

1 **Experimental and numerical study of a two-body heaving**
2 **wave energy converter with different power take-off models**

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12 **Abstract**

13 Wave energy is one of the most difficult energies to be captured among marine renewables. With
14 the technical progress, wave energy converters (WECs) are being tested in relatively deeper waters,
15 which makes floating concepts almost the only choice. In this paper, a two-body heaving WEC
16 where the wave energy is absorbed through the relative motion between the outer annular and the
17 inner cylindrical buoys is studied. Both experimental and numerical studies are adopted for regular
18 wave conditions. In the physical model test, a hydraulic system is used to achieve constant power
19 take-off (PTO) damping force. Numerical simulations, validated against experimental data, are
20 applied using both the frequency domain and the time domains analyses. Different types of PTOs,

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21 including constant, linear and nonlinear damping forces, are undertaken to evaluate the
22 hydrodynamic and power absorption performance of such device.

23 **Keywords:** two-body heaving wave energy converter; physical model test; numerical
24 simulation; power take-off.

25 **1 Introduction**

26 According to the *Vision for International Deployment of Ocean Energy* by OES (2018), the global
27 potential of ocean renewable energy which could be developed is about 748GW, and the level of its
28 consumption could reduce up to 5.2 billion tons of CO₂ emission by 2050. Utilization of ocean
29 energy resources will contribute to the world's future sustainable power supply. Although wave
30 energy converters (WECs) extracting energy from the ocean surface are facing issues of both safety
31 and efficiency, they still represent a very remarkable share of the overall global supply in the future.
32 Various forms of WECs have been developed worldwide, though none of them has been stood out
33 as a definitive choice (Ji et al., 2020). From the perspective of marine resources, offshore regions
34 have relatively more abundant and stable wave energy than near shore in most sea areas (Castro and
35 Chiang, 2020; Zheng et al., 2014). Thereby, with the technical progress, floating WECs are
36 becoming a hotspot.

37 Generally, a floating WEC system consists of floats which react each other to harness energy from
38 the relative motion in between. Compared to fixed ones, it is more flexible as it less affected by
39 water depth and is not limited by the power take-off (PTO) forms. Besides, it is easier to deploy and
40 maintain. For the multi-body form, it leads to multiple resonant characteristic which broadens the
41 operational sea condition. Taking the advantages, many floating WECs have been investigated.

42 Some devices absorb energy using relative motion of the sealed hull and heavy counterweights
43 inside (Crowley et al., 2018), whose concept amplifies the energy under low frequency and small
44 amplitude waves. One of the representatives is the Wello Penguin, and the full-scale prototype of
45 which has been tested in the European Marine Energy Center (Tethys, 2019). Another WEC under
46 this principle is SEAREV, which utilizes the interaction between the pendular wheel and the hull to
47 produce electricity, and where the centre of gravity of the wheel is off-centred (Cordonnier et al.,
48 2015). The approach of using wave curvature along its propagation direction to extract energy is
49 also considered, and most of which incorporate hydraulic system to convert energy. The raft-type
50 one, such as McCabe Wave Pump, which articulates three rectangular floating pontoons, uses the
51 relative motion in pitch to drive the energy conversion system (Liu et al., 2018). This methodology
52 also underpins the design of Pelamis, whose four rafts move adjacently and drive the hydraulic
53 PTOs to absorb energy (Henderson, 2006).

54 Although the above novel WECs have been well studied, point absorbers utilizing heave motion for
55 wave energy conversion still occupy a certain proportion in wave energy devices. This kind of WEC
56 typically employs floats to react against each other, to generate mechanical energy which is
57 extracted by means of PTO. Powerbuoy, a cost-effective two-body heaving WEC, utilizes relative
58 motion between a float and a spar with a heave plate to drive the push rod and convert wave energy
59 (Van Rij et al., 2017). Inspired by the Powerbuoy, the U.S. Department of Energy funded Reference
60 Model 3 (Neary et al., 2014). Apart of that, many other two-body heaving WECs with different
61 functions are invented, some of which are connected to the grid (Rusu and Onca, 2017) and some
62 are used to power mobile devices (Dai et al., 2017; Shi et al., 2019). Optimization of size and shape
63 is an important aspect of improving energy capture. Son et al. (2014; 2016) designed a shaped

64 bottom named 'Berkeley Wedge' to increase the dynamic response. A streamlined submerged body
65 is adopted in Wavebob to achieve greater relative velocity (Windt et al., 2018). A heave plate is
66 presented in Aegir Dynamo to give the steady reference to the floats (Al-Habaibeh et al., 2010). In
67 addition, Martin et al. (2020) investigated cylinder, sphere and plate type floats, giving sequence of
68 the response performances from the best to the worst. Beatty et al. (2015; 2019) compared different
69 shapes of submerged floats on both Response Amplitude Operator (RAO) and PTO summarily.

70 The dimension of a float, including draught and diameter, also affects its hydrodynamic response
71 (Amiri et al., 2016). Mass ratio between floats has an effect on the energy absorption as well (Liang
72 and Zuo, 2017), which controls the relative motion with PTO together. Aside from the design of the
73 floats, PTO is another vital aspect to WEC. Electromechanical PTO (Castro and Chiang, 2020),
74 linear generator (Tan et al., 2020), hydraulic PTO (Xu et al., 2019) and other different forms of PTO
75 emerge in endlessly. Electromechanical PTO utilizes rack and pinion system combined generator to
76 realize power conversion. The electrical resistance of the generator and the radius of the pinions
77 have an effect on the PTO damping. Some mechanical PTO systems would assemble a variable
78 inertia flywheel to reach the resonance condition. In this case, the rotational inertia and inertia disc
79 radius should be in consideration (Hernández et al., 2017). Linear generators directly link to the
80 wave without any motion transmission, and several key design parameters that will significantly
81 affect the PTO damping coefficient, i.e., the width of the coil and the radius of the central shaft,
82 which in turn affect the energy capture performance. In the hydraulic system, the damping
83 coefficient of hydraulic cylinder and pressure drop determine the PTO damping together. The piston
84 displacement and accumulator define the PTO stiffness (Negandari et al., 2018). In a word, both the
85 magnitudes of PTO damping and PTO stiffness have an influence on the peak value of capture

86 power. What's more, PTO damping also affects the resonance point (Falnes, 1999). Subsequently,
87 different methods are used to find the optimal PTO. Liang and Zuo (2017) found closed-form
88 solutions for both optimal and suboptimal PTO designs. Jin et al. (2019) used a linear frequency
89 domain model to control generator damping and stiffness actively. It is believed that under optimal
90 shape and PTO designs, when wave frequency is within the natural frequencies of the two floats,
91 two-float heaving WEC can produce more power than single-float one. Nevertheless, the passive
92 motion of a WEC under wave excitation could not get an ideal energy efficiency. As the
93 development of control strategy, passive loading, equivalent saturation control and maximum stroke
94 control are applied, providing the possibility of obtaining higher average power (Van den Berg et
95 al., 2011). Multi resonant control and Q-learning algorithm are also introduced to maximize the
96 energy harvesting (Abdelkhalik and Zou, 2019; Anderlini et al., 2018). Generalized analytical phase
97 control conditions (Bubbar and Buckham, 2020) and unlatching control strategy could also increase
98 average energy absorption (Henriques et al., 2012).

99 This paper focuses on a fundamental study of a two-body heaving WEC under regular waves,
100 seeking the maximum power extraction with optimal PTO damping. The dynamic performance is
101 primarily dependent on the PTO force. Most previous studies have used various methods to apply
102 linear PTO damping force, such as mechanical method (Martin et al., 2020) or linear generator (Tan
103 et al., 2020). Although linear PTO damping is simple and common for analysis, constant PTO force
104 is also another major damping form. Thus, in this paper, the physical model test is conducted to
105 explore the hydrodynamic performance of the WEC, and constant PTO damping force is applied to
106 the model via a specially designed hydraulic control. In numerical simulation part, both linear and
107 non-linear PTO damping forces are determined. The results presented in this paper are all model

108 scale values.

