loads all the data and then curve
865 fits. experimental2.py was therefore used, but this is left
866 as it has more options that were attempted, but ultimately
867 not used"""
868
869 import numpy as np
870 import os
871 import hyperspy.api as hs
872 import hyperspy
873 import matplotlib.pyplot as plt
874 import scipy.ndimage as nd
875 from datetime import date
876 from tqdm import trange
877 # import logging

```

```

874 import logging
875 import pickle
876 from scipy.stats import linregress
877 from time import time
878 from scipy.stats import skew, kurtosis
879 from scipy.stats import goodness_of_fit
880 import concatenate
881 import sys
882 sys.path.append('c:\\\\Users\\\\krist\\\\OneDrive - NTNU\\\\Semestre\\\\10
883 - 2023 V r\\\\Master\\\\Data handling\\\\fitting')
884 import signal_lines
885 import Params
886 import Model
887 import utils
888 import MyFunc
889 from Model import hbar, e_0          # note that hbar is taken in
890           eV
891 import rcparams
892 rcparams.main()                      # sets the rcParams for
893           plotting
894 logging.basicConfig(level=logging.DEBUG)
895 # For automatic loading of the scripts
896 # This is needed in order for the MyFunc.py to be updated when
897           needed
898
899 tic = time()
900
901
902 def run_experimental(T, **kwargs):
903     """
904     ##### Set the parameters here #####
905     T           = str(T)
906     if 'energy_interval' in kwargs.keys():
907         energy_interval = kwargs['energy_interval']
908     else:
909         energy_interval = (10.,16.)
910     if 'deconvolve' in kwargs.keys():
911         deconvolve = kwargs['deconvolve']
912     else:
913         deconvolve = False
914     if 'modifier' in kwargs.keys():
915         modifier = kwargs['modifier']
916     else:
917         modifier      = 'zero_loss'
918         # modifier      = 'gaussian'
919         # modifier      = 'fourier_ratio' # not really a
920         modifier, but it works as a quick fix

```

```

916     if 'curve_FWHM' in kwargs.keys():
917         curve_FWHM = kwargs['curve_FWHM'] # The
918         FWHM for the deconvolution function
919     else:
920         curve_FWHM      = 0.2
921     if 'cut_zlp' in kwargs.keys():
922         cut_zlp = kwargs['cut_zlp']
923     else:
924         cut_zlp      = 2.
925     if 'add_constant' in kwargs.keys():
926         add_constant = kwargs['add_constant']
927     else:
928         add_constant = False
929     if 'subtract_median' in kwargs.keys():
930         subtract_median = kwargs['subtract_median']
931     else:
932         subtract_median = True
933     if 'align_on_plasmon_peak' in kwargs.keys():
934         align_on_plasmon_peak = kwargs['align_on_plasmon_peak']
935     else:
936         align_on_plasmon_peak = False
937     if 'weighted' not in kwargs.keys():
938         kwargs['weighted'] = False
939     if 'add_pixels' in kwargs.keys():
940         add_pixels = kwargs['add_pixels']
941     else:
942         add_pixels=False
943     if 'derivative' in kwargs.keys():
944         derivative = kwargs['derivative']
945     else:
946         derivative = False
947     if 'components' in kwargs.keys():
948         components = kwargs['components']
949     else:
950         if 'signal_line' in kwargs.keys():
951             plasmon = kwargs['signal_line']
952         else:
953             if align_on_plasmon_peak:
954                 para = Params.splitvoigt_PP
955             else:
956                 para = Params.splitvoigt
957             plasmon = signal_lines.SplitVoigt(para,'Plasmon1')
958             components = [plasmon]
959     if add_constant:
960         c      = signal_lines.Constant(Params.const, 'Constant')
961     components.append(c)

```

```

961
962     # class_type = str(type(plasmon))[21:-2]
963     class_type = str(type(components[0]))[21:-2]
964     # dict_name = f'{class_type}_{modifier}_energyInterval{energy_interval}_FWHM{curve_FWHM}_cutZLP{cut_zlp}_w'+'o'* (not
965     # add_constant) + f'constant'+'_alignedPP_'
966     # align_on_plasmon_peak+'medianSubtracted'*subtract_median+
967     # _unweighted'* (not kwargs['weighted'])+f'{T}C'
968     dict_name = utils.produce_filename(T,kwargs)
969     # if 'add_to_dict_name' in kwargs.keys():
970     #     dict_name += kwargs['add_to_dict_name']
971     #####
972     ##### Loading the data
973     #####
974
975     all_data = []
976     hyper = []
977     hyper_l = []
978     ll = []
979     x = []
980     x_tot = []
981     s = []
982
983     for i in range(1,11):
984         # print(str(i))
985         if T == 'RM':
986             temp = hs.load(r'Cambridge_7\S1_'+T+'Temp_0.01
987             s_58nmpixelSTEM SI'+str(i)+'.dm4')
988             elif T=='40':
989                 temp = hs.load(r'Cambridge_7\S1_'+T+'CTemp_0.01
990                 s_58nmpixel.060'+str(i)+'.dm4')
991                 else:
992                     temp = hs.load(r'Cambridge_7\S1_'+T+'CTemp_0.01
993                     s_58nmpixel'+str(i)+'.dm4')
994
995                     try:
996                         l_ = temp[-2]
997                         h_ = temp[-1]
998                         l_.align_zero_loss_peak(also_align=[h_], print_stats
999                         =False, show_progressbar=False, signal_range=(-1.,1.))
1000                         l = l_.isig[:12.]
1001
1002                             # To make sure that the pixels are defined over same
1003                             # energy interval
1004                             h = h_.isig[10.:19.]
1005
1006                             # To make sure that the pixels are defined over same
1007                             # energy interval
1008                             if deconvolve:

```

```

992             x_test, _, s_ = concatenate.merge(hl=h_, ll=l_,
993 modifier=modifier, FWHM=curve_FWHM, cut_zlp=cut_zlp)
994         else:
995             x_test, s_, _ = concatenate.merge(hl=h_, ll=l_,
996 modifier=modifier, FWHM=curve_FWHM, cut_zlp=cut_zlp)
997             if s_.axes_manager['Energy loss'].offset > -0.4:
998                 # print(f'Skipped {i} during loading due to ')
999                 logging.debug(f'Skipped {i} during loading due
1000 to offset value')
1001             continue
1002         x_tot.append(x_test)
1003         if derivative:
1004             s.append(s_.isig[-1.5:20.].isig[:1900].
1005 derivative(axis=-1))
1006             all_data.append(s_.derivative(axis=-1).data)
1007         else:
1008             s.append(s_.isig[-1.5:20.].isig[:1900])
1009             all_data.append(s_.data)
1010             hyper.append( h)                                # if
1011             not sliced, the different scans will not have the same
1012 dimension
1013             hyper_l.append( l)                                # if
1014             not sliced, the different scans will not have the same
1015 dimension
1016             # all_data.append(s_.derivative(axis=-1).data)
1017             # used to be h
1018             ll.append( l.data)                                # if
1019             not sliced, the different scans will not have the same
1020 dimension
1021             # start = h.axes_manager['Energy loss'].offset
1022             # stop = h.axes_manager['Energy loss'].scale*h.
1023 axes_manager['Energy loss'].size + start
1024             # x.append(np.linspace(start,stop, num = h.
1025 axes_manager['Energy loss'].size))
1026             # break
1027         except FileNotFoundError:
1028             print((T,i), 'FileNotFoundException')
1029         except TypeError:
1030             print((T,i), 'TypeError, typically EELS spectrum not
1031 subscriptable')
1032         except ValueError:
1033             print((T,i), 'ValueError')
1034         x      = x_tot # overriding to update for new programming
1035         logging.info(f'Straight after loading (and deconvolution if
1036 performed): x_high = {x[0][-1]}, x_low = {x[0][0]}' )
1037         # print(f'Straight after loading (and deconvolution if
1038 performed): x_high = {x[0][-1]}, x_low = {x[0][0]}' )

```

```

1023     #####
1024
1025     if align_on_plasmon_peak:
1026         temp           = [signal.isig[10:]].inav[20:,20:].deepcopy
1027         () for signal in s]
1028         [temp[i].align_zero_loss_peak(print_stats=False,
1029         show_progressbar=False) for i in trange(len(temp))]
1030         s = temp
1031
1032     #####
1033
1034     ##### Adding the pixels together #####
1035
1036     if add_pixels:
1037         offset_low   = np.array([i.axes_manager['Energy loss'].
1038         offset for i in s]).max()
1039         adjusted     = [i.isig[offset_low:] for i in s]
1040         offset_high = np.array([i.axes_manager['Energy loss']].
1041         offset + i.axes_manager['Energy loss'].size * i.axes_manager
1042         ['Energy loss'].scale for i in adjusted]).min()
1043         adjusted     = [i.isig[:offset_high] for i in adjusted]
1044         size = adjusted[0].axes_manager['Energy loss'].size
1045         scale = adjusted[0].axes_manager['Energy loss'].scale
1046         hyper = adjusted[0]
1047         deleted = 0
1048         for i in range(1,len(adjusted)):
1049             try:
1050                 hyper += adjusted[i]
1051             except ValueError:
1052                 print(i,'ValueError')
1053                 del adjusted[i-deleted]
1054                 deleted+=1
1055                 continue
1056             except IndexError:
1057                 print(i,'IndexError')
1058
1059         if subtract_median:
1060             all_data = np.array([i.data for i in adjusted]).mean
1061             (axis=0)
1062             all_data -= np.median(all_data, axis=-1).reshape((
1063             all_data.shape[0], all_data.shape[1],1))
1064             all_data = [all_data]
1065         else:
1066             all_data = [np.array([i.data for i in adjusted]).
1067             mean(axis=0)]

```

```

1061
1062     x = [np.linspace(start=offset_low, stop=offset_high, num
1063     =size)]
1064     logging.info(f'After summing over the scans: x_high={x
1065     [0][-1]}, x_low={x[0][0]}' )
1066     # print(f'After summing over the scans: x_high={x
1067     [0][-1]}, x_low={x[0][0]}' )
1068
1069 ##### Run the curve-fitting algorithm and get the
1070 results #####
1071 # np.ones(shape=(data.shape[0],data.shape[1]))
1072 print('Starting curve fitting')
1073 for i in range(len(all_data)):
1074     if 'mask' in kwargs.keys():
1075         mask = kwargs['mask']
1076     else:
1077         mask = np.zeros((all_data[i].shape[0],all_data[i].
1078 shape[1]))
1079         mask[10:,10:] = 1
1080         mask = mask > 0.5
1081     m = Model.Model(all_data[i],x[i])
1082     skewness = (np.mean(skew(all_data[i][mask]), bias=
1083     True, axis=-1)), np.std(skew(all_data[i][mask]), bias=
1084     True, axis=-1))
1085     kurt = (np.mean(kurtosis(all_data[i][mask]), bias=
1086     True, axis=-1)), np.std(kurtosis(all_data[i][mask]), bias=
1087     True, axis=-1))
1088     for comp in components:
1089         m.append(comp)
1090     if np.any(np.isnan(all_data[i])):
1091         mask[np.any(np.isnan(all_data[i]), axis=-1)] = False
1092         m.append_labels(masks=mask)
1093     else:
1094         m.segment_model(mask = mask)
1095     res, pcov = m.multifit(num_workers=8, use_parallel=False
1096     , **kwargs)
1097         # TODO: pcov now only for last scan!
1098         r = Model.Result(m,res)
1099         r.save_model(dict_name+str(i) ,mask, kurtosis=kurt,
1100 skewness=skewness)
1101         test = r.data.copy()
1102         j = 0
1103         # print(T,'C')
1104         for sym in r.model.symbols.keys():
1105             if r.model.symbols[sym].bmin == -1:
1106                 # print(f'{sym}: is fixed to {fwhm_gauss[i]}')
1107                 continue

```