109 The paper is organized as follows. In Section 2, a two-body heaving WEC physical model test is
110 established, where the constant PTO damping force is introduced in the model test. Section 3
111 introduces the method of simulation. Section 4 illustrates the comparative analysis of experimental
112 and numerical results, where the simulation model is validated by free decay, RAO and power
113 capture tests. Section 5 discusses the results of the numerical dynamics model, revealing the
114 characteristics of the WEC. Finally, section 6 draws the conclusion of the study, giving
115 methodological comparison and suggestion to the two-body heaving WEC optimization, and
116 expectation of the further work as well.

117 **2 Physical model test**

118 **2.1 Model setup**

119 A 1:9 scale model test of the proposed WEC based on the Froude similarity law is conducted in
120 Shandong Provincial Key Laboratory of Ocean Engineering, as shown in Fig. 1. The wave tank is
121 60 m long, 36 m wide, 1.5 m deep, and features a piston-type wave maker which can generate waves
122 with heights ranging from 0.05 m to 0.25 m, and periods ranging from 0.5 s to 2.5 s in both regular
123 and irregular wave conditions. The tank absorbs wave energy with a sloped porous medium at the
124 end and vertical ones in front of the flanks to minimize the wave reflection. The layout of the wave
125 tank for the physical model test is shown in Fig. 2. The water depth of the physical model test is
126 1.10 m, and the model is placed 30 m from the wave maker and 7 m from one side of the flanks, so
127 that the wall effect can be ignored.

128 Wave gauges are used to measure the fluctuation of the water surface. An NDI Optotrack Certus is

129 used to record the motions of the floats. Before the physical model test, all measurements are
130 calibrated. Fig. 3 shows the positions of the wave gauges and the NDI.

131 The WEC consists of two coaxial floats, which are neutrally buoyant and only have the heave
132 motions. The model set is a non-mooring system. As shown in Fig. 4, the outer float moves along
133 three guide rods (①②③), while the inner one moves along the other two (④⑤). A frame is
134 assembled outside of the model set and fixed to the tank, to make sure the model stays in an upright
135 position. The geometric parameters of the model are given in Table 1 in detail.

136 In the test, a hydraulic system is applied to provide the PTO damping. As shown in Fig. 5, it controls
137 constant PTO force by adjusting hydraulic pressure. In the hydraulic system, relief valve (6) is used
138 to keep steady pressure and regulate pressure. Solenoid operated directional valve (7) is intended to
139 exhaust gas from the hydraulic circuit. Proportional valve (9) is assembled to control the hydraulic
140 pressure. The error between the PTO force provided by the hydraulic system and the desired value
141 is within 10%, and the error of the period is within 1%. It can be considered that the PTO damping
142 force applied by the hydraulic system can meet the test requirement. More detailed analysis of the
143 PTO force is shown in Section 4.3.

144 **2.2 Regular wave condition**

145 Sea state of North China presents the characteristics of small wave height and short wave period.
146 As one of the most representative sea, Zhaitang Island, has been selected as a marine energy test
147 center in North China. The most frequent occurring wave conditions of this island include the wave
148 heights of 0.25–0.75 m and the wave periods of 3.0–4.0 s (Liu et al., 2017). Considering capacity
149 of the wave tank, and inherent performance of the WEC, the model test selects the conditions with

150 a wave height range between 0.075-0.20 m and a period between 1.05-2.30 s.

151 **3 Numerical model**

152 The numerical model is established using the boundary element method software Ansys-Aqwa to
153 obtain the inviscid hydrodynamic coefficients, such as added mass and radiation damping. Panel
154 methods are used to analyze the hydrodynamic behavior in waves (Fig. 6). The numerical model
155 simulates the hydrostatic restoring force, radiation force and wave excitation force. The additional
156 user-defined viscous force, which is obtained by experimental data, can be optionally included in
157 the equation of motion. Furthermore, the PTO damping force is also imposed through user defined
158 process.

159 The governing equation of the two-body heaving WEC in frequency domain is,

$$160 \begin{cases} [-\omega^2(m_1 + A_{11} + A_1) + i\omega(B_{11} + B_{vis1}) + C_1]Z_1 + (-\omega^2A_{12} + i\omega B_{12})Z_2 = F_{e1} \\ [-\omega^2(m_2 + A_{22} + A_2) + i\omega(B_{22} + B_{vis2}) + C_2]Z_2 + (-\omega^2A_{21} + i\omega B_{21})Z_1 = F_{e2} \end{cases} \quad (1)$$

161 where, subscript 1 denotes to the outer float, and 2 represents the inner float, respectively; m_i stands
162 for the mass of a float; A_{ij} is the added mass and A_i is the amended added mass; B_{ij} is the
163 radiation damping coefficient; C is the restoring force coefficient; B_{visi} is the linearized viscous
164 damping coefficient; F_{ei} is the complex amplitude of the exciting force on a float which causes its
165 heave motion; Z_i is the complex amplitude of the heave motion of a float.

166 In frequency domain analysis, PTO damping force is not taken into account, and all the motions are
167 harmonic. The amended added mass and linearized viscous damping coefficients are complemented
168 by free decay test which is discussed in Section 4.1.

169 The set of dynamic equations describing two floats' heave motions in time domain can be expressed
170 as follows,

171
$$\begin{cases} [m_1 + A_{11}(\infty) + A_1]\ddot{z}_1(t) + A_{12}(\infty)\ddot{z}_2(t) + k_{11}(t) * \dot{z}_1(t) + k_{12}(t) * \dot{z}_2(t) + B_{vis1}\dot{z}_1(t) + C_1z_1(t) = f_{e1}(t) + f_{PTO}(t) \\ [m_2 + A_{22}(\infty) + A_2]\ddot{z}_2(t) + A_{21}(\infty)\ddot{z}_1(t) + k_{22}(t) * \dot{z}_2(t) + k_{21}(t) * \dot{z}_1(t) + B_{vis2}\dot{z}_2(t) + C_2z_2(t) = f_{e2}(t) - f_{PTO}(t) \end{cases}$$

172 (2)

173 where, f_{ei} is the exciting force on a float which causes its heave motion; $z(t)$ is the heave motion
 174 of a float, where $\dot{z}(t)$ and $\ddot{z}(t)$ are the velocity and the acceleration, respectively; the symbol $(*)$
 175 denotes the operation of convolution. f_{PTO} is the PTO damping force. Further, $k_{ij}(t)$ is the
 176 radiation-force impulse-response function, which is the inverse Fourier transform of

177
$$K_{ij}(\omega) = i\omega[A_{ij}(\omega) - A_{ij}(\infty)] + B_{ij}(\omega) \quad (3)$$

178 where $A_{ij}(\infty)$ is the added mass when $\omega = \infty$.

179 In time domain analysis, PTO damping force are calculated at each timestep by user-defined routine.
 180 The PTO damping force yields $f_{PTO} = -B \text{sgn}(\dot{z}_1 - \dot{z}_2) |\dot{z}_1 - \dot{z}_2|^{\frac{1}{n}}$, where B is the damping
 181 coefficient, and n is the PTO nonlinear factor. If $n = 1$, the PTO force is linear, and if $n \rightarrow \infty$, the
 182 PTO force is a Coulomb PTO and its value is constant. In fact, when $n > 100$, the nonlinear PTO
 183 force can be very close to a constant PTO force, which is shown in Fig. 7. The results of the time
 184 domain analysis only remain the steady-state response.

185 **4 Model validation and comparison of numerical and experimental**

186 **results**

187 In this section, the numerical model is calibrated by heave decay physical model test and validated
 188 by regular wave physical model tests. The conditions regarding to the PTO damping force are
 189 different in the frequency domain and time domain analyses. For power capture test, the model test
 190 only considers the constant PTO damping force, while the time-domain numerical model simulates
 191 both linear and nonlinear PTO damping force, including the case of constant PTO damping force.

192 The frequency-domain numerical model only considers the cases with no PTO. Table 2 lists the
193 differences in these cases.

194 **4.1 Heave decay with no PTO**

195 Free decay test is taken to check the natural periods, and the hydrodynamic coefficients of the model
196 could be determined at these natural periods. Here, the free decay test is carried out in still water for
197 heave motions. The heave decay tests of the two floats are carried out separately. One of the floats
198 is forced to an initial displacement, along the vertical axis, and then is allowed to return naturally to
199 its equilibrium position. During the process, the decay motion of the float is recorded, while the
200 natural period in heave is measured as the period between peaks of the recorded motion. At the same
201 time, the other float keeps at the equilibrium position. In the data analysis, the total added mass a
202 and the damping coefficient b are determined from the exponential decay curve, which are
203 calculated in the following equations (4) and (5).

$$204 \quad a = \frac{k}{\omega_0^2} - m \quad (4)$$

$$205 \quad b = \frac{2C\xi}{\omega_0} \quad (5)$$

206 where $\omega_0 = \frac{2\pi}{T}$ is the natural frequency of oscillation, C is the restoring force coefficient, m is
207 the mass of a float and ξ is the damping ratio.

208 Based on the decay test data of the physical model and the hydrodynamic coefficients calculated by
209 the software, numerical model is calibrated by modifying the amended added mass and the
210 linearized viscous damping coefficient to match the real motion of the floats considering the viscous
211 effects. Fig. 8 shows the results of a free oscillation test in heave of the two floating bodies. The
212 decay curves from the experimental and numerical results are quite close. As shown in Fig. 8, the

213 natural period of the outer float is approximately 1.29 s whereas the natural period of the inner one
214 is approximately 1.82 s.

215 **4.2 Heave motion RAO with no PTO in regular waves**

216 Regular wave test without the influence of PTO is undertaken to obtain the heave motion RAO and
217 to assess the validity of the numerical simulation modified by free-decay test as discussed in Section
218 4.1. Without the experimental data, the heave motion RAO of this system will be heavily over-
219 predicted. The model tests are performed with two floating bodies coupled. The heave motion RAO,
220 which is defined as the amplitude of the body's heave motion, normalized by the wave amplitude,
221 and also regarded as a function of wave period, is used to evaluate the hydrodynamic performance
222 of the two-body heaving WEC. The experimental RAOs are obtained from regular wave tests
223 without PTO damping force. The ratio between heave motion response and wave signals is obtained
224 over a range of wave periods under regular wave conditions. Noting that, the numerical results
225 conduct two simulation results in frequency and time domain, respectively.

226 Experimental and numerical heave motion RAOs of the two bodies without PTO damping force are
227 plotted in Fig. 9, which are in good agreement. Fig. 9(a) shows that the outer float has a good wave
228 following property in most periods. The responses have a great promotion from 1.05 s to 1.30 s, and
229 remain approximately 1.0 afterwards. Fig. 9(b) reveals that the inner float is more sensitive to the
230 wave period than the outer one, which has an obvious resonance range. That may be due to the
231 streamline shape of the float. At 1.80 s, the resonance condition is approached where the heave
232 displacement amplitude is three times of the incident wave. The relative heave motion RAO is
233 defined as the ratio of the amplitude of the instantaneous relative movement of the outer and inner

234 floats to the wave amplitude, which is shown in Fig. 9(c). It also obtains the peak value at the period
235 of 1.80 s. Since the outer float has no distinct peak value of RAO, the relative motion appears to
236 have a similar tendency to that of the inner float. When the wave period is small, the inner float has
237 a weak response, and the performance of the system mainly depends on the response of the outer
238 one. As the period increases to the resonance range of the inner float, it interacts strongly with wave,
239 which leads to a large relative motion. When the wave period continues increasing, the motions of
240 two floats tend to be consistent, and the relative motion is reduced. Regardless of large wave periods,
241 the relative motion has a better performance than any of the two single floats, especially when the
242 period is small. The opposite phenomenon occurs under large wave periods.

243 In summary, the two-body heaving WEC has better response than a single-body one under shorter
244 wave periods, the motion characteristics of the outer and inner floats are mixed to supply a greater
245 contribution to the united system. The heave motion RAO curves show that this model is suitable
246 for the selected sea state with the outer float as the absorber. The numerical model is valid under
247 free oscillation when supplied with experimentally derived added mass and damping coefficients.

248 **4.3 PTO force and relative float motion in regular waves**

249 The cases of constant PTO damping force supplied by the hydraulic system in physical model test
250 are taken into account. Fig. 10(a) and (b) illustrate the comparisons of the changes of PTO damping
251 force with time and with relative motion. The data are recorded from time domain numerical model
252 and physical model test. Time series of relative motion is shown in Fig. 10(c). Numerical results are
253 more stable and higher than experimental ones, but the periodicity is good. Fig. 10(d) and (e) show
254 the comparisons of the amplitudes of PTO damping force, relative motion and relative velocity

255 between physical model test and time domain simulation results under two regular wave cases with
256 $T=1.30$ s, $H=0.175$ m (case 1) and $H=0.20$ m (case 2), respectively. The trends of curves between
257 them are quite similar. PTO damping force is measured by pressure sensor in physical model test,
258 while the numerical model results are the configured values. As shown in Fig. 10, the PTO damping
259 force applied to the WEC in the physical model test is relatively smaller compared to the numerical
260 model. However, focusing on WEC relative motion and relative velocity, numerical results perform
261 better than experimental ones. That may cause by the friction loss of the hydraulic system. The PTO
262 damping force applied in the model test is obtained by multiplying the pressure difference of the
263 hydraulic cylinder by the cavity area, ignoring the friction of the hydraulic system. Therefore, the
264 measured PTO damping force in the physical model test is less than the actual damping force applied
265 to the floats and the relative motion is smaller than numerical ones. It is believed that in the full-
266 scale model, the proportion of friction loss is small, so that it is not taken into account in the
267 following numerical model. In Fig. 10(f) and (g), the comparisons of capture width ratio (CWR)
268 under case 1 and 2 are investigated, which are the combination of PTO force and relative motion
269 (shown in Fig. 10(d) and (e)). There is a certain difference in absolute values, but the trends of the
270 curve are consistent.

271 For more proof, the comparisons of the optimal coefficient between physical model test and time
272 domain simulation are illustrated in Fig. 10(h). In addition to case 1 and 2, case 3 ($T=1.55$ s, $H=0.175$
273 m) and case 4 ($T=1.55$ s, $H=0.20$ m) are added for comparison. It can be seen that there is a good
274 match between the physical and simulation model in terms of the optimal coefficient. It can also be
275 proved that the constant PTO damping force can be applied in the numerical simulation.

276 Based on the above discussions, the numerical dynamics model agrees well with experimental

277 results in terms of displacements, RAO, PTO force and power capture. The numerical dynamics
278 model is considered to be valid for the purposes of this study. To further explore the properties of
279 WEC, the below part is discussed by simulation model data. What's more, in order to show the
280 performance in more detail, the simulation extends the wave conditions to the wave height between
281 0.05-0.08 m and the wave period between 0.80-2.60 s.

282 **5 Power capture performance with different PTOs for regular wave** 283 **conditions**

284 Sensitivity study of power capture performance of the two-body heaving WEC is conducted in this
285 section using time-domain numerical simulation models. Power capture performance under
286 different PTO models, wave heights and periods are compared. Phase angle between different PTO
287 damping forces are also contrasted. All of the analyses are based on the time domain simulation
288 results. Nondimensional numbers are suggested to obtain the laws of performance for upscaled
289 WEC. Table 3 shows the dimensional and the corresponding nondimensional quantities used in this
290 section.

291 **5.1 Power capture performance for different wave heights**

292 Data from time domain model results with constant PTO damping are given in Fig. 11. The wave
293 heights are $H=0.10-0.20$ m, with the interval of 0.02 m, and $T=1.40$ s. The data plotted are
294 mechanical power capture and CWR as a function of PTO damping force. The energy captured is
295 calculated by multiplying PTO damping with relative motion during the time history. As shown in
296 Fig. 11(a), the higher wave height is, the more power produces. The optimal PTO damping force
297 increase with the increase of the wave height. Meanwhile, the CWR is also proportional to the wave

298 height, as shown in Fig. 11(b). Normalize the PTO damping force by $\rho g D^2 H$, where ρ is wave
299 density, g is gravitational acceleration, D is the diameter of the outer float and H is the wave
300 height. For one certain wave period, the curve of mean power is consistent with the change of
301 normalized force under various wave heights. In respect of linear PTO damping force condition, the
302 mean power and its corresponding CWR increase first and then decrease, along with the increase of
303 damping coefficient, as shown in Fig. 12. The captured power is proportional to the wave height,
304 while the curves of CWR at different wave heights are coincident. In contrast to constant PTO force,
305 the optimal PTO damping coefficient is fixed for one wave period. What's more, the energy
306 acquisition of the device under linear damping is more moderate with the change of damping
307 coefficient.