```

1096
1097         logging.debug(f'{T}: {sym}: {np.nanmean(test[:, :, j][
1098             mask])}      {np.nanstd(test[:, :, j][mask])}')
1099         j+=1
1100         # np.save(f'numpy_results\{dict_name}\', res)
1101         np.save(os.path.join('numpy_results', dict_name+f'_{i}\'), 
1102             res)
1103         np.save(os.path.join('numpy_results', dict_name+f'_{i}\'
1104             _std'), pcov)
1105         print('Saved')
1106
1107     #####
1108
1109     # ##### Some plotting #####
1110     # filter_zlp = 500*(not align_on_plasmon_peak)
1111     # divide = 2
1112     # i,j = -1,-2
1113     # y = m.data[i:,j:].mean(axis=(0,1))-m.data[i:,j:].mean(axis
1114     # =(0,1))[-1000:].max()/divide
1115     # y = y[filter_zlp: ].copy()
1116     # # print(np.where(y==y.max()))
1117     # if not modifier=='fourier_ratio':
1118     #     max = np.where(y==y.max())[0][0]
1119     #     # print(max[0])
1120     #     x_low_ = np.where(np.abs(y[:max]) == np.abs(y[:max]).min
1121     #     )[0][0]
1122     #     x_high_ = np.where(np.abs(y[max:]) == np.abs(y[max:]).min
1123     #     )[0][0]+max
1124
1125     #     x_low = x[-1][x_low_+filter_zlp]
1126     #     x_high = x[-1][x_high_+filter_zlp]
1127
1128     #     logging.debug(T+C', 'Numerical FWHM:', (x_high-x_low)
1129     #     'eV')
1130     #     np.save(os.path.join('numpy_results', dict_name+'_
1131     numerical'), np.array(x_high-x_low))
1132
1133     # plot_s = slice(1500,-100)
1134
1135     # ##### Plotting the fit #####
1136
1137     # plt.figure()
1138     # i,j = -3,-3
1139     # plt.scatter(x[-1][plot_s],m.data[i,j][plot_s]-m.data[i,j
1140     # ][-1000: ].max() /
1141     divide, s=1, alpha=1, label = '1 Pixel')

```

```

1133     # # plt.scatter()
1134     # plt.scatter(x[-1][plot_s],m.data[i:,j:,plot_s].mean(axis
1135     =(0,1))-m.data[i:,j:].mean(axis=(0,1))[-1000:].max()      /
1136     divide, s=1, alpha=1, label='Multiple pixels')
1137     # plt.scatter(x[-1][plot_s],m.data[mask,plot_s].mean(axis
1138     =(0)) - m.data[mask].mean(axis=(0))[-1000:].max()      /
1139     divide , s=1, alpha=1, label='All pixels')
1140     # plt.plot(x[-1][plot_s] , MyFunc.model_function(x[-1][
1141     plot_s],*res[i,j])- MyFunc.model_function(x[-1][plot_s],*res[
1142     i,j]).max()/divide , label='Other' ,color='red')
1143
1144     # if not modifier=='fourier_ratio':
1145     #     plt.scatter(x_low, 0)
1146     #     plt.scatter(x_high,0)
1147     # plt.plot(x[-1][plot_s], np.zeros(x[-1][plot_s].shape)
1148     #'--')
1149     # # plt.scatter(x[-1],m.data[i:,j:].mean(axis=(0,1))-MyFunc.
1150     model_function(x[-1],*res[i,j]))
1151     # plt.title(f'{i,j}'+'_FWHM='+ str(curve_FWHM) +'_'+T+'C')
1152     # plt.legend()
1153     # plt.savefig(f'Figures\Curve_fit\{dict_name}.png')
1154     # # for sym,val in zip(m.symbols, res[i,j]):
1155     # #     print(sym,val)
1156
1157     # known_params = {}
1158     # count = 0
1159     # for sym,val in zip(m.symbols, res[i,j]):
1160     #     # print(sym,val)
1161     #     known_params[str(sym)] = (val, count)
1162     #     count+=1
1163
1164     # #####
1165     # # ##### Plotting the distribution of a parameter #####
1166     # if class_type=='Lorentzian' or class_type =='Fano':
1167     #     param = 'fwhmPlasmon1'
1168     #     param_index = known_params[param][-1]
1169     #     plt.figure()
1170     #     plt.hist(res[:, :, param_index][mask], bins=50)
1171     #     plt.title(param+'_FWHM='+ str(curve_FWHM) +'_'+T+'C')
1172     #     plt.savefig(f'Figures\{dict_name}.png')
1173     #     # plt.colorbar()
1174
1175     # if class_type=='SplitVoigt':
1176     #     param = 'fwhm2Plasmon1'
1177     #     param_index = known_params[param][-1]
1178     #     plt.figure()

```

```

1171     #     plt.hist(res[:, :, param_index][mask], bins=50, range
1172     #     =(0,1))
1173     #     plt.title(param)
1174     #     plt.savefig(f'Figures\{dict_name}.png')
1175     #     # plt.colorbar()
1176     #     None
1177     #     param = 'fwhm1Plasmon1'
1178     #     param_index = known_params[param][-1]
1179     #     plt.figure()
1180     #     plt.hist(res[:, :, param_index][mask], bins=50, range
1181     #     =(0,1))
1182     #     plt.title(param)
1183     #     plt.savefig(f'Figures\{dict_name}.png')
1184     #     # plt.colorbar()
1185     #     #####
1186     # ##### Plotting the map over the values #####
1187     # plt.figure()
1188     # plt.imshow(res[:, :, param_index])
1189     # plt.title(param)
1190
1191     # toc = time()
1192     # print(f'At {T}C, it took {(toc-tic)/60} min')
1193
1194     # index_a = 1500
1195     # TODO: improve these
1196     # index_b = 1700
1197     # TODO: improve these
1198     # assert modifier == 'zero_loss' or not deconvolve or
1199     # modifier=='fourier_ratio'
1200     # i,j=30,30
1201     # reg = (linregress(MyFunc.model_function(x[-1][index_a:
1202     index_b],*res[i,j]),m.data[i,j,index_a:index_b]))
1203     # reg_ = np.zeros(shape=mask.shape, dtype=object)
1204     # for i in range(res.shape[0]):
1205     #     for j in range(res.shape[1]):
1206     #         if not mask[i,j]:
1207     #             continue
1208     #         reg_[i,j] = (linregress(MyFunc.model_function(x
1209     [-1][index_a:index_b],*res[i,j]),m.data[i,j,index_a:index_b])
1210     )
1211     # np.save(os.path.join('Lin_results',dict_name),reg_)

```

```

1210 """
1211 """ experimental2.py
1212 Some parameters not availavle compared to experimental.py, but
1213     deals with memory better"""
1214 import numpy as np
1215 import os
1216 import hyperspy
1217 import matplotlib.pyplot as plt
1218 import scipy.ndimage as nd
1219 from datetime import date
1220 from tqdm import trange
1221 # import logging
1222 import logging
1223 import pickle
1224 from scipy.stats import linregress
1225 from time import time
1226 from scipy.stats import skew, kurtosis
1227 from scipy.stats import goodness_of_fit
1228 import concatenate
1229 import sys
1230 sys.path.append('c:\\\\Users\\\\krist\\\\OneDrive - NTNU\\\\Semestre\\\\10
1231 - 2023 V r\\\\Master\\\\Data handling\\\\fitting')
1232 import signal_lines
1233 import Params
1234 import Model
1235 import utils
1236 import MyFunc
1237 from Model import hbar, e_0          # note that hbar is taken in
1238           eV
1239 import rcparams
1240 rcparams.main()                      # sets the rcParams for
1241           plotting
1242
1243 tic = time()
1244 prel_mask = utils.mask      # preliminary mask, taking out vacuum
1245           and contamination
1246
1247 def run_experimental(T, **kwargs):
1248     """
1249     ##### Set the parameters here #####
1250     T              = str(T)
1251     if 'energy_interval' in kwargs.keys():
1252         energy_interval = kwargs['energy_interval']
1253     else:

```

```

1252     energy_interval = (10.,16.)
1253     if 'deconvolve' in kwargs.keys():
1254         deconvolve = kwargs['deconvolve']
1255     else:
1256         deconvolve = False
1257     if 'modifier' in kwargs.keys():
1258         modifier = kwargs['modifier']
1259     else:
1260         modifier      = 'zero_loss'
1261         # modifier      = 'gaussian'
1262         # modifier      = 'fourier_ratio' # not really a
1263         modifier, but it works as a quick fix
1264     if 'curve_FWHM' in kwargs.keys():
1265         curve_FWHM = kwargs['curve_FWHM']           # The
1266         FWHM for the reconvolution function
1267     else:
1268         curve_FWHM      = 0.2
1269     if 'cut_zlp' in kwargs.keys():
1270         cut_zlp = kwargs['cut_zlp']
1271     else:
1272         cut_zlp      = 2.
1273     if 'add_constant' in kwargs.keys():
1274         add_constant = kwargs['add_constant']
1275     else:
1276         add_constant      = False
1277     if 'subtract_median' in kwargs.keys():
1278         subtract_median = kwargs['subtract_median']
1279     else:
1280         subtract_median = True
1281     if 'align_on_plasmon_peak' in kwargs.keys():
1282         align_on_plasmon_peak = kwargs['align_on_plasmon_peak']
1283     else:
1284         align_on_plasmon_peak = False
1285     if 'weighted' not in kwargs.keys():
1286         kwargs['weighted'] = False
1287     if 'add_pixels' in kwargs.keys():
1288         add_pixels = kwargs['add_pixels']
1289     else:
1290         add_pixels=False
1291     if 'derivative' in kwargs.keys():
1292         derivative = kwargs['derivative']
1293     else:
1294         derivative = False
1295     if 'components' in kwargs.keys():
1296         components = kwargs['components']
1297     else:
1298         if 'signal_line' in kwargs.keys():

```

```

1297         plasmon = kwargs['signal_line']
1298     else:
1299         if align_on_plasmon_peak:
1300             para = Params.splitvoigt_PP
1301         else:
1302             para = Params.splitvoigt
1303             plasmon = signal_lines.SplitVoigt(para, 'Plasmon1')
1304             components = [plasmon]
1305     if add_constant:
1306         c           = signal_lines.Constant(Params.const, 'Constant'
1307     )
1308     components.append(c)
1309
1310     # class_type = str(type(plasmon))[21:-2]
1311     class_type = str(type(components[0]))[21:-2]
1312     # dict_name = f'{class_type}_{modifier}_energyInterval{energy_interval}_FWHM{curve_FWHM}_cutZLP{cut_zlp}_w'+'o'* (not add_constant) + f'constant'+'_alignedPP_'
1313     align_on_plasmon_peak+'medianSubtracted'*subtract_median+'_unweighted'* (not kwargs['weighted'])+f'{T}C'
1314     dict_name = utils.produce_filename(T,kwargs)
1315     x          = []
1316     x_tot      = []
1317     s          = []
1318     for i in range(1,11):
1319         # print(str(i))
1320         if T == 'RM':
1321             fname = r'Cambridge_7\S1_'+T+'Temp _ 0.01
1322             s_58nmpixelSTEM SI'+str(i)+'.dm4'
1323         elif T=='40':
1324             fname = r'Cambridge_7\S1_'+T+'CTemp _ 0.01
1325             s_58nmpixel.060'+str(i)+'.dm4'
1326         else:
1327             fname = r'Cambridge_7\S1_'+T+'CTemp _ 0.01
1328             s_58nmpixel'+str(i)+'.dm4'
1329         try:
1330             temp = hs.load(fname)
1331             l_ = temp[-2]
1332             h_ = temp[-1]
1333             l_.align_zero_loss_peak(also_align=[h_], print_stats=False, show_progressbar=False, signal_range=(-1.,1.))
1334             # l = l_.isig[:12.]
1335
1336             # To make sure that the pixels are defined over same energy interval
1337             # h = h_.isig[10.:19.]
```

```

# To make sure that the pixels are defined over
same energy interval
1331     if deconvolve:
1332         x_test, _, scan = concatenate.merge(hl=h_, ll=l_,
1333                                     modifier=modifier, FWHM=curve_FWHM, cut_zlp=cut_zlp)
1334     else:
1335         x_test, scan, _ = concatenate.merge(hl=h_, ll=l_,
1336                                     modifier=modifier, FWHM=curve_FWHM, cut_zlp=cut_zlp)
1337         if scan.axes_manager['Energy loss'].offset > -0.4:
1338             print(f'Skipped {(T,i)} during loading due to
1339 offset value')
1340             logging.debug(f'Skipped {i} during loading due
1341 to offset value')
1342             continue
1343         scan = scan.isig[-1.5:20.].isig[:1900]
1344         if derivative:
1345             scan = scan.derivative(axis=-1)
1346
1347         ##### Switching to np arrays
1348         x = scan.axes_manager['Energy loss'].offset +scan.
1349 axes_manager['Energy loss'].scale *np.arange(scan.
1350 axes_manager['Energy loss'].size)
1351         data = scan.data
1352
1353         ##### Loading and
1354         deconvolving and derivative done, starting curve fit
1355         #####
1356         if 'mask' in kwargs.keys():
1357             mask = kwargs['mask']
1358         else:
1359             mask = prel_mask
1360             print(data.shape, x.shape)
1361             m = Model.Model(data, x)
1362             for comp in components:
1363                 m.append(comp)
1364             if np.any(np.isnan(data)):
1365                 mask[np.any(np.isnan(data), axis=-1)] = False
1366                 m.append_labels(masks=mask)
1367             else:
1368                 m.segment_model(mask = mask)
1369             res, pcov = m.multifit(num_workers=8, use_parallel=
1370 False, **kwargs)
1371             r = Model.Result(m,res)
1372             r.save_model(dict_name+str(i) ,mask)
1373             for j,sym in enumerate(r.model.symbols.keys()):
1374                 if r.model.symbols[sym].bmin == -1:
1375                     continue

```