308 **5.2 Optimal PTO damping coefficients for the two-body heaving WEC**

309 The curves of optimal coefficients of PTO damping force, which is defined as the corresponding
310 damping coefficient when the maximum energy is reached at certain period, are shown in Fig. 13.
311 For constant damping force, as seen in Fig. 13(a) and (b), the optimal coefficient goes up to a peak
312 value and subsequently declines with the increase of wave period. In particular, the optimal damping
313 coefficient reaches the maximum value under the optimal capture wave period. Obviously, higher
314 wave heights will have larger optimal coefficient. When the optimal coefficient normalized by
315 $\rho g D^2 H$, it only depends on the wave period. Accordingly, the optimal coefficient of constant
316 damping is proportional to wave height. Fig. 13(c) and (d) show that the optimal damping coefficient
317 of linear PTO damping force decreases first with respect to wave period and normalized wave period,
318 and then increases again when the period gets to the synchronous motion of the two floats. Wave
319 height has no effect on the optimal coefficient under linear PTO damping force. It can be seen that

320 the optimized damping coefficients can be so different under different periods.

321 The characteristics of optimal damping and power capture are so different due to hydrodynamic
322 properties, as expected. For constant PTO damping force, both wave height and period affect the
323 power capture and optimal coefficient, whereas linear damping, only wave period does.

324 **5.3 Sensitivity study for regular wave conditions**

325 When the PTO damping coefficient is fixed, constant PTO damping force exhibits different
326 capacities under different wave height, while linear PTO damping force remains consistent, as
327 shown in Fig. 14. This is because, as discussed above, the optimal coefficient that can be obtained
328 at different wave heights are different with constant PTO damping force, and in regard to linear
329 system, only the incident wave period matters.

330 Observing from Fig. 9(c), the peak value of relative heave motion RAO appears at 1.80 s. However,
331 the high relative motion amplitude is associated with low power production, as shown in Fig. 14(a)
332 and (c). Therefore, blindly increasing the relative motion does not improve energy production. Due
333 to the influence of PTO damping force, the maximum response interval of the system changes,
334 moving towards the shorter wave periods. The results of the model show that the system has a
335 good power capacity under small wave periods between 1.40-1.60 s, that is, between the
336 resonance periods of the two bodies. The ability of power capture gets worse when wave period
337 exceeds the optimal range, and synchronism occurs.

338 The power capture performance with different constant PTO damping force varies with the wave
339 height. Observing the curves of mean power in Fig. 15(a), as the PTO damping rising, the curvature
340 of energy growth ascends. For its corresponding CWR, as shown in Fig. 15(b), CWR falls down

341 with the increase of the wave height under small constant PTO damping force, then as the force
342 increases, the degree of decrease becomes slower and turns to an upward tendency, and from the
343 convex curve to concave curve, gradually. Therefore, in the power capture range, as the constant
344 PTO force goes up, the energy obtained will ascend at a faster rate as the wave height goes up. Differ
345 from constant force, wave height improves the mean power capture monotonically under linear
346 damping, while the corresponding CWR does not change, as shown in Fig. 15(c) and (d).

347 **5.4 Power capture performance with different PTO models**

348 The comparisons of optimal mean power capture as a function of wave period between constant and
349 linear PTO forces are shown in Fig. 16. When the wave period is shorter than 1.80 s, the resonance
350 period of the inner float, linear PTO damping force are better than the constant one. When the period
351 exceeds 1.80 s, the difference is very tiny. Wave height has no impact on the optimal power. For
352 CWR comparison, wave height has no influence on the maximum power capture.

353 As introduced in Section 3, the PTO damping force yields $f_{PTO} = -B \operatorname{sgn}(\dot{z}_1 - \dot{z}_2) |\dot{z}_1 - \dot{z}_2|^{\frac{1}{n}}$. The
354 PTO nonlinear factor n is introduced to evaluate power capacity, which is used to characterize the
355 nonlinearity of PTO damping force. In order to better display in the figure, here, assuming that the
356 PTO nonlinear factor of constant PTO damping force is 150. As shown in Fig. 17, the maximum
357 power increases first and then drops down with the increase of n , and reaches the maximum when
358 $n = 1$, that is, the linear damping. In this comparative analysis, the PTO damping coefficient is set
359 to an optimized value that is specific to each PTO and wave conditions. It indicates that the linear
360 damping is the optimal PTO damping form to this system.

361 Contour plots of mean power production and CWR as a function of H and T for constant and

362 linear PTO damping forces are shown in Fig. 18. The optimum passive PTO damping force was
363 chosen for each wave condition. Linear PTO damping force has a better power capture performance
364 than constant force under the same sea state, but with very little difference. In that sense, changing
365 the type of PTO damping force can only improve the capacity to a limited extent. The main power
366 band is concentrated between the heave natural periods of the two bodies. Wave period
367 corresponding to the peak value of the mean power is longer than the period corresponding to the
368 peak value of CWR.

369 **5.5 Phase angle with different PTO models**

370 The phase angle is used to describe the phase difference between the outer and the inner floats. The
371 phase angle is larger in the absence of PTO forces than that in the presence of PTO force. The
372 maximum phase angle without PTO force is 125° at 1.60 s, and approximately zero phase angle
373 above 2.0 s. As shown in Fig. 19, it is very difficult to make these two bodies move in opposite
374 direction under passive damping. They get a maximum phase angle at 1.80 s, which is the natural
375 period of the inner float. The PTO damping force reduces phase angle by 55%. When the wave
376 period exceeds 1.80 s, the phase angle plunged to zero.

377 **6 Conclusions**

378 In this paper, a 1:9 scaled model of a two-body heaving WEC is studied both experimentally
379 and numerically, with focus on the power absorption performance and the motion
380 characteristics of the WEC. The linearized viscous damping coefficients derived from the
381 physical decay model test is used to calibrate the numerical model to achieve a good
382 comparison to the experimental results. The numerical result under regular wave condition has

383 a good agreement with that of the physical model test. Sensitive study of this system is
384 conducted using the validated time-domain numerical models, and the conclusion that larger
385 relative movement does not mean better power capture is drawn. Linear PTO damping system
386 is more suitable for this system with consideration of the PTO nonlinear factor n . For the
387 different wave heights that were considered in this study, the WEC system behaves linearly.

388 The hydrodynamic performance of this system depends on the coupling property of the two
389 floats and has a decent motion response in resonance range. The power production is markedly
390 affected by the wave period. Phase angle between the two floats is not the only indicator of
391 power capacity. In terms of power capture, the system has a good performance when the wave
392 period is between the two floats' natural periods and especially at the average value of their
393 natural periods. As a result, a larger separation between two natural periods is beneficial,
394 leading to a wider and higher capture band.

395 The paper addresses the dynamic performance of the two-body WEC, with no active phase
396 control, which will be applied in the future.

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402 **References**

403 Abdelkhalik, O., Zou, S., 2019. Control of small two-body heaving wave energy converters for ocean

404 measurement applications. *Renew. Energy* 132, 587-595.
405 <https://doi.org/10.1016/j.renene.2018.08.004>.

406 Al-Habaibeh, A., Su, D., McCague, J., Knight, A., 2010. An innovative approach for energy generation
407 from waves. *Energy Conversion and Management* 51 (8), 1664-1668.
408 <https://doi.org/10.1016/j.enconman.2009.11.041>.

409 Amiri, A., Panahi, R., Radfar, S., 2016. Parametric study of two-body floating-point wave absorber. *J.*
410 *Mar. Sci. Appl.* 15 (1), 41-49. <https://doi.org/10.1007/s11804-016-1342-1>.

411 Anderlini, E., Forehand, D.I.M., Bannon, E., Xiao, Q., Abusara, M., 2018. Reactive control of a two-
412 body point absorber using reinforcement learning. *Ocean. Eng.* 148, 650-658.
413 <https://doi.org/10.1016/j.oceaneng.2017.08.017>.