```

1367         logging.debug(f'{T}: {sym}: {np.nanmean(r.data
1368 [:,:,j][mask])} {np.nanstd(r.data[:, :, j][mask])}')
1369         j+=1
1370         np.save(os.path.join('numpy_results', dict_name+f'_{i
1371 }'), res)
1372         np.save(os.path.join('numpy_results', dict_name+f'_{i
1373 }_std'), pcov)
1374         print(f'Saved {(T,i)}')
1375
1376     except FileNotFoundError:
1377         print((T,i), 'FileNotFoundException')
1378     except TypeError:
1379         print((T,i), 'TypeError, typically EELS spectrum not
1380 subscriptable')
1381     except ValueError as e:
1382         print((T,i), 'ValueError:', e)
1383
1384 """plot_class.py
1385 Big thanks to Emil F. Christiansen for providing a notebook this
1386 .py file is heavily based on. Changed slightly to suit my
1387 purposes better"""
1388 import matplotlib.pyplot as plt
1389 from matplotlib.colors import SymLogNorm
1390 import seaborn as sns
1391 import pandas as pd
1392
1393 class SplitLorentzSegmented():
1394     """
1395         A class for storing Lorentz fit parameters of a plasmon peak
1396         after segmentation
1397     """
1398     def __init__(self, data, name=None, temperature=None):
1399         """
1400             Create a VoigtFit instance
1401
1402             Parameters:
1403             -----
1404             data: numpy.ndarray of shape (M, 4)
1405                 Data array with plasmon energy, left fwhm, right
1406                 fwhm and integral (area) in 1st, 2nd, 3rd and 4th position of
1407                 third axis, respectively.
1408             name: str
1409                 The name of the dataset

```

```

1405     temperature: float
1406         The temperature used in the experiment in K
1407     """
1408
1409     self.data = data
1410     self.name = name
1411     self.temperature = temperature
1412     self._parameter_mapping = {'Ep': self.energy, 'fwhm1':
1413     self.fwhm1, 'fwhm2': self.fwhm2, 'integral': self.integral}
1414
1415     def __repr__(self):
1416         return f'{self.__class__.__name__}({self.data!r}, name={
1417         self.name!r}, temperature={self.temperature!r})'
1418
1419     def __str__(self):
1420         return f'{self.__class__.__name__} with shape {self.data
1421         .shape}:\n{self.dataframe.describe()}'
```

1422

```

1423     @property
1424     def energy(self):
1425         return self.data[:, 0]
1426
1427     @property
1428     def fwhm1(self):
1429         return self.data[:, 1]
1430
1431     @property
1432     def fwhm2(self):
1433         return self.data[:, 2]
1434
1435
1436     @property
1437     def integral(self):
1438         return self.data[:, 3]
1439
1440
1441     @property
1442     def dataframe(self):
1443         df = pd.DataFrame(self.data.reshape((self.data.shape[0],
1444             self.data.shape[1])), columns=['Ep', 'fwhm1', 'fwhm2',
1445             'integral'])
1446         df.insert(0, 'T', self.temperature)
1447         return df
1448
1449     def get_parameter(self, parameter):
1450         """
1451             Return a parameter from the fit
1452         """

```

```

1447         return self._parameter_mapping.get(parameter, None)
1448
1449     def plot(self, kind, parameters, *args, **kwargs):
1450         """
1451             Plot the fit results
1452
1453             Parameters:
1454             -----
1455             kind: str
1456                 The kind of plot. Should be either "img" or "hist"
1457             parameters: str or list
1458                 The parameter(s) to plot.
1459             *args: Optional arguments passed to plotting functions
1460             **kwargs: Optional keyword arguments passed to plotting
functions
1461
1462             Returns:
1463             -----
1464             Returns the figure and axes generated by the plotting
functions
1465             """
1466
1467         if parameters is None:
1468             parameters = ['Ep', 'fwhm1', 'fwhm2', 'integral']
1469         else:
1470             if isinstance(parameters, str):
1471                 parameters = [parameters]
1472             else:
1473                 parameters = list(parameters)
1474
1475
1476         if kind == 'img':
1477             fig, axes = plt.subplots(ncols=len(parameters),
nrows=1, sharex=True, sharey=True)
1478             fig.suptitle(f'{self.name}')
1479             if len(parameters) == 1:
1480                 axes = list([axes])
1481             kwargs['norm'] = kwargs.get('norm', SymLogNorm(0.01))
1482         )
1483             for ax, parameter in zip(axes, parameters):
1484                 ax.imshow(self.get_parameter(parameter), *args,
**kwargs)
1485                     ax.set_title(f'{parameter}')
1486             elif kind == 'hist':
1487                 fig, axes = plt.subplots(ncols=len(parameters),
nrows=1)
1488                 fig.suptitle(f'{self.name}')
1489                 if len(parameters) == 1:

```

```

1488         axes = list([axes])
1489         df = self.dataframe
1490         for ax, parameter in zip(axes, parameters):
1491             sns.histplot(df, x=parameter, ax=ax, **kwargs)
1492     else:
1493         return NotImplementedError(f'Kind {kind} is not
1494 implemented')
1495     return fig, axes
1496
1497 class VoigtSegmented():
1498     """
1499         A class for storing Voigt fit parameters of a plasmon peak
1500         after segmentation
1501     """
1502     def __init__(self, data, name=None, temperature=None):
1503         """
1504             Create a VoigtFit instance
1505
1506             Parameters:
1507             -----
1508             data: numpy.ndarray of shape (M, 4)
1509                 Data array with plasmon energy, Gaussian fwhm,
1510                 Lorentzian fwhm and integral (area) in 1st, 2nd, 3rd and 4th
1511                 position of third axis, respectively.
1512             name: str
1513                 The name of the dataset
1514             temperature: float
1515                 The temperature used in the experiment in K
1516         """
1517         self.data = data
1518         self.name = name
1519         self.temperature = temperature
1520         self._parameter_mapping = {'Ep': self.energy, 'Gfwhm':
1521             self.Gfwhm, 'Lfwhm': self.Lfwhm, 'integral': self.integral}
1522
1523     def __repr__(self):
1524         return f'{self.__class__.__name__}({self.data!r}, name={
1525             self.name!r}, temperature={self.temperature!r})'
1526
1527     def __str__(self):
1528         return f'{self.__class__.__name__} with shape {self.data
1529             .shape}:\n{self.dataframe.describe()}'
```

```

1528     @property
1529     def Gfwhm(self):
1530         return self.data[:, 1]
1531
1532     @property
1533     def Lfwhm(self):
1534         return self.data[:, 2]
1535
1536     @property
1537     def integral(self):
1538         return self.data[:, 3]
1539
1540
1541
1542     @property
1543     def dataframe(self):
1544         df = pd.DataFrame(self.data.reshape((self.data.shape[0],
1545                                         self.data.shape[1])), columns=['Ep', 'Gfwhm', 'Lfwhm', 'integral'])
1546         df.insert(0, 'T', self.temperature)
1547         return df
1548
1549     def get_parameter(self, parameter):
1550         """
1551             Return a parameter from the fit
1552         """
1553         return self._parameter_mapping.get(parameter, None)
1554
1555     def plot(self, kind, parameters, *args, **kwargs):
1556         """
1557             Plot the fit results
1558
1559             Parameters:
1560             -----
1561             kind: str
1562                 The kind of plot. Should be either "img" or "hist"
1563             parameters: str or list
1564                 The parameter(s) to plot.
1565             *args: Optional arguments passed to plotting functions
1566             **kwargs: Optional keyword arguments passed to plotting
1567             functions
1568
1569             Returns:
1570             -----
1571                 Returns the figure and axes generated by the plotting
1572                 functions
1573         """

```

```

1571     if parameters is None:
1572         parameters = ['Ep', 'Gfwhm', 'Lfwhm', 'integral']
1573     else:
1574         if isinstance(parameters, str):
1575             parameters = [parameters]
1576         else:
1577             parameters = list(parameters)
1578
1579
1580     if kind == 'img':
1581         fig, axes = plt.subplots(ncols=len(parameters),
1582                                 nrows=1, sharex=True, sharey=True)
1583         fig.suptitle(f'{self.name}')
1584         if len(parameters) == 1:
1585             axes = list([axes])
1586         kwargs['norm'] = kwargs.get('norm', SymLogNorm(0.01))
1587     )
1588     for ax, parameter in zip(axes, parameters):
1589         ax.imshow(self.get_parameter(parameter), *args,
1590                  **kwargs)
1591         ax.set_title(f'{parameter}')
1592     elif kind == 'hist':
1593         fig, axes = plt.subplots(ncols=len(parameters),
1594                                 nrows=1)
1595         fig.suptitle(f'{self.name}')
1596         if len(parameters) == 1:
1597             axes = list([axes])
1598         df = self.dataframe
1599         for ax, parameter in zip(axes, parameters):
1600             sns.histplot(df, x=parameter, ax=ax, **kwargs)
1601     else:
1602         return NotImplementedError(f'Kind {kind} is not
1603 implemented')
1604     return fig, axes
1605
1606
1607 class SplitVoigtSegmented():
1608     """
1609     A class for storing Lorentz fit parameters of a plasmon peak
1610     after segmentation
1611     """
1612     def __init__(self, data, name=None, temperature=None):
1613         """
1614         Create a VoigtFit instance
1615
1616         Parameters:

```

```

1612     -----
1613     data: numpy.ndarray of shape (M, 4)
1614         Data array with plasmon energy, left fwhm, right
1615         fwhm and integral (area) in 1st, 2nd, 3rd and 4th position of
1616         third axis, respectively.
1617         name: str
1618             The name of the dataset
1619         temperature: float
1620             The temperature used in the experiment in K
1621         """
1622         self.data = data
1623         self.name = name
1624         self.temperature = temperature
1625         self._parameter_mapping = {'Ep': self.energy, 'fwhm1':
1626         self.fwhm1, 'fwhm2': self.fwhm2, 'eta' : self.eta, 'integral':
1627         self.integral}
1628
1629     def __repr__(self):
1630         return f'{self.__class__.__name__}({self.data!r}, name={
1631         self.name!r}, temperature={self.temperature!r})'
1632
1633     def __str__(self):
1634         return f'{self.__class__.__name__} with shape {self.data
1635         .shape}:\n{self.dataframe.describe()}'
```

1636 @property  
1637 def energy(self):  
1638 return self.data[:, 0]

1639 @property  
1640 def fwhm1(self):  
1641 return self.data[:, 1]

1642 @property  
1643 def fwhm2(self):  
1644 return self.data[:, 2]

1645 @property  
1646 def eta(self):  
1647 return self.data[:, 3]

1648 @property  
1649 def integral(self):  
1650 return self.data[:, 4]

```

1653     @property
1654     def dataframe(self):
1655         df = pd.DataFrame(self.data.reshape((self.data.shape[0],
1656                                             self.data.shape[1])), columns=['Ep', 'fwhm1', 'fwhm2', 'eta',
1657                                             'integral'])
1658         df.insert(0, 'T', self.temperature)
1659         return df
1660
1661
1662     def get_parameter(self, parameter):
1663         """
1664             Return a parameter from the fit
1665         """
1666         return self._parameter_mapping.get(parameter, None)
1667
1668
1669     def plot(self, kind, parameters, *args, **kwargs):
1670         """
1671             Plot the fit results
1672
1673             Parameters:
1674             -----
1675             kind: str
1676                 The kind of plot. Should be either "img" or "hist"
1677             parameters: str or list
1678                 The parameter(s) to plot.
1679             *args: Optional arguments passed to plotting functions
1680             **kwargs: Optional keyword arguments passed to plotting
1681             functions
1682
1683             Returns:
1684             -----
1685             Returns the figure and axes generated by the plotting
1686             functions
1687         """
1688         if parameters is None:
1689             parameters = ['Ep', 'fwhm1', 'fwhm2', 'integral']
1690         else:
1691             if isinstance(parameters, str):
1692                 parameters = [parameters]
1693             else:
1694                 parameters = list(parameters)
1695
1696
1697         if kind == 'img':
1698             fig, axes = plt.subplots(ncols=len(parameters),
1699                                     nrows=1, sharex=True, sharey=True)
1700             fig.suptitle(f'{self.name}')
1701             if len(parameters) == 1:

```

```

1695         axes = list([axes])
1696         kwargs['norm'] = kwargs.get('norm', SymLogNorm(0.01)
1697     )
1698     for ax, parameter in zip(axes, parameters):
1699         ax.imshow(self.get_parameter(parameter), *args,
1700                  **kwargs)
1701         ax.set_title(f'{parameter}')
1702     elif kind == 'hist':
1703         fig, axes = plt.subplots(ncols=len(parameters),
1704                                 nrows=1)
1705         fig.suptitle(f'{self.name}')
1706         if len(parameters) == 1:
1707             axes = list([axes])
1708         df = self.dataframe
1709         for ax, parameter in zip(axes, parameters):
1710             sns.histplot(df, x=parameter, ax=ax, **kwargs)
1711     else:
1712         return NotImplementedError(f'Kind {kind} is not
1713 implemented')
1714     return fig, axes
1715
1716
1717 class FanoSegmented():
1718     """
1719     A class for storing Fano fit parameters of a plasmon peak
1720     after segmentation
1721     """
1722     def __init__(self, data, name=None, temperature=None):
1723         """
1724             Create a VoigtFit instance
1725
1726             Parameters:
1727             -----
1728             data: numpy.ndarray of shape (M, 4)
1729                 Data array with plasmon energy, fwhm, integral (area
1730                 ), phi and background in 1st, 2nd, 3rd, 4th and 5th position
1731                 of third axis, respectively.
1732             name: str
1733                 The name of the dataset
1734             temperature: float
1735                 The temperature used in the experiment in K
1736         """
1737         self.data = data
1738         self.name = name
1739         self.temperature = temperature
1740         self._parameter_mapping = {'Ep': self.energy, 'fwhm1':
1741             self.fwhm1, 'phi': self.phi, 'integral': self.integral, '

```

```

background' : self.background}

1734
1735     def __repr__(self):
1736         return f'{self.__class__.__name__}({self.data!r}, name={
1737             self.name!r}, temperature={self.temperature!r})'
1738
1739     def __str__(self):
1740         return f'{self.__class__.__name__} with shape {self.data
1741             .shape}:\n{self.dataframe.describe()}'
```

1742

```

1743     @property
1744     def energy(self):
1745         return self.data[:, 0]

1746     @property
1747     def fwhml(self):
1748         return self.data[:, 1]

1749
1750     @property
1751     def integral(self):
1752         return self.data[:, 2]

1753     @property
1754     def phi(self):
1755         return self.data[:, 3]

1756
1757
1758     @property
1759     def background(self):
1760         return self.data[:, 4]

1761
1762
1763
1764
1765     @property
1766     def dataframe(self):
1767         df = pd.DataFrame(self.data.reshape((self.data.shape[0],
1768             self.data.shape[1])), columns=['Ep', 'fwhml', 'integral',
1769             'phi', 'background'])
1770         df.insert(0, 'T', self.temperature)
1771         return df

1772
1773     def get_parameter(self, parameter):
1774         """
1775             Return a parameter from the fit
1776         """
1777
1778         return self._parameter_mapping.get(parameter, None)
```

```

1776
1777     def plot(self, kind, parameters, *args, **kwargs):
1778         """
1779             Plot the fit results
1780
1781             Parameters:
1782             -----
1783             kind: str
1784                 The kind of plot. Should be either "img" or "hist"
1785             parameters: str or list
1786                 The parameter(s) to plot.
1787             *args: Optional arguments passed to plotting functions
1788             **kwargs: Optional keyword arguments passed to plotting
functions
1789
1790             Returns:
1791             -----
1792                 Returns the figure and axes generated by the plotting
functions
1793             """
1794             if parameters is None:
1795                 parameters = ['Ep', 'fwhml', 'phi', 'integral',
'background']
1796             else:
1797                 if isinstance(parameters, str):
1798                     parameters = [parameters]
1799                 else:
2000                     parameters = list(parameters)
2001
2002
2003             if kind == 'img':
2004                 fig, axes = plt.subplots(ncols=len(parameters),
nrows=1, sharex=True, sharey=True)
2005                 fig.suptitle(f'{self.name}')
2006                 if len(parameters) == 1:
2007                     axes = list([axes])
2008                 kwargs['norm'] = kwargs.get('norm', SymLogNorm(0.01))
2009             )
2010                 for ax, parameter in zip(axes, parameters):
2011                     ax.imshow(self.get_parameter(parameter), *args,
**kwargs)
2012                         ax.set_title(f'{parameter}')
2013             elif kind == 'hist':
2014                 fig, axes = plt.subplots(ncols=len(parameters),
nrows=1)
2015                 fig.suptitle(f'{self.name}')
2016                 if len(parameters) == 1:

```

```

1816         axes = list([axes])
1817         df = self.dataframe
1818         for ax, parameter in zip(axes, parameters):
1819             sns.histplot(df, x=parameter, ax=ax, **kwargs)
1820     else:
1821         return NotImplementedError(f'Kind {kind} is not
1822 implemented')
1823     return fig, axes
1824
1825
1826
1827 class VolumePlasmon():
1828     """
1829     A class for storing Voigt fit parameters of a plasmon peak
1830     after segmentation
1831     """
1832     def __init__(self, data, name=None, temperature=None):
1833         """
1834         Create a Volume Plasmon instance
1835
1836         Parameters:
1837         -----
1838         data: numpy.ndarray of shape (M, 4)
1839             Data array with plasmon energy, fwhm, and integral (
1840             area) in 1st, 2nd and 3re position of third axis,
1841             respectively.
1842         name: str
1843             The name of the dataset
1844         temperature: float
1845             The temperature used in the experiment in C
1846         """
1847         self.data = data
1848         self.name = name
1849         self.temperature = temperature
1850         self._parameter_mapping = {'Ep': self.energy, 'fwhm':
1851             self.fwhm, 'integral': self.integral}
1852
1853     def __repr__(self):
1854         return f'{self.__class__.__name__}({self.data!r}, name={
1855             self.name!r}, temperature={self.temperature!r})'
1856
1857     def __str__(self):
1858         return f'{self.__class__.__name__} with shape {self.data
1859             .shape}:\n{self.dataframe.describe()}'
```

@property

```

1856     def energy(self):
1857         return self.data[:, 0]
1858
1859     @property
1860     def fwhm(self):
1861         return self.data[:, 1]
1862
1863     @property
1864     def integral(self):
1865         return self.data[:, 2]
1866
1867
1868
1869     @property
1870     def dataframe(self):
1871         df = pd.DataFrame(self.data.reshape((self.data.shape[0],
1872                                             self.data.shape[1])), columns=['Ep', 'fwhm', 'integral'])
1873         df.insert(0, 'T', self.temperature)
1874         return df
1875
1876     def get_parameter(self, parameter):
1877         """
1878             Return a parameter from the fit
1879         """
1880         return self._parameter_mapping.get(parameter, None)
1881
1882     def plot(self, kind, parameters, *args, **kwargs):
1883         """
1884             Plot the fit results
1885
1886             Parameters:
1887             -----
1888             kind: str
1889                 The kind of plot. Should be either "img" or "hist"
1890             parameters: str or list
1891                 The parameter(s) to plot.
1892             *args: Optional arguments passed to plotting functions
1893             **kwargs: Optional keyword arguments passed to plotting
1894             functions
1895
1896             Returns:
1897             -----
1898             Returns the figure and axes generated by the plotting
1899             functions
1900         """
1901         if parameters is None:
1902             parameters = ['Ep', 'Gfwhm', 'Lfwhm', 'integral']

```

```

1900     else:
1901         if isinstance(parameters, str):
1902             parameters = [parameters]
1903         else:
1904             parameters = list(parameters)
1905
1906
1907         if kind == 'img':
1908             fig, axes = plt.subplots(ncols=len(parameters),
1909                                     nrows=1, sharex=True, sharey=True)
1910             fig.suptitle(f'{self.name}')
1911             if len(parameters) == 1:
1912                 axes = list([axes])
1913             kwargs['norm'] = kwargs.get('norm', SymLogNorm(0.01))
1914         )
1915         for ax, parameter in zip(axes, parameters):
1916             ax.imshow(self.get_parameter(parameter), *args,
1917                       **kwargs)
1918             ax.set_title(f'{parameter}')
1919         elif kind == 'hist':
1920             fig, axes = plt.subplots(ncols=len(parameters),
1921                                     nrows=1)
1922             fig.suptitle(f'{self.name}')
1923             if len(parameters) == 1:
1924                 axes = list([axes])
1925             df = self.dataframe
1926             for ax, parameter in zip(axes, parameters):
1927                 sns.histplot(df, x=parameter, ax=ax, **kwargs)
1928             else:
1929                 return NotImplementedError(f'Kind {kind} is not
1930 implemented')
1931         return fig, axes
1932
1933
1934
1935
1936 """plot_class.py
1937 Big thanks to Emil F. Christiansen for providing a notebook this
1938 .py file is heavily based on. Changed slightly to suit my
1939 purposes better"""
1940
1941 import matplotlib.pyplot as plt
1942 from matplotlib.colors import SymLogNorm
1943 import seaborn as sns
1944 import pandas as pd
1945
1946 class SplitLorentzSegmented():
1947     """
1948         A class for storing Lorentz fit parameters of a plasmon peak
1949         after segmentation

```

```

1939 """
1940     def __init__(self, data, name=None, temperature=None):
1941         """
1942             Create a VoigtFit instance
1943
1944         Parameters:
1945         -----
1946         data: numpy.ndarray of shape (M, 4)
1947             Data array with plasmon energy, left fwhm, right
1948             fwhm and integral (area) in 1st, 2nd, 3rd and 4th position of
1949             third axis, respectively.
1950         name: str
1951             The name of the dataset
1952         temperature: float
1953             The temperature used in the experiment in K
1954
1955         self.data = data
1956         self.name = name
1957         self.temperature = temperature
1958         self._parameter_mapping = {'Ep': self.energy, 'fwhm1':
1959             self.fwhm1, 'fwhm2': self.fwhm2, 'integral': self.integral}
1960
1961     def __repr__(self):
1962         return f'{self.__class__.__name__}({self.data!r}, name={
1963             self.name!r}, temperature={self.temperature!r})'
1964
1965     def __str__(self):
1966         return f'{self.__class__.__name__} with shape {self.data
1967             .shape}:\n{self.dataframe.describe()}'"
1968
1969     @property
1970     def energy(self):
1971         return self.data[:, 0]
1972
1973     @property
1974     def fwhm1(self):
1975         return self.data[:, 1]
1976
1977     @property
1978     def fwhm2(self):
1979         return self.data[:, 2]
1980
1981     @property
1982     def integral(self):
1983         return self.data[:, 3]

```

```

1981
1982     @property
1983     def dataframe(self):
1984         df = pd.DataFrame(self.data.reshape((self.data.shape[0],
1985                                     self.data.shape[1])), columns=['Ep', 'fwhm1', 'fwhm2', '
1986                                     integral'])
1987         df.insert(0, 'T', self.temperature)
1988         return df
1989
1990     def get_parameter(self, parameter):
1991         """
1992             Return a parameter from the fit
1993         """
1994         return self._parameter_mapping.get(parameter, None)
1995
1996     def plot(self, kind, parameters, *args, **kwargs):
1997         """
1998             Plot the fit results
1999
2000             Parameters:
2001             -----
2002             kind: str
2003                 The kind of plot. Should be either "img" or "hist"
2004             parameters: str or list
2005                 The parameter(s) to plot.
2006             *args: Optional arguments passed to plotting functions
2007             **kwargs: Optional keyword arguments passed to plotting
2008             functions
2009
2010             Returns:
2011             -----
2012             Returns the figure and axes generated by the plotting
2013             functions
2014             """
2015             if parameters is None:
2016                 parameters = ['Ep', 'fwhm1', 'fwhm2', 'integral']
2017             else:
2018                 if isinstance(parameters, str):
2019                     parameters = [parameters]
2020                 else:
2021                     parameters = list(parameters)
2022
2023
2024             if kind == 'img':
2025                 fig, axes = plt.subplots(ncols=len(parameters),
2026                                         nrows=1, sharex=True, sharey=True)
2027                 fig.suptitle(f'{self.name}')

```