414 Beatty, S.J., Bocking, B., Bubbar, K., Buckham, B.J., Wild, P., 2019. Experimental and numerical
415 comparisons of self-reacting point absorber wave energy converters in irregular waves. *Ocean. Eng.*
416 173, 716-731. <https://doi.org/10.1016/j.oceaneng.2019.01.034>.

417 Beatty, S.J., Hall, M., Buckham, B.J., Wild, P., Bocking, B., 2015. Experimental and numerical
418 comparisons of self-reacting point absorber wave energy converters in regular waves. *Ocean. Eng.*
419 104, 370-386. <http://dx.doi.org/10.1016/j.oceaneng.2015.05.027>.

420 Bubbar, K., Buckham, B., 2020. On establishing generalized analytical phase control conditions in two
421 body self-reacting point absorber wave energy converters. *Ocean. Eng.* 197.
422 <https://doi.org/10.1016/j.oceaneng.2019.106879>.

423 Castro, F.A., Chiang, L.E., 2020. Design optimization and experimental validation of a two-body Wave
424 Energy Converter with adjustable Power Take-Off parameters. *Energy for Sustainable Development*
425 56, 19-32. <https://doi.org/10.1016/j.esd.2020.02.007>.

426 Cordonnier, J., Gorintin, F., De Cagny, A., Clement, A.H., Babarit, A., 2015. SEAREV: Case study of
427 the development of a wave energy converter. *Renew. Energy* 80, 40-52.
428 <https://doi.org/10.1016/j.renene.2015.01.061>.

429 Crowley, S., Porter, R., Taunton, D.J., Wilson, P.A., 2018. Modelling of the WITT wave energy converter.
430 *Renew. Energy* 115, 159-174. <https://doi.org/10.1016/j.renene.2017.08.004>.

431 Dai, Y.M., Chen, Y.Z., Xie, L.H., 2017. A study on a novel two-body floating wave energy converter.
432 *Ocean. Eng.* 130, 407-416. <http://dx.doi.org/10.1016/j.oceaneng.2016.11.049>.

433 Falnes J. 1999. Wave-energy conversion through relative motion between two single-mode oscillating
434 bodies. *J Offshore Mech Arctic Eng.* 121:32–38. <https://doi.org/10.1115/1.2829552>.

435 Henderson, R., 2006. Design, simulation, and testing of a novel hydraulic power take-off system for the
436 Pelamis wave energy converter. *Ocean. Eng.* 31 (2), 271-283.
437 <https://doi.org/10.1016/j.renene.2005.08.021>.

438 Henriques, J.C.C., Lopes, M.F.P., Gomes, R.P.F., Gato, L.M.C., Falcao, A.F.O., 2012. On the annual
439 wave energy absorption by two-body heaving WECs with latching control. *Renew. Energy* 45, 31-
440 40. <https://doi.org/10.1016/j.renene.2012.01.102>.

441 Hernández, F., Chiang, L. E., & Corbalán, P. (2017). A general architecture for electric
442 powermanagement of small scale NCRE converters: Design methodology and validation. *Energy*
443 for Sustain. Development, Vol., 41, 128–138. <https://doi.org/10.1016/j.esd.2017.09.001>.

444 Ji, X.Y., Al Shami, E., Monty, J., Wang, X., 2020. Modelling of linear and non-linear two-body wave
445 energy converters under regular and irregular wave conditions. *Renew. Energy* 147, 487-501.
446 <https://doi.org/10.1016/j.renene.2019.09.010>.

447 Jin, P., Zhou, B.Z., Goteman, M., Chen, Z.F., Zhang, L., 2019. Performance optimization of a coaxial-

448 cylinder wave energy converter. *Energy* 174, 450-459. <https://doi.org/10.1016/j.energy.2019.02.189>.

449 Liang, C.W., Zuo, L., 2017. On the dynamics and design of a two-body wave energy converter. *Renew.*
450 *Energy* 101, 265-274. <http://dx.doi.org/10.1016/j.renene.2016.08.059>.

451 Liu, C., Yang, Q., Bao, G., 2018. Influence of hydraulic power take-off unit parameters on power capture
452 ability of a two-raft-type wave energy converter. *Ocean. Eng.* 150, 69-80.
453 <https://doi.org/10.1016/j.oceaneng.2017.12.063>.

454 Liu, Z., Qu, N., Shi, H.D., 2017. Experimental study on hydrodynamic performance of a wave energy
455 converter within multi-heaving-buoys. *International J. Energy Res.* 41 (9), 1351-1366.
456 <https://doi.org/10.1002/er.3725>.

457 Martin, D., Li, X.F., Chen, C.A., Thiagarajan, K., Ngo, K., Parker, R., Zuo, L., 2020. Numerical analysis
458 and wave tank validation on the optimal design of a two-body wave energy converter. *Renew.*
459 *Energy* 145, 632-641. <https://doi.org/10.1016/j.renene.2019.05.109>.

460 Neary, V.S., Lawson, M., Previsic, M., Copping, A., Hallett, K.C., LaBonte, A., Rieks, J., Murray, D.,
461 2014. Methodology for Design and Economic Analysis of Marine Energy Conversion (MEC)
462 Technologies. Sandia National Lab.(SNL-NM), Albuquerque, NM (United States).

463 Negandari, M., Dalayeli, H., Moghadas, M.H., 2018. Design of a Two-Body Wave Energy Converter by
464 Incorporating the Effect of Hydraulic Power Take-Off Parameters. *Journal of Marine Science and*
465 *Technology-Taiwan* 26 (4), 496-507. [https://doi.org/10.6119/JMST.201808_26\(4\).0003](https://doi.org/10.6119/JMST.201808_26(4).0003).

466 Ocean Energy System. 2018. Spotlight on ocean energy.

467 Rusu, L., Onea, F., 2017. The performance of some state-of-the-art wave energy converters in locations
468 with the worldwide highest wave power. *Renew. Sustain. Energy Rev.* 75, 1348-1362.
469 <https://doi.org/10.1016/j.rser.2016.11.123>.

470 Shi, H.D., Dong, X.C., Feng, L., Han, Z., 2019. Experimental Study on the Hydrodynamic Performance
471 of a Heaving Buoy Assembled on a Net Cage Platform. *J. Ocean Univ. China* 18 (5), 1031-1040.
472 <https://doi.org/10.1007/s11802-019-4028-x>.

473 Son, D., Belissen, V., Yeung, R.W., 2016. Performance validation and optimization of a dual coaxial-
474 cylinder ocean-wave energy extractor. *Renew. Energy* 92, 192-201.
475 <https://doi.org/10.1016/j.renene.2016.01.032>.

476 Son, D., Yeung, R.W., 2014. Performance Predictions and Validation of a Two Coaxial-Cylinder System
477 as a Wave-Energy Extractor. 33rd International Conference on Ocean, Offshore and Arctic Eng.,
478 2014, Vol 7.

479 Tan, Y.M., Liu, N., Lin, K.J., Zhang, Z.G., 2020. Frequency Domain Modeling of a Halbach PM Linear
480 Generator Based Two-Body Point Absorber for Wave Energy Conversion. *Front. Energy Res.* 8.
481 <https://doi.org/10.3389/fenrg.2020.00019>.

482 Tethys, 2019. Wello Penguin at EMEC. <https://tethys.pnnl.gov/annex-iv-sites/wello-penguin-emec>
483 (Accessed 23 July 2020)

484 Van den Berg, J., Ricci, P., Touzon, I., Villate, J.L., Tedeschi, E., 2011. Control strategies for improving
485 power performance of two-body heaving wave energy devices, The Twenty-first International
486 Offshore and Polar Engineering Conference. International Society of Offshore and Polar Engineers.

487 Van Rij, J., Yu, Y.-H., Edwards, K., Mekhiche, M., 2017. Ocean power technology design optimization.
488 *International Journal of Marine Energy* 20, 97-108. <https://doi.org/10.1016/j.ijome.2017.07.010>.

489 Windt, C., Davidson, J., Ringwood, J.V., 2018. High-fidelity numerical modelling of ocean wave energy
490 systems: A review of computational fluid dynamics-based numerical wave tanks. *Renew. Sustain.*
491 *Energy Rev.* 93, 610-630. <https://doi.org/10.1016/j.rser.2018.05.020>.