```

2023         if len(parameters) == 1:
2024             axes = list([axes])
2025             kwargs['norm'] = kwargs.get('norm', SymLogNorm(0.01)
2026         )
2027         for ax, parameter in zip(axes, parameters):
2028             ax.imshow(self.get_parameter(parameter), *args,
2029             **kwargs)
2030             ax.set_title(f'{parameter}')
2031     elif kind == 'hist':
2032         fig, axes = plt.subplots(ncols=len(parameters),
2033         nrows=1)
2034         fig.suptitle(f'{self.name}')
2035         if len(parameters) == 1:
2036             axes = list([axes])
2037             df = self.dataframe
2038             for ax, parameter in zip(axes, parameters):
2039                 sns.histplot(df, x=parameter, ax=ax, **kwargs)
2040         else:
2041             return NotImplementedError(f'Kind {kind} is not
2042             implemented')
2043         return fig, axes
2044
2045 class VoigtSegmented():
2046     """
2047     A class for storing Voigt fit parameters of a plasmon peak
2048     after segmentation
2049     """
2050     def __init__(self, data, name=None, temperature=None):
2051         """
2052             Create a VoigtFit instance
2053
2054             Parameters:
2055             -----
2056             data: numpy.ndarray of shape (M, 4)
2057                 Data array with plasmon energy, Gaussian fwhm,
2058                 Lorentzian fwhm and integral (area) in 1st, 2nd, 3rd and 4th
2059                 position of third axis, respectively.
2060             name: str
2061                 The name of the dataset
2062             temperature: float
2063                 The temperature used in the experiment in K
2064         """
2065         self.data = data
2066         self.name = name
2067         self.temperature = temperature
2068         self._parameter_mapping = {'Ep': self.energy, 'Gfwhm':
2069             self.Gfwhm, 'Lfwhm': self.Lfwhm, 'integral': self.integral}

```

```

2062
2063     def __repr__(self):
2064         return f'{self.__class__.__name__}({self.data!r}, name={
2065             self.name!r}, temperature={self.temperature!r})'
2066
2067     def __str__(self):
2068         return f'{self.__class__.__name__} with shape {self.data
2069             .shape}:\n{self.dataframe.describe()}'
```

2070

```

2071     @property
2072     def energy(self):
2073         return self.data[:, 0]
```

2074

```

2075     @property
2076     def Gfwhm(self):
2077         return self.data[:, 1]
```

2078

```

2079     @property
2080     def Lfwhm(self):
2081         return self.data[:, 2]
```

2082

```

2083     @property
2084     def integral(self):
2085         return self.data[:, 3]
```

2086

```

2087     @property
2088     def dataframe(self):
2089         df = pd.DataFrame(self.data.reshape((self.data.shape[0],
2090             self.data.shape[1])), columns=['Ep', 'Gfwhm', 'Lfwhm',
2091             'integral'])
2092         df.insert(0, 'T', self.temperature)
2093         return df
```

2094

```

2095     def get_parameter(self, parameter):
2096         """
2097             Return a parameter from the fit
2098         """
2099         return self._parameter_mapping.get(parameter, None)
```

2100

```

2101     def plot(self, kind, parameters, *args, **kwargs):
2102         """
2103             Plot the fit results
2104
2105             Parameters:
2106             -----

```

```

2105     kind: str
2106         The kind of plot. Should be either "img" or "hist"
2107     parameters: str or list
2108         The parameter(s) to plot.
2109     *args: Optional arguments passed to plotting functions
2110     **kwargs: Optional keyword arguments passed to plotting
2111         functions
2112
2113     Returns:
2114     -----
2115     Returns the figure and axes generated by the plotting
2116     functions
2117     """
2118     if parameters is None:
2119         parameters = ['Ep', 'Gfwhm', 'Lfwhm', 'integral']
2120     else:
2121         if isinstance(parameters, str):
2122             parameters = [parameters]
2123         else:
2124             parameters = list(parameters)
2125
2126     if kind == 'img':
2127         fig, axes = plt.subplots(ncols=len(parameters),
2128                                 nrows=1, sharex=True, sharey=True)
2129         fig.suptitle(f'{self.name}')
2130         if len(parameters) == 1:
2131             axes = list([axes])
2132         kwargs['norm'] = kwargs.get('norm', SymLogNorm(0.01))
2133     )
2134     for ax, parameter in zip(axes, parameters):
2135         ax.imshow(self.get_parameter(parameter), *args,
2136         **kwargs)
2137         ax.set_title(f'{parameter}')
2138     elif kind == 'hist':
2139         fig, axes = plt.subplots(ncols=len(parameters),
2140                                 nrows=1)
2141         fig.suptitle(f'{self.name}')
2142         if len(parameters) == 1:
2143             axes = list([axes])
2144         df = self.dataframe
2145         for ax, parameter in zip(axes, parameters):
2146             sns.histplot(df, x=parameter, ax=ax, **kwargs)
2147     else:
2148         return NotImplementedError(f'Kind {kind} is not
2149         implemented')
2150     return fig, axes

```

```

2145
2146
2147
2148 class SplitVoigtSegmented():
2149     """
2150         A class for storing Lorentz fit parameters of a plasmon peak
2151         after segmentation
2152     """
2153     def __init__(self, data, name=None, temperature=None):
2154         """
2155             Create a VoigtFit instance
2156
2157             Parameters:
2158             -----
2159             data: numpy.ndarray of shape (M, 4)
2160                 Data array with plasmon energy, left fwhm, right
2161                 fwhm and integral (area) in 1st, 2nd, 3rd and 4th position of
2162                 third axis, respectively.
2163             name: str
2164                 The name of the dataset
2165             temperature: float
2166                 The temperature used in the experiment in K
2167         """
2168         self.data = data
2169         self.name = name
2170         self.temperature = temperature
2171         self._parameter_mapping = {'Ep': self.energy, 'fwhm1':
2172             self.fwhm1, 'fwhm2': self.fwhm2, 'eta' : self.eta, 'integral':
2173             self.integral}
2174
2175     def __repr__(self):
2176         return f'{self.__class__.__name__}({self.data!r}, name={
2177             self.name!r}, temperature={self.temperature!r})'
2178
2179     def __str__(self):
2180         return f'{self.__class__.__name__} with shape {self.data
2181             .shape}:\n{self.dataframe.describe()}'
```

```

2185     def fwhm2(self):
2186         return self.data[:, 2]
2187
2188     @property
2189     def eta(self):
2190         return self.data[:, 3]
2191
2192     @property
2193     def integral(self):
2194         return self.data[:, 4]
2195
2196
2197
2198     @property
2199     def dataframe(self):
2200         df = pd.DataFrame(self.data.reshape((self.data.shape[0],
2201             self.data.shape[1])), columns=['Ep', 'fwhml', 'fwhm2', 'eta',
2202             'integral'])
2203         df.insert(0, 'T', self.temperature)
2204         return df
2205
2206
2207     def get_parameter(self, parameter):
2208         """
2209             Return a parameter from the fit
2210         """
2211         return self._parameter_mapping.get(parameter, None)
2212
2213     def plot(self, kind, parameters, *args, **kwargs):
2214         """
2215             Plot the fit results
2216
2217             Parameters:
2218             -----
2219             kind: str
2220                 The kind of plot. Should be either "img" or "hist"
2221             parameters: str or list
2222                 The parameter(s) to plot.
2223             *args: Optional arguments passed to plotting functions
2224             **kwargs: Optional keyword arguments passed to plotting
2225             functions
2226
2227             Returns:
2228             -----
2229             Returns the figure and axes generated by the plotting
2230             functions
2231             """
2232             if parameters is None:

```

```

2228         parameters = ['Ep', 'fwhml', 'fwhm2', 'integral']
2229     else:
2230         if isinstance(parameters, str):
2231             parameters = [parameters]
2232         else:
2233             parameters = list(parameters)
2234
2235
2236     if kind == 'img':
2237         fig, axes = plt.subplots(ncols=len(parameters),
2238                                 nrows=1, sharex=True, sharey=True)
2239         fig.suptitle(f'{self.name}')
2240         if len(parameters) == 1:
2241             axes = list([axes])
2242         kwargs['norm'] = kwargs.get('norm', SymLogNorm(0.01))
2243     )
2244     for ax, parameter in zip(axes, parameters):
2245         ax.imshow(self.get_parameter(parameter), *args,
2246                    **kwargs)
2247         ax.set_title(f'{parameter}')
2248     elif kind == 'hist':
2249         fig, axes = plt.subplots(ncols=len(parameters),
2250                                 nrows=1)
2251         fig.suptitle(f'{self.name}')
2252         if len(parameters) == 1:
2253             axes = list([axes])
2254         df = self.dataframe
2255         for ax, parameter in zip(axes, parameters):
2256             sns.histplot(df, x=parameter, ax=ax, **kwargs)
2257     else:
2258         return NotImplementedError(f'Kind {kind} is not
2259 implemented')
2260     return fig, axes
2261
2262
2263
2264
2265
2266
2267
2268 class FanoSegmented():
2269     """
2270     A class for storing Fano fit parameters of a plasmon peak
2271     after segmentation
2272     """
2273     def __init__(self, data, name=None, temperature=None):
2274         """
2275         Create a VoigtFit instance
2276
2277         Parameters:
2278         -----
2279         data: numpy.ndarray of shape (M, 4)

```

```

2269             Data array with plasmon energy, fwhm, integral (area
2270             ), phi and background in 1st, 2nd, 3rd, 4th and 5th position
2271             of third axis, respectively.
2272             name: str
2273                 The name of the dataset
2274             temperature: float
2275                 The temperature used in the experiment in K
2276             """
2277             self.data = data
2278             self.name = name
2279             self.temperature = temperature
2280             self._parameter_mapping = {'Ep': self.energy, 'fwhm1':
2281             self.fwhm1, 'phi': self.phi, 'integral': self.integral, 'background':
2282             self.background}
2283
2284             def __repr__(self):
2285                 return f'{self.__class__.__name__}({self.data!r}, name={self.name!r}, temperature={self.temperature!r})'
2286
2287             def __str__(self):
2288                 return f'{self.__class__.__name__} with shape {self.data.shape}:\n{self.dataframe.describe()}'
```

2289

```

2290             @property
2291             def energy(self):
2292                 return self.data[:, 0]
2293
2294
2295             @property
2296             def fwhm1(self):
2297                 return self.data[:, 1]
2298
2299
2300             @property
2301             def integral(self):
2302                 return self.data[:, 2]
2303
2304
2305             @property
2306             def phi(self):
2307                 return self.data[:, 3]
2308
2309             @property
2310             def background(self):
2311                 return self.data[:, 4]
```

```

2310     @property
2311     def dataframe(self):
2312         df = pd.DataFrame(self.data.reshape((self.data.shape[0],
2313                                         self.data.shape[1])), columns=['Ep', 'fwhml', 'integral',
2314                                         'phi', 'background'])
2315         df.insert(0, 'T', self.temperature)
2316         return df
2317
2318     def get_parameter(self, parameter):
2319         """
2320             Return a parameter from the fit
2321         """
2322         return self._parameter_mapping.get(parameter, None)
2323
2324     def plot(self, kind, parameters, *args, **kwargs):
2325         """
2326             Plot the fit results
2327
2328             Parameters:
2329             -----
2330             kind: str
2331                 The kind of plot. Should be either "img" or "hist"
2332             parameters: str or list
2333                 The parameter(s) to plot.
2334             *args: Optional arguments passed to plotting functions
2335             **kwargs: Optional keyword arguments passed to plotting
2336             functions
2337
2338             Returns:
2339             -----
2340             Returns the figure and axes generated by the plotting
2341             functions
2342             """
2343             if parameters is None:
2344                 parameters = ['Ep', 'fwhml', 'phi', 'integral',
2345                               'background']
2346             else:
2347                 if isinstance(parameters, str):
2348                     parameters = [parameters]
2349                 else:
2350                     parameters = list(parameters)
2351
2352
2353             if kind == 'img':
2354                 fig, axes = plt.subplots(ncols=len(parameters),
2355                                         nrows=1, sharex=True, sharey=True)
2356                 fig.suptitle(f'{self.name}')

```