492 Xu, Q.L., Li, Y., Yu, Y.H., Ding, B.Y., Jiang, Z.Y., Lin, Z.L., Cazzolato, B., 2019. Experimental and
493 numerical investigations of a two-body floating-point absorber wave energy converter in regular
494 waves. *Journal of Fluids and Structures* 91. <https://doi.org/10.1016/j.jfluidstructs.2019.03.006>.

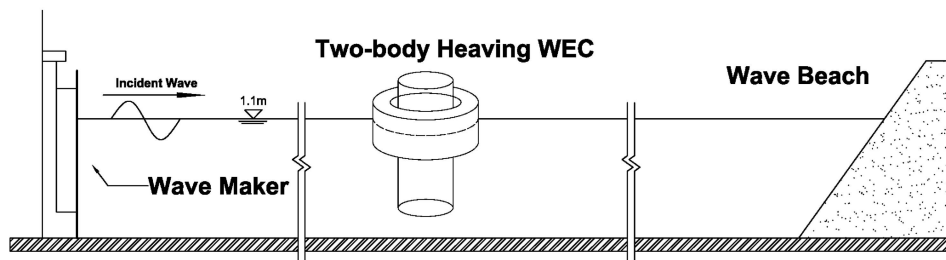
495 Zheng, C.W., Zhou, L., Jia, B.K., Pan, J., Li, X., 2014. Wave characteristic analysis and wave energy
496 resource evaluation in the China Sea. *Journal of Renew. and Sustain. Energy* 6 (4).
497 <http://dx.doi.org/10.1063/1.4885842>.



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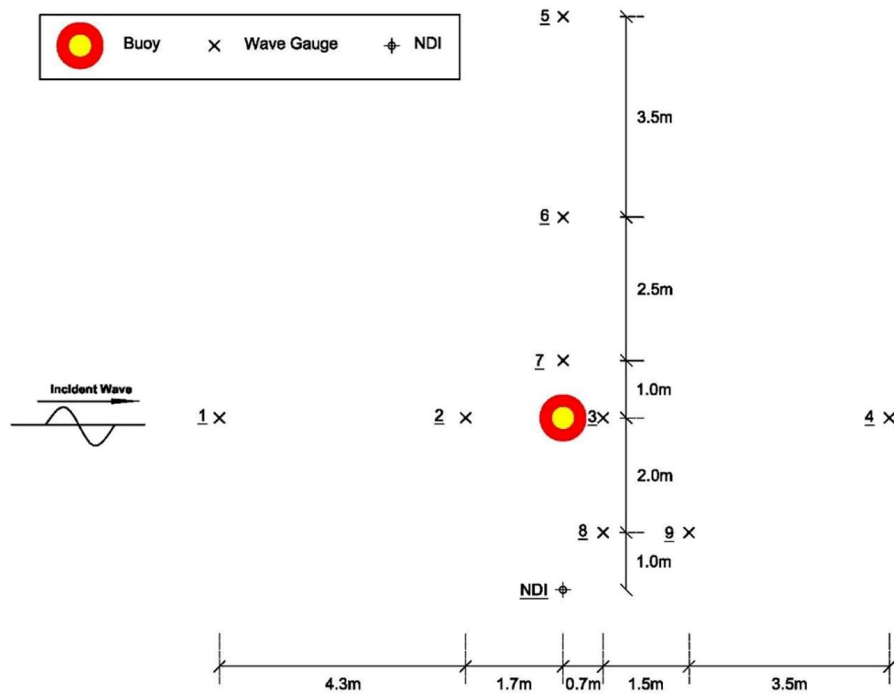
Fig. 1. Scale model in wave tank.



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Fig. 2. Wave tank layout for the model test.



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Fig. 3. The deployment of wave gauges and NDI.



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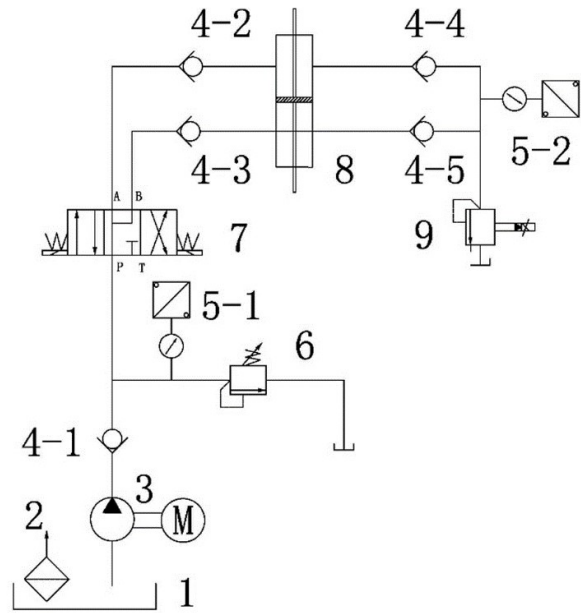
Fig. 4. The two-body heaving WEC.



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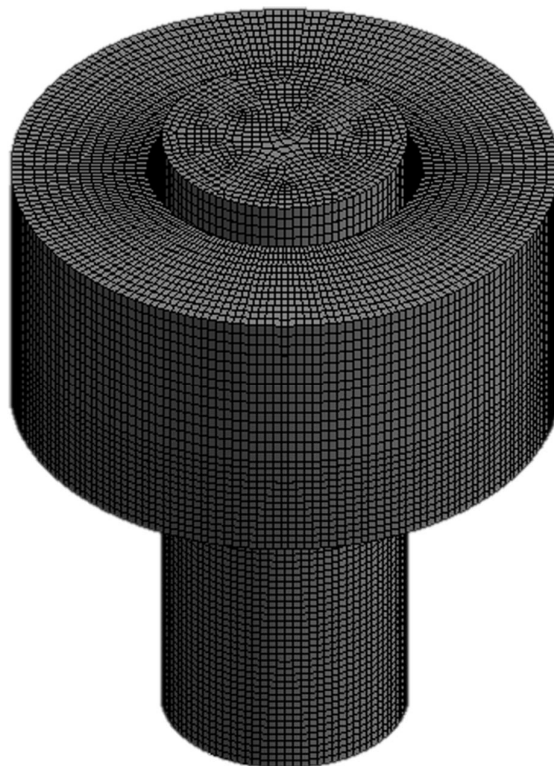
Fig. 5. (a) The hydraulic station.



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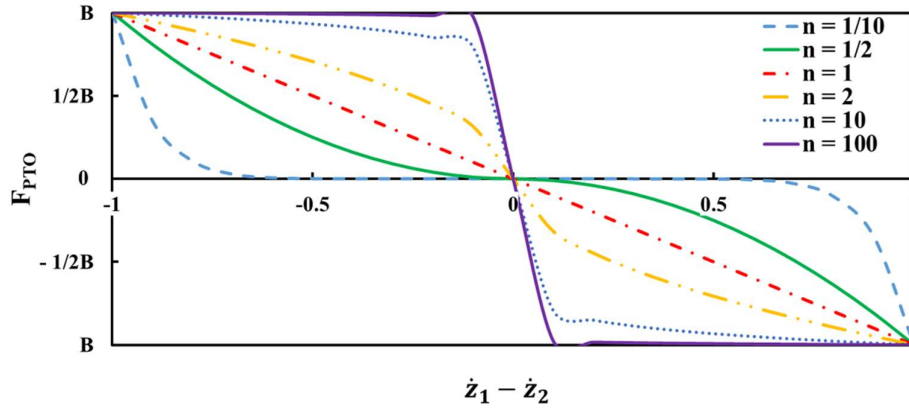
Fig. 5. (b) Schematic diagram of the hydraulic system.



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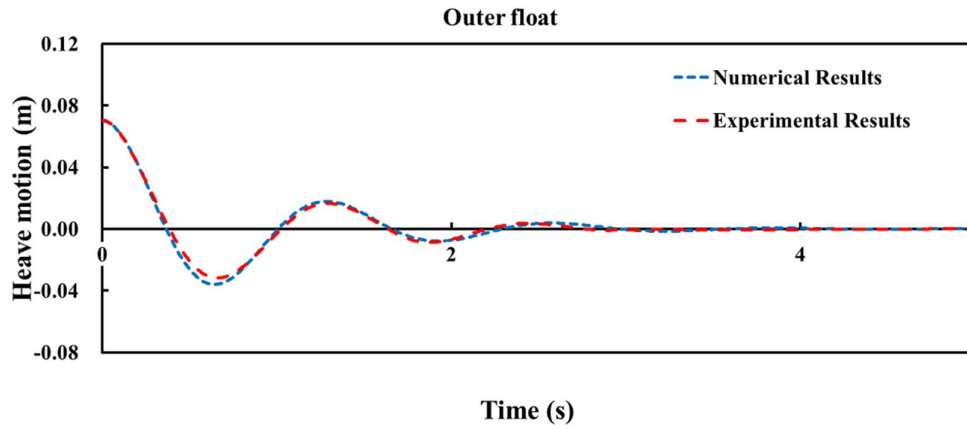
Fig. 6. The mesh used to simulate the heaving two-body WEC.



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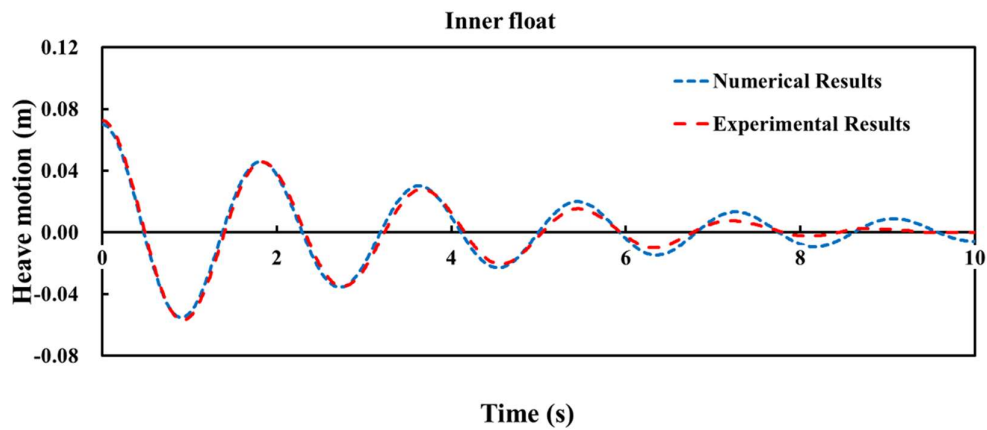
Fig. 7. Nonlinear PTO force for different n .



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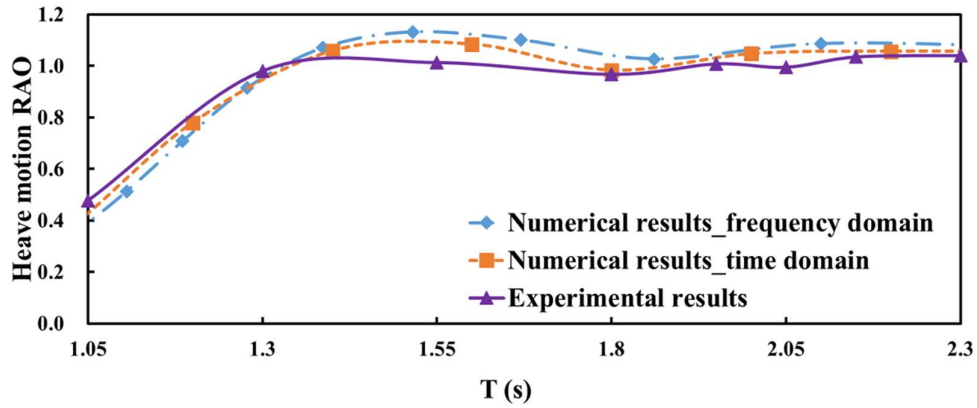
Fig. 8. (a) Heave decay curves of the outer float.



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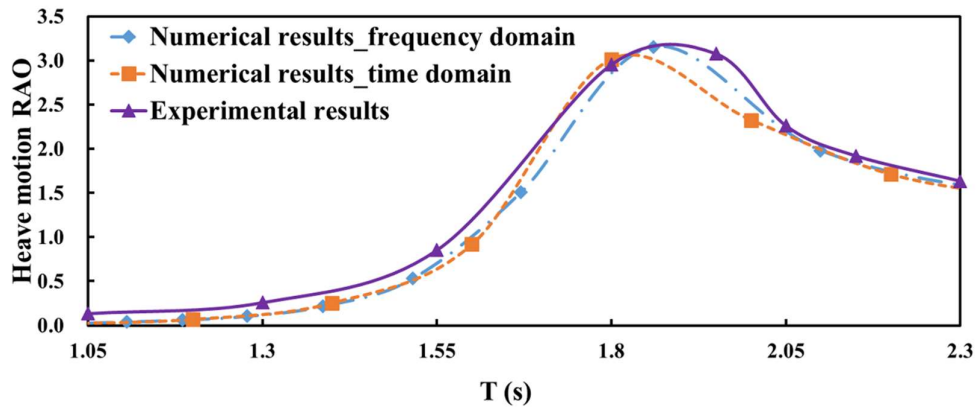
Fig. 8. (b) Heave decay curves of the inner float.



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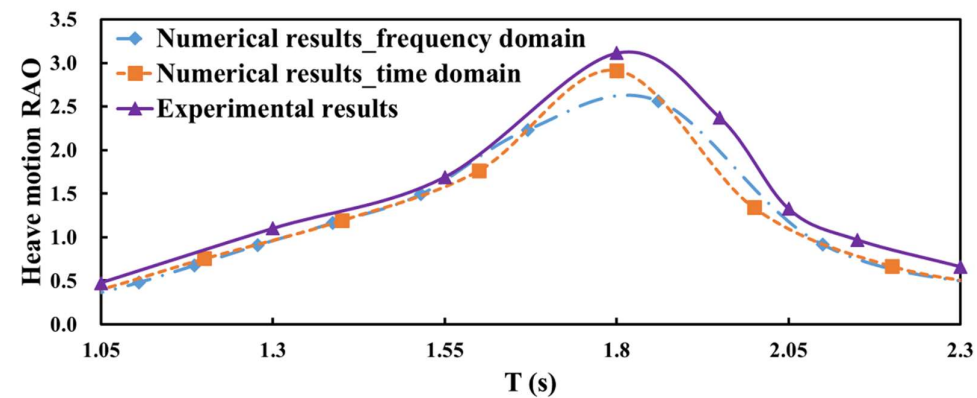
Fig. 9. (a) Heave motion RAO of the outer float.



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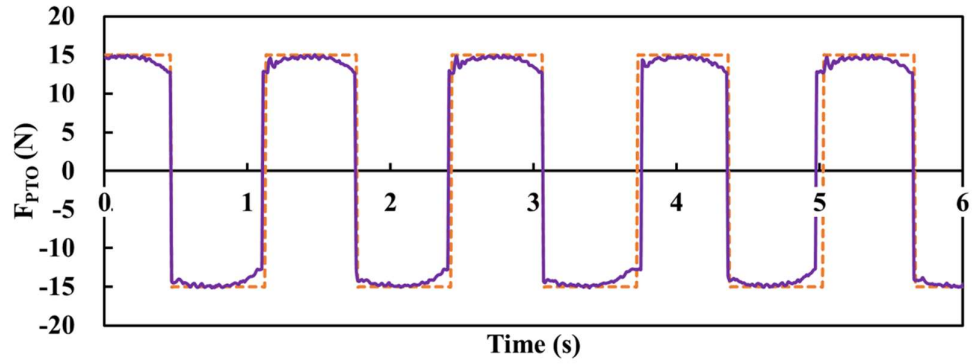
Fig. 9. (b) Heave motion RAO of the inner float.



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Fig. 9. (c) Relative heave motion RAO.