```

2351         if len(parameters) == 1:
2352             axes = list([axes])
2353             kwargs['norm'] = kwargs.get('norm', SymLogNorm(0.01)
2354         )
2355         for ax, parameter in zip(axes, parameters):
2356             ax.imshow(self.get_parameter(parameter), *args,
2357             **kwargs)
2358                 ax.set_title(f'{parameter}')
2359             elif kind == 'hist':
2360                 fig, axes = plt.subplots(ncols=len(parameters),
2361             nrows=1)
2362                 fig.suptitle(f'{self.name}')
2363             if len(parameters) == 1:
2364                 axes = list([axes])
2365                 df = self.dataframe
2366                 for ax, parameter in zip(axes, parameters):
2367                     sns.histplot(df, x=parameter, ax=ax, **kwargs)
2368             else:
2369                 return NotImplementedError(f'Kind {kind} is not
2370             implemented')
2371             return fig, axes
2372
2373
2374
2375
2376
2377
2378
2379
2380
2381
2382
2383
2384
2385
2386
2387
2388
2389
2390 class VolumePlasmon():
2391     """
2392     A class for storing Voigt fit parameters of a plasmon peak
2393     after segmentation
2394     """
2395     def __init__(self, data, name=None, temperature=None):
2396         """
2397             Create a Volume Plasmon instance
2398
2399             Parameters:
2400             -----
2401             data: numpy.ndarray of shape (M, 4)
2402                 Data array with plasmon energy, fwhm, and integral (
2403                 area) in 1st, 2nd and 3rd position of third axis,
2404                 respectively.
2405             name: str
2406                 The name of the dataset
2407             temperature: float
2408                 The temperature used in the experiment in C
2409             """
2410             self.data = data
2411             self.name = name

```

```

2391         self.temperature = temperature
2392         self._parameter_mapping = {'Ep': self.energy, 'fwhm':
2393             self.fwhm, 'integral': self.integral}
2394
2395     def __repr__(self):
2396         return f'{self.__class__.__name__}({self.data!r}, name={
2397             self.name!r}, temperature={self.temperature!r})'
2398
2399     def __str__(self):
2400         return f'{self.__class__.__name__} with shape {self.data
2401 .shape}:\n{self.dataframe.describe()}'
```

2402

```

2403     @property
2404     def energy(self):
2405         return self.data[:, 0]
2406
2407     @property
2408     def fwhm(self):
2409         return self.data[:, 1]
2410
2411     @property
2412     def integral(self):
2413         return self.data[:, 2]
2414
2415     @property
2416     def dataframe(self):
2417         df = pd.DataFrame(self.data.reshape((self.data.shape[0],
2418             self.data.shape[1])), columns=['Ep', 'fwhm', 'integral'])
2419         df.insert(0, 'T', self.temperature)
2420         return df
2421
2422     def get_parameter(self, parameter):
2423         """
2424             Return a parameter from the fit
2425         """
2426         return self._parameter_mapping.get(parameter, None)
2427
2428     def plot(self, kind, parameters, *args, **kwargs):
2429         """
2430             Plot the fit results
2431
2432             Parameters:
2433             -----
2434             kind: str
2435                 The kind of plot. Should be either "img" or "hist"

```

```

2434     parameters: str or list
2435         The parameter(s) to plot.
2436     *args: Optional arguments passed to plotting functions
2437     **kwargs: Optional keyword arguments passed to plotting
functions
2438
2439     Returns:
2440     -----
2441     Returns the figure and axes generated by the plotting
functions
2442     """
2443     if parameters is None:
2444         parameters = ['Ep', 'Gfwhm', 'Lfwhm', 'integral']
2445     else:
2446         if isinstance(parameters, str):
2447             parameters = [parameters]
2448         else:
2449             parameters = list(parameters)
2450
2451
2452     if kind == 'img':
2453         fig, axes = plt.subplots(ncols=len(parameters),
nrows=1, sharex=True, sharey=True)
2454         fig.suptitle(f'{self.name}')
2455         if len(parameters) == 1:
2456             axes = list([axes])
2457         kwargs['norm'] = kwargs.get('norm', SymLogNorm(0.01))
2458     )
2459     for ax, parameter in zip(axes, parameters):
2460         ax.imshow(self.get_parameter(parameter), *args,
**kwargs)
2461         ax.set_title(f'{parameter}')
2462     elif kind == 'hist':
2463         fig, axes = plt.subplots(ncols=len(parameters),
nrows=1)
2464         fig.suptitle(f'{self.name}')
2465         if len(parameters) == 1:
2466             axes = list([axes])
2467             df = self.dataframe
2468             for ax, parameter in zip(axes, parameters):
2469                 sns.histplot(df, x=parameter, ax=ax, **kwargs)
2470             else:
2471                 return NotImplementedError(f'Kind {kind} is not
implemented')
2472     return fig, axes
2473

```

```

2474 ##### A .ipynb file used to compile and run the
2475     algorithm
2476 %matplotlib qt
2477 import numpy as np
2478 import sys
2479 sys.path.append('c:\\\\Users\\\\krist\\\\OneDrive - NTNU\\\\Semestre\\\\10
2480 - 2023 V r\\\\Master\\\\Data handling\\\\fitting')
2481 import methods
2482 # import voigt_methods
2483 import splitLorentz_methods
2484 import pickle
2485 import pandas as pd
2486 from Model import e_0, hbar
2487 import experimental
2488 import plot_class
2489 import summation_and_deconvolve
2490 import analyse_results
2491 import seaborn as sns
2492 import matplotlib.pyplot as plt
2493 from functools import cache
2494 from tqdm import tqdm
2495 import MyFunc
2496 from scipy.stats import linregress
2497 %load_ext autoreload
2498 %autoreload 2
2499
2500 ##### New cell
2501 @cache
2502                                     # Convenient, but sort of
2503                                     # hardcoded. The loading can and perhaps should have been dealt
2504                                     # with in a more proper way, but it worked as a quick work
2505                                     # around
2506 def load_data(temperatures, mask_num, approved_files):
2507                                     # overriding the function so that
2508                                     # it can be rerun multiple times, requires input as tuple
2509                                     # return summation_and_deconvolve.load_data(list(temperatures)
2510                                     , mask_num, list(approved_files))
2511
2512 ##### New cell
2513 # example
2514 import splitLorentz_methods
2515 import lorentzian_methods
2516 import volumeplasmon_methods
2517 import fano_methods
2518 # all_methods = splitLorentz_methods.all_methods[:12]
2519
2520 # all_methods = voigt_methods.all_methods

```

```

2512 # all_methods = volumeplasmon_methods.all_methods
2513 # all_methods = [all_methods[0]]
2514 # all_methods = [splitLorentz_methods.method1000]
2515
2516 # all_methods = [splitLorentz_methods.methodpreSplitLorentz]
2517 # all_methods = [splitLorentz_methods.methodpreFano      ]
2518 # all_methods = [splitLorentz_methods.methodpreLorentz   ]
2519 # all_methods = [splitLorentz_methods.methodpreVolume   ]
2520 # all_methods = [splitLorentz_methods.method1000,
2521     splitLorentz_methods.method1001]
2522
2523 all_methods = splitLorentz_methods.tenseg[-4:]
2524 all_methods[0].compile()
2525 ##### New cell      - where the algorithm is ran after a
2526           pixel-wise scan
2527 delta = 0 # To be removed
2528 narrow_linregress_interval = False
2529 for method in all_methods:
2530     print(f'Starting method {method.name}')
2531     exclude_scans = method.part2['exclude_scans']
2532     _, approved_files = analyse_results.get_numpy_arrays(**method
2533 .part1)
2534
2535     pre_segmentation_data = {}
2536     _[23] = _['RM']
2537     for T in _.keys():
2538         if T=='RM' :
2539             _[23] =_[T]
2540             continue
2541             l1 = np.array([i[0] for i in approved_files])
2542             l2 = np.array([i[1] for i in approved_files])
2543             pre_segmentation_data[int(T)] =_[T]
2544
2545             pre_segmentation_data[int(T)] =_[T]
2546             # arrays=arrays_
2547             segments, mean_value,sigma_value, mask_num =
2548             summation_and_deconvolve.segment(pre_segmentation_data,
2549             pre_plasmon = method.part1['components'][0],
2550             num_params_pre_segment=method.part1['num_params_pre_segment',
2551             ],**method.part2)
2552             temperatures = tuple(method.part2['temperatures'])
2553             experimental_data, offset,scale = load_data(temperatures,
2554             mask_num,tuple(approved_files))
2555             mask_num = min(mask_num,4)
2556             # assert False

```

```

2551 ##### Skipping linear regression for pre-segmentation data
2552 ##########
2553 ##### Pre-segmentation #####
2554 d = experimental_data
2555 val = np.array(pre_segmentation_data[23]).swapaxes(0,1).
2556 swapaxes(1,2)[:, :, :mask_num]
2557
2558 sh = (val.shape[0], val.shape[1], val.shape[2])
2559 x = 10 + np.arange(d[23].shape[-1])*0.010
2560 #Hardcoded
2561
2562 ##### Running for
2563 premethods
2564 if method.name[:3] == 'pre':
2565     for T in tqdm(experimental_data.keys()):
2566         linreg = np.zeros(shape=sh, dtype=object)
2567         val = np.array(pre_segmentation_data[T]).swapaxes
2568 (0,1).swapaxes(1,2)[:, :, :mask_num]
2569         for i in range(sh[0]):
2570             for j in range(sh[1]):
2571                 for k in range(sh[2]):
2572                     if val[i,j,k,0] == 0:
2573                         continue
2574                     # linreg[i,j,k] = linregress(d[T][i,j,k],
2575 [MyFunc.model_function(x, *val[i,j,k])).rvalue**2
2576                     if narrow_linregress_interval:
2577                         raise NotImplementedError('Are you
2578 sure you dont want to change this?')
2579                     delta = 1
2580                     c = np.nanmean(val[:, :, :, 0], where
2581 = (val[:, :, :, 0] > 0))
2582                     I = ((c-delta) < x) & (x < (c+delta)
2583 )
2584                     linreg[i,j,k] = linregress(d[T][i,j,
2585 k][I], MyFunc.model_function(x[I], *val[i,j,k])).rvalue**2
2586                 else:
2587                     curve_fitted_interval = (method.
2588 part2['energy_interval'][0] < x) & (x < method.part2['
2589 energy_interval'][1])
2590                     linreg[i,j,k] = linregress(d[T][i,j,
2591 k][curve_fitted_interval], MyFunc.model_function(x[
2592 curve_fitted_interval], *val[i,j,k])).rvalue**2
2593
2594 np.save(f'linregress_result\{method.name}
2595 _rvalueSquared_{T}'+narrow_linregress_interval+'(narrowed'+
2596 str(delta)), linreg, allow_pickle=True )

```

```

2582         continue
2583
2584
2585 ##### Filtering out empty segments
2586     labels = {}
2587     for T in temperatures:
2588         labels[T] = []
2589         for i in range(segments[T].max()+1):
2590             if (segments[T]==i).sum()!=0:
2591                 labels[T].append(i)
2592             else:
2593                 print(T,i)
2594     x = offset + np.arange(experimental_data[method.part2['temperatures']][0]).shape[-1])*scale
2595     new_fit = summation_and_deconvolve.sum_and_deconvolve(
2596         experimental_data, segments, pre_segmentation_data, mask_num,
2597         labels, mean_value, x, array=pre_segmentation_data, **method.
2598         part2) # Dont really know why I take in the array twice???
2599 ##### Running curve fit #####
2600     popt,pcov = summation_and_deconvolve.fit_curves(new_fit,x,
2601     mean_value,labels,**method.part2) # inclusion of only
2602     energies within energy interval taken care of in fit_curves()
2603
2604 ##### Linear regression analysis #####
2605     d = new_fit
2606     val = np.array(new_fit[23])
2607
2608     # sh = (val.shape[0],val.shape[1],val.shape[2])
2609     x = 10 + np.arange(np.array(d[23]).shape[-1])*0.010
2610             #Hardcoded
2611
2612     for T in tqdm(experimental_data.keys()):
2613         linreg = np.zeros(shape=np.array(new_fit[T]).shape[0],
2614                           dtype=object)
2615         val = np.array(new_fit[T])
2616         for i in range(1,len(new_fit[T])):
2617             # Skipping the
2618             # 0th
2619             en_interval = (method.part2['energy_interval'][0] <
2620             x) & (x < method.part2['energy_interval'][1])
2621             linreg[i] = linregress(val[i][en_interval] ,MyFunc.
2622             model_function(x[en_interval], *popt[T][i])).rvalue**2
2623             # raise RunTimeError('Husk p bytte navn!')
2624             np.save(f'linregress_result\{method.name}_rvalueSquared_
2625             {T}'+(exclude_scans)*'_excluded_scans',linreg, allow_pickle=
2626             True )

```

```

2614 ####### Visualize a plot #####
2615 # plt.figure()
2616 # plt.plot(x,new_fit[200][3])
2617 # plt.plot(x, MyFunc.model_function(x,*popt[200][3]))
2618 ##### Converting the results to a DataFrame to make
2619 plotting easier #####
2620 temp = np.array
2621 ([23,80,120,160,200,240,280,320,360,400,450,500])
2622 class_type = method.part2['class_type']
2623 df = pd.DataFrame()
2624 df_cov = pd.DataFrame()
2625 data={}
2626 for T in temp:
2627     covariance = np.diagonal(np.array(pcov[T]), axis1=1,
2628     axis2=2)
2629     if class_type == 'Fano':
2630         data[T] = plot_class.FanoSegmented(np.array(popt[T])
2631 ,str(T), int(T))
2632         data_cov[T] = plot_class.FanoSegmented(covariance,
2633 str(T), int(T))
2634     if class_type == 'SplitVoigt':
2635         data[T] = plot_class.SplitVoigtSegmented(np.array(
2636 popt[T]),str(T), int(T))
2637         data_cov[T] = plot_class.SplitVoigtSegmented(
2638 covariance,str(T), int(T))
2639     if class_type =='DerivativeLorentzian' or class_type==''
2640 SplitLorentzian':
2641         data[T] = plot_class.SplitLorentzSegmented(np.array(
2642 popt[T]),str(T), int(T))
2643         data_cov[T] = plot_class.SplitLorentzSegmented(
2644 covariance,str(T), int(T))
2645     if class_type =='Voigt':                                #
2646 NOTE: Integral is not really integral as of now
2647         data[T] = plot_class.VoigtSegmented(np.array(popt[T]
2648 ],str(T), int(T)))
2649         data_cov[T] = plot_class.VoigtSegmented(covariance,
2650 str(T), int(T))
2651     if class_type =='VolumePlasmon' or class_type==''
2652 Lorentzian':
2653         data[T] = plot_class.VolumePlasmon(np.array(popt[T])
2654 ,str(T), int(T))
2655         data_cov[T] = plot_class.VolumePlasmon(covariance,
2656 str(T), int(T))

```

```

2645         df      = df.append(data[T].dataframe, ignore_index=True)
2646         df_cov = df_cov.append(data_cov[T].dataframe,
2647                               ignore_index=True)
2648         # unfiltered_df = df.copy()
2649
2650         ##### Filtering end-result #####
2651         param_list = ['Ep', 'fwhm1', 'fwhm2']                                # Split
2652         Lorentz
2653         if class_type=='Fano':
2654             param_list = ['Ep', 'fwhm1']
2655         if class_type == 'Voigt':
2656             param_list = ['Ep', 'Lfwhm']
2657         if class_type == 'VolumePlasmon' or class_type == 'Lorentzian':
2658             param_list = ['Ep', 'fwhm']
2659         for param in param_list:
2660             for T in temp:
2661                 q_low  = df[(df['T'] == int(T))[param].quantile(0.1)
2662                 q_high = df[(df['T'] == int(T))[param].quantile(0.9)
2663                 df_cov = df_cov[((df[param]>= q_low) & (df[param]
2664                 <= q_high)) | (df['T']!=int(T))]           # Has to be
2665                 before the next line
2666                 df      = df[((df[param]>= q_low) & (df[param] <=
2667                 q_high)) | (df['T']!=int(T))]
2668                 # df_cov = df_cov[((df[param]>= q_low) & (df[param]
2669                 ] <= q_high)) | (df['T']!=int(T))]
2670                 # zlp = zlp[((df[param]>= q_low) & (df[param] <=
2671                 q_high)) | (df['T']!=int(T))]
2672                 if (class_type=='SplitLorentzian') or (class_type==
2673                 DerivativeLorentzian'):
2674                     df['FWHM'] = (df['fwhm1'] + df['fwhm2'])/2 # add
2675                     additional column for mean FWHM      # for Lorentzian split
2676                 elif class_type == 'Fano':
2677                     df['FWHM'] = df['fwhm1']                         # for Fano
2678                 elif class_type == 'Voigt':
2679                     df['FWHM'] = df['Lfwhm']
2680                 elif class_type == 'VolumePlasmon' or class_type == 'Lorentzian':
2681                     df['FWHM'] = df['fwhm']
2682                 else:
2683                     raise NotImplementedError
2684         df['Resistivity'] = (e_0/hbar*df['Ep']**2/df['FWHM'])**-1

```

```

2677     df.to_pickle('dataframes_results\\' + method.name +'
2678         excluded_scans'*exclude_scans) # save the method for future reference
2679     df_cov.to_pickle('dataframes_results\\' + method.name +'
2680         excluded_scans'*exclude_scans+'_cov') # save the method for future reference
2681     f = open(f'dataframes_results\\' + method.name + '
2682         excluded_scans'*exclude_scans+'_part1.pkl', 'wb')
2683     pickle.dump(method.part1,f)
2684     f.close()
2685     f = open(f'dataframes_results\\' + method.name + '
2686         excluded_scans'*exclude_scans+'_part2.pkl', 'wb')
2687     pickle.dump(method.part2,f)
2688     f.close()
2689 ##### Plot resistivity #####
2690     # plt.figure()
2691     # sns.scatterplot(df, x='T', y='Resistivity')
2692
2693
2694
2695
2696
2697
2698
2699
2700
2701
2702
2703
2704
2705
2706
2707
2708
2709
2710
2711
2712 """ summation_and_deconvolve.py """
import experimental
from tqdm import tqdm, trange
import numpy as np
import matplotlib.pyplot as plt
import logging
import sys
sys.path.append('c:\\\\Users\\\\krist\\\\OneDrive - NTNU\\\\Semestre\\\\10
- 2023 V r\\\\Master\\\\Data handling\\\\fitting')
import signal_lines
import Params
import Model
# import MyFuncCopy
from skimage.segmentation import slic
import utils

lowest_ev = 10.
temperatures = ['80', '120', '160', '200', '240', '280', '320', '360', '
400', '450', '500']

def segment(arrays, pre_plasmon, num_params_pre_segment, **
kwargs):
    """This function takes in the results and segments them.
    Class type is the class used pre-segmentation.

```

```

2713     num_params_pre_segment is the number of parameters used for
2714     the curve fitting prior to segmentation.
2715     """
2716     class_type = str(type(pre_plasmon))[21:-2]
2717     ##### Getting rid of outliers #####
2718     if 'temperatures' in kwargs.keys():
2719         temp = kwargs['temperatures']
2720     else:
2721         temp = ['80','120','160','200','240','280','320','360','
2722         400','450','500']
2723     if 'filter_outliers_max' in kwargs.keys():
2724         filter_outliers_max = kwargs['filter_outliers_max']
2725     else:
2726         filter_outliers_max = None
2727     if 'filter_outliers_min' in kwargs.keys():
2728         filter_outliers_min = kwargs['filter_outliers_min']
2729     else:
2730         filter_outliers_min = None
2731     if 'additional_constraints' in kwargs.keys():
2732         additional_constraints = kwargs['additional_constraints']
2733     else:
2734         additional_constraints = False
2735     if 'pixels_per_segment' in kwargs.keys():
2736         pixels_per_segment = kwargs['pixels_per_segment']
2737     else:
2738         pixels_per_segment = 20
2739     if 'num_param' in kwargs.keys():
2740         num_param = kwargs['num_param']
2741     else:
2742         num_param = 3                                     # number of
2743                                         parameters to be included. Ended up with 1 for the thesis (i
2744                                         .e. only center) as this provided the best results
2745     if 'compactness' in kwargs.keys():
2746         compactness = kwargs['compactness']
2747     else:
2748         compactness = 1
2749
2750     mask__ = arrays[temp[0]][0][:,:,0] > 10          #
2751     To be removed
2752     mask_num = np.array([len(arrays[T]) for T in temp]).min()
2753                                         #TODO: make the code susceptible for ragged
2754                                         arrays
2755
2756     segments = { }

```

```

2751     n_sigma = kwargs['n_sigma']
2752         # sigma value for including/excluding pixels.
2753             # all pixels included in the
2754 analysis. Should be same for all known values when using
2755 SplitVoigt
2756
2757 mask      = np.zeros(shape=(mask__.shape[0],mask__.shape[1],
2758 mask_num), dtype=bool)      # img_sizeX, img_sizeY, number of
2759 scans
2760 mean_value = {}
2761 sigma_value = {}
2762 _           = {}
2763
2764 for k,T in enumerate(temp):
2765     mask      = np.zeros(shape=(mask__.shape[0],mask__.shape
2766 [1],len(arrays[T])), dtype=bool)      # img_sizeX, img_sizeY,
2767 number of scans
2768     # print(mask_num)
2769     _[T] = np.array(arrays[T][:]).swapaxes(0,1).swapaxes
2770 (1,2)
2771     # for i in range(mask_num):
2772     for i in range(len(arrays[T])):
2773         if (int(T), i) in utils.exclude and kwargs['
2774 exclude_scans']:
2775             mask[:, :, i] = False
2776             print((T,i), 'excluded')
2777             continue
2778         if additional_constraints:
2779             if class_type == 'SplitLorentzian':
2780                 mask_ = (arrays[T][i][:,:,0] > 13) & (
2781 arrays[T][i][:,:,1] > 0.3) & (arrays[T][i][:,:,2] > 0.3) &
2782 (arrays[T][i][:,:,1] < 0.8) & (arrays[T][i][:,:,2] < 1.2)
2783                 elif class_type == 'Fano':
2784                     mask_ = (arrays[T][i][:,:,0] > 13) #& (
2785 arrays[T][i][:,:,1] > 0.3) & (arrays[T][i][:,:,1] < 0.8)
2786                 elif class_type == 'Lorentzian':
2787                     mask_ = (arrays[T][i][:,:,0] > 13) &
2788 (arrays[T][i][:,:,1] > 0.1)
2789                 elif class_type == 'VolumePlasmon':
2790                     mask_ = (arrays[T][i][:,:,0] > 13) &
2791 (arrays[T][i][:,:,1] > 0.1)
2792
2793             else:
2794                 mask_ = arrays[T][i][:,:,0] > 10
2795             mask[:, :, i] = mask_==1
2796             for param in range(num_params_pre_segment):
2797                 arr = arrays[T][i][:,:,param]

```

```

2784             arr_ = arr[mask_]
2785             if (filter_outliers_min is not None) and (
2786                 filter_outliers_min is not None):
2787                 mean_value[(T,i,param)] = np.mean(np.sort(
2788                     arr_)[filter_outliers_min:-filter_outliers_max]) # excluding
2789                     extreme outliers
2790                     sigma_value[(T,i,param)] = np.std( np.sort(
2791                         arr_)[filter_outliers_min:-filter_outliers_max]) # excluding
2792                     extreme outliers
2793                     else:
2794                         mean_value[(T,i,param)] = np.nanmean(arr,
2795                         where=mask_)
2796                         sigma_value[(T,i,param)] = np.nanstd( arr,
2797                         where=mask_)
2798                         mask[:, :, i] = mask[:, :, i] & (np.abs
2799                         ((arr - mean_value[(T,i,param)])) < n_sigma*sigma_value[(T,
2800                             i,param)])
2801                         valid_num_pixels = mask.sum()
2802                         n_segements = int(valid_num_pixels/
2803                             pixels_per_segment)
2804                         segments[T] = slic(_[T][:, :, :, :num_param], n_segments=
2805                             n_segements, compactness=compactness, mask=mask,
2806                             enforce_connectivity=kwarg['enforce_connectivity'],
2807                             min_size_factor=0.9), max_size_factor=1.4)
2808                         print(T,f'C. valid_num_pixels={valid_num_pixels},
2809                             n_segments={n_segements}; Result: n_segments={segments[T].max
2810                             ()}, pixels_per_segment={valid_num_pixels/segments[T].max()})
2811
2812                         return segments, mean_value, sigma_value, mask_num
2813
2814
2815
2816
2817
2818
2819
2820 def load_data(temp, mask_num, approved_files):
2821     """Mask_num is to be removed"""
2822     import hyperspy.api as hs
2823     # mask_num = min(mask_num, 4)
2824     T = temp
2825     last_index = 1000
2826     experimental_data = {}
2827     for T in tqdm(temp):
2828         if (T == 23) or (T=='RM'):
2829             lowloss = []
2830             highloss = []
2831
2832             for i in range(1,12):
2833                 # print(f'Checking {(T,i)}')
2834                 if ('RM',i) not in approved_files:

```

```

2815             logging.warning(f'Skipping{(T,i)}')
2816             continue
2817             lowloss.append(hs.load(r'Cambridge_7\S1_RMTemp_'
2818             _ 0.01s_58nmpixelSTEM SI'+str(i)+'.dm4')[-2]) #  

2819             special note! The scans before were not good enough  

2820             highloss.append(hs.load(r'Cambridge_7\S1_RMTemp_'
2821             _ 0.01s_58nmpixelSTEM SI'+str(i)+'.dm4')[-1]) #  

2822             special note! The scans before were not good enough  

2823             # logging.warning('Note to self: tar kun med  

2824             scannene fom scan 6 for romtemperatur, mulig det blir ragged  

2825             da')
2826             else:
2827                 lowloss = []
2828                 highloss= []
2829                 for i in range(1,12):
2830                     if (str(T),i) not in approved_files:
2831                         logging.warning(f'Skipping{(T,i)}')
2832                         continue
2833                         lowloss.append(hs.load(r'Cambridge_7\S1_'+str(T)
2834                         _ '+CTemp _ 0.01s_58nmpixel'+str(i)+'.dm4')[-2])# for i in
2835                         range(1,mask_num+1])
2836                         highloss.append(hs.load(r'Cambridge_7\S1_'+str(T)
2837                         _ '+CTemp _ 0.01s_58nmpixel'+str(i)+'.dm4')[-1])# for i in
2838                         range(1,mask_num+1])
2839                         for l,h in zip(lowloss,highloss):
2840                             l.align_zero_loss_peak(also_align=[h],
2841                             show_progressbar=False, print_stats=False, signal_range
2842                             =(-1.,1.))
2843                             print(T,(highloss))
2844                             scale = h.axes_manager['Energy loss'].scale
2845                             experimental_data[T] = np.array([h.isig[lowest_ev:]].isig
2846                             [:last_index].data for h in highloss]).swapaxes(0,1).
2847                             swapaxes(1,2)
2848                             # logging.debug(T+ [h.isig[lowest_ev:]].axes_manager['
2849                             Energy loss'].offset for h in highloss])
2850                             print(T,[h.isig[lowest_ev:]].axes_manager['Energy loss'].
2851                             offset for h in highloss])
2852                             print(T,[h.isig[lowest_ev:]].axes_manager['Energy loss'].
2853                             scale for h in highloss])
2854                             return experimental_data, lowest_ev,scale
2855
2856
2857
2858
2859
2860 def sum_and_deconvolve(experimental_data,segments, arrays,
2861 mask_num, labels, mean_value,x, array=None, class_type='
2862 SplitVoigt', **kwargs):
2863     ##### Summing the pixels and deconvolving
2864     #####

```

```

2842     from scipy.optimize import curve_fit
2843     import MyFunc
2844     import utils
2845     import hyperspy.api as hs
2846     if 'temperatures' in kwargs.keys():
2847         temperatures = kwargs['temperatures']
2848     if 'subtract_center' in kwargs.keys():
2849         subtract_center = kwargs['subtract_center']
2850         assert array is not None
2851     else:
2852         subtract_center = False
2853     known_params = {}
2854     offset = x[0]
2855     scale = (x[-1]-x[0])/x.shape[0]
2856     # class_type = str(type(plasmon))[21:-2]
2857     if class_type=='SplitVoigt':
2858         known_params = {'center': 0,
2859                         'fwhm1' : 1,
2860                         'fwhm2' : 2,
2861                         'frac' : 3,
2862                         'area' : 4}
2863     if class_type=='Fano':
2864         known_params = {'center': 0,
2865                         'fwhm' : 1,
2866                         'area' : 2,
2867                         'phi' : 3,
2868                         'background' : 4}
2869     if 'FWHM' in kwargs.keys():
2870         FWHM = kwargs['FWHM'] # eV
2871     else:
2872         FWHM = 0.08
2873     deconv_func = signal_lines.gaussian((np.arange(x.shape[0])-x.shape[0]/2)*scale,0,FWHM,1) #Keep it centered
2874     modifier = kwargs['modify']
2875     new_fit = {}
2876     for T in temperatures:
2877         # print(T)
2878         if T==23:
2879             zlp = hs.load(r'Cambridge_7\S1_RMTemp _ 0.01s_58nmpixelSTEM SI'+str(1)+'.dm4')[2]
2880             else:
2881                 zlp = hs.load(r'Cambridge_7\S1_'+str(T)+'_CTemp _ 0.01s_58nmpixel'+str(1)+'.dm4')[2]
2882             zlp = np.roll(zlp.data[0,0,:x.shape[0]],300) # centering the ZLP prior to deconvolution, but hardcoded, taken from vacuum

```

```

2883     new_fit[T] = []
2884     arr = np.array(arrays[int(T)]).swapaxes(0,1).swapaxes
2885     (1,2)[:, :, :mask_num]
2886     for i in range(0, segments[int(T)].max() + 1):
2887         if subtract_center:
2888             # Not used for end-results
2889             y = []
2890             for j in range(experimental_data[T].shape[0]):
2891                 for k in range(experimental_data[T].shape
2892 [1]):
2893                     for l in range(experimental_data[T].
2894 shape[2]):
2895                         if segments[T][j,k,l] != i:
2896                             continue
2897                         center = np.array(array[T]).swapaxes
2898 (0,1).swapaxes(1,2)[j,k,l,0].copy()           # the center has
2899 index 0. TODO: this is not very convenient programming
2900                         y.append(np.roll(experimental_data[T]
2901 [j,k,l], -int(np.round(center/scale)+500))) ## Adding same
2902 value for all pixels. Not supposed to find the center now
2903 anyway
2904
2905     y = np.array(y).mean(axis=0)
2906
2907     else:
2908         y = experimental_data[T][segments[T]==i].mean(
2909 axis=0).copy()                      # Take the mean over the
2910 pixels
2911     y -= np.median(np.roll(y, 200)[200:2*200])
2912
2913     Hardcoded. Median is taken for 10<E<12
2914     if modifier == 'fourier_ratio':
2915         # deconvolve the summed regions
2916         s = np.fft.ifft(utils.fourier_ratio(y, zlp,
2917 deconv_func))
2918     elif modifier == 'zero_loss':
2919         y += zlp
2920         s = np.fft.ifft(np.fft.fft(zlp))          *np.log(
2921 np.fft.fft(y)/np.fft.fft(zlp)))
2922     elif modifier == 'gaussian':
2923         y += zlp
2924         s = np.fft.ifft(np.fft.fft(deconv_func)*np.log(
2925 np.fft.fft(y)/np.fft.fft(zlp)))
2926     else:
2927         s = y
2928     new_fit[T].append(s.real)
2929
2930 return new_fit

```

```

2913
2914
2915 def fit_curves(new_fit,x_, mean_value,labels, class_type='
2916 SplitVoigt',**kwargs):
2917     ##### running the curve fit again #####
2918     from scipy.optimize import curve_fit
2919     import MyFunc
2920     from tqdm import tqdm
2921     if class_type=='SplitVoigt':
2922         known_params = {'center': 0,
2923                         'fwhm1' : 1,
2924                         'fwhm2' : 2,
2925                         'frac' : 3,
2926                         'area' : 4}
2927         bounds      = ([13.8,0,0,0,0],[16,3,3,1,1e5])
2928         init_guess = [15,0.5,0.5,1,1e3]
2929     if class_type=='Fano':
2930         known_params = {'center': 0,
2931                         'fwhm' : 1,
2932                         'area' : 2,
2933                         'phi' : 3,
2934                         'background' : 4}
2935         bounds      = ([14,0,0,-2*np.pi,0],[16,3,1e6,2*np.pi,1])
2936         init_guess = [15,0.2,2e2,1,0.5]
2937
2938     if (class_type=='SplitLorentzian') or (class_type=='DerivativeLorentzian'):
2939         known_params = {'center': 0,
2940                         'fwhm1' : 1,
2941                         'fwhm2' : 2,
2942                         'area' : 3}
2943         bounds      = ([13.8,0,0,0],[16,3,3,1e5])
2944         init_guess = [15,0.5,0.5,1e3]
2945         if class_type=='DerivativeLorentzian':
2946             init_guess = [15,0.5,0.5,1e1]
2947     if class_type=='Voigt':
2948         known_params = {'center': 0,
2949                         'Gfwhm' : 1,
2950                         'Lwhm' : 2,
2951                         'area' : 3}
2952         bounds      = ([13.8,0,0,0],[16,3,3,1e5])
2953         init_guess = [15,0.1,0.5,1e3]
2954     if (class_type=='VolumePlasmon') or (class_type=='Lorentzian
2955 ') :
2956         known_params = {'center': 0,
2957                         'fwhm' : 1,
2958                         'area' : 2}

```

```

2957     bounds      = ([13.8,0,0],[16,3,1e5])
2958     init_guess = [15,0.5,1e3]
2959
2960     weight = kwargs['weighted']
2961     popt = {}
2962     pcov = {}
2963     for T in tqdm(new_fit.keys()):
2964         popt[T] = []
2965         pcov[T] = []
2966         for i,y_ in enumerate(new_fit[T]):
2967
2968             try:
2969                 if i not in labels[T] or i ==0:
2970                     # dont include the 0th as it is all
2971                     the pixels that were excluded from the scans
2972                     raise RuntimeError
2973
2974                 if kwargs['num_bin'] != -1:
2975                     y_ = np.histogram(x_,weights=y_, bins =
2976                         kwargs['num_bin'])[0]
2977                     x_ = np.histogram(x_,weights=y_, bins =
2978                         kwargs['num_bin'])[1][:-1]
2979
2980                 else:
2981                     y_ = y_.copy()
2982                     x_ = x_.copy()
2983
2984                 en_interval = (kwargs['energy_interval'][0] < x)
2985                 & (x < kwargs['energy_interval'][1])
2986                 x = x[en_interval]
2987                     # Not considering
2988                     what is outside of the pre-defined energy interval
2989                     y = y[en_interval]
2990                     # Not considering
2991                     what is outside of the pre-defined energy interval
2992                     if True:
2993                         # print(kwargs['derivative'])
2994                         if kwargs['derivative'] or class_type==''
2995                             DerivativeLorentzian': # TODO: fix all weighting
2996                             specifications
2997                             if class_type != 'DerivativeLorentzian':
2998                                 raise NotImplementedError
2999                             y = utils.derivative(y,x)
3000                             if weight:
3001                                 logging.warning('Taking the weighted
3002 fit of a derivative.')
3003                             sigma = 1/y**2*weight+1*(not weight)

```

```

2992                     popt_, pcov_ = curve_fit(MyFunc.
2993             model_function, x,y, p0=init_guess, bounds=bounds, sigma=
2994                 sigma, absolute_sigma=True)
2995             if kwargs['temp_adjusted_interval']:
2996                 if class_type != 'SplitLorentzian':
2997                     raise NotImplementedError('Only
2998                         available for SplitLorentzian')
2999                     en_interval = ((popt_[0]-0.5)< x) & (x <
3000                         popt_[0]))
3001                     bounds_      = ([14.4,0,0,0], [popt_[0],1,1,1
3002                         e5])
3003                     init_guess_ = [popt_[0],0.1,0.1,1e3]
3004                     x = x[en_interval]
3005                     y = y[en_interval]
3006                     popt_, pcov_ = curve_fit(MyFunc.
3007             model_function, x,y, p0=init_guess_, bounds=bounds_, sigma=
3008                 sigma[en_interval], absolute_sigma=True)
3009                     popt[T].append(popt_)
3010                     pcov[T].append(pcov_)
3011             except RuntimeError:
3012                 popt[T].append(np.array([np.nan]*len(
3013                     known_params.keys())))
3014                 pcov[T].append(np.ones( (len(known_params.keys()
3015                     ()),len(known_params.keys())) ) *np.nan)
3016                     print(T,i,'Runtimeerror, adding nans')
3017                     continue
3018             return (popt, pcov)

```