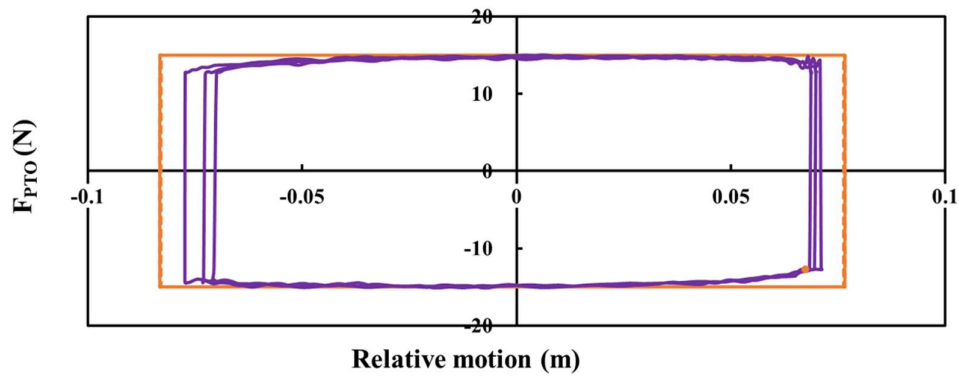


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--- Numerical results — Experimental results

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Fig. 10. (a) Constant PTO damping force time series.

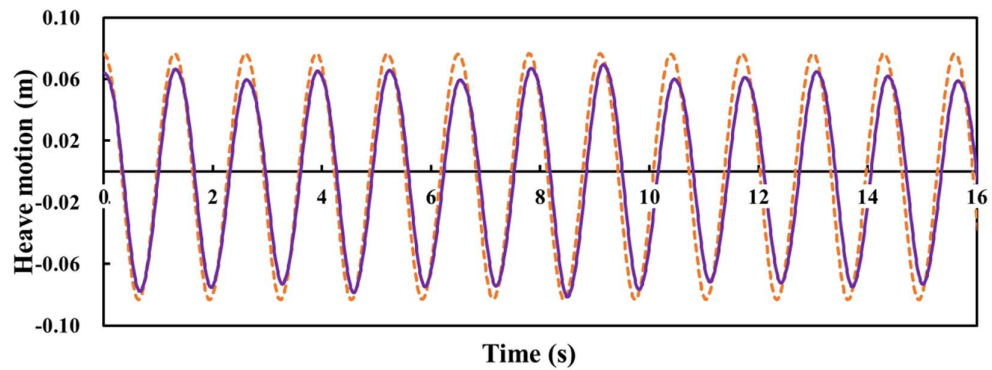


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--- Numerical results — Experimental results

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Fig. 10. (b) Constant PTO damping force as a function of the relative heave motion.



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--- Numerical results — Experimental results

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Fig. 10. (c) Relative motion time series.

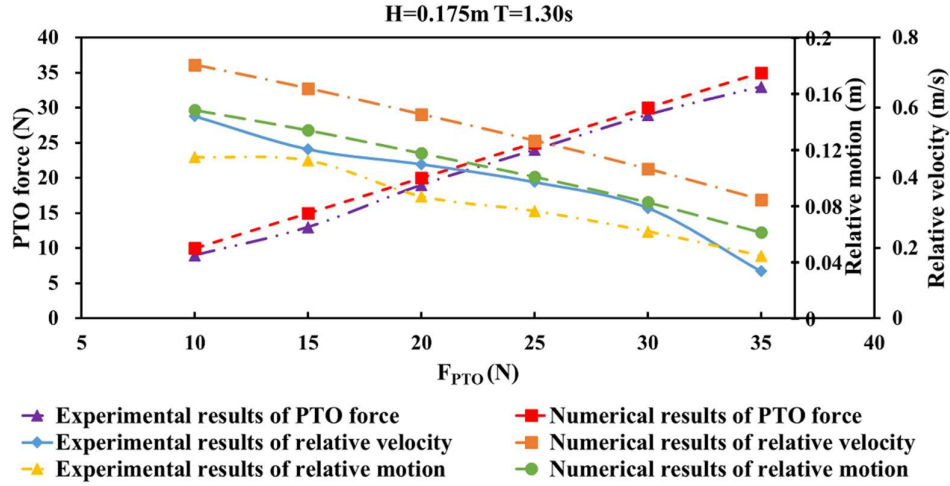


Fig. 10. (d) The comparison of experimental and numerical results under case 1.

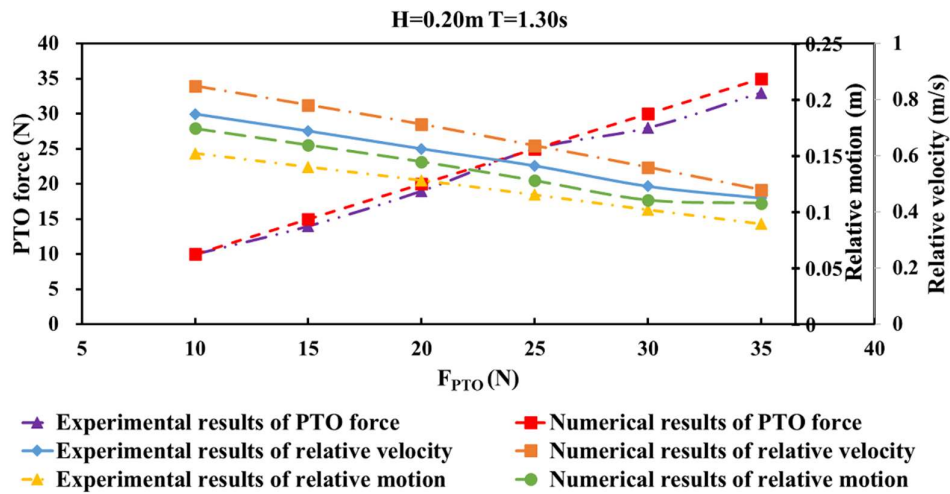


Fig. 10. (e) The comparison of experimental and numerical results under case 2.

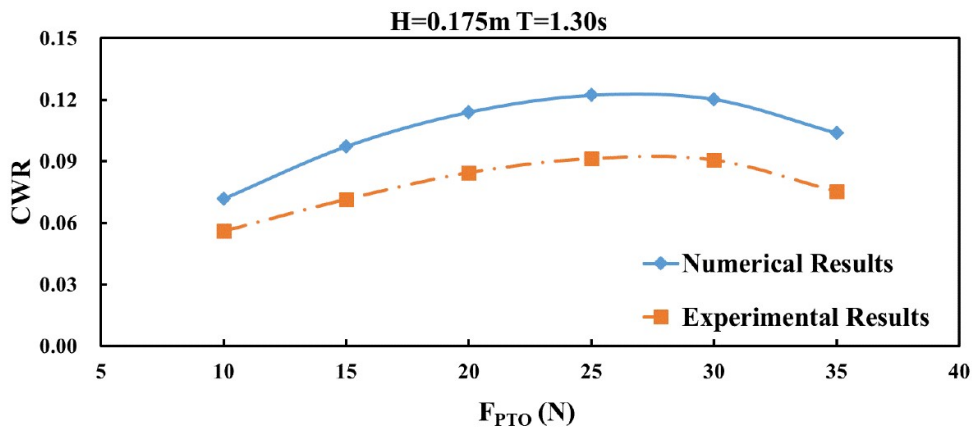
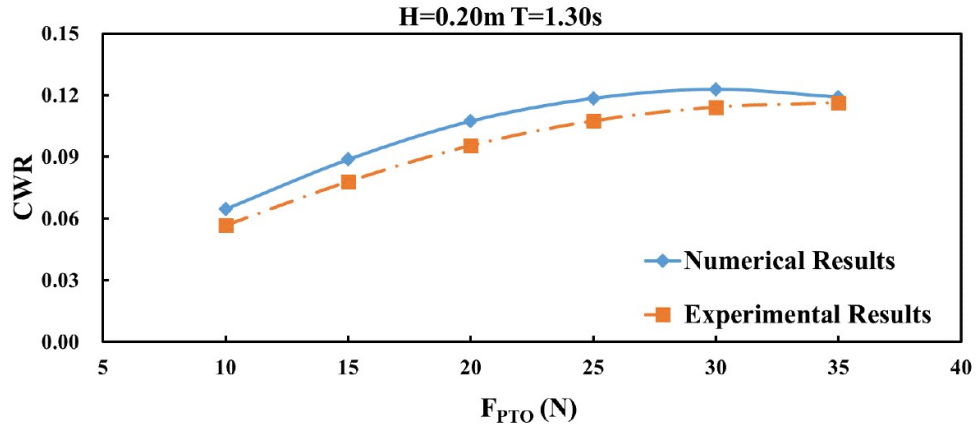


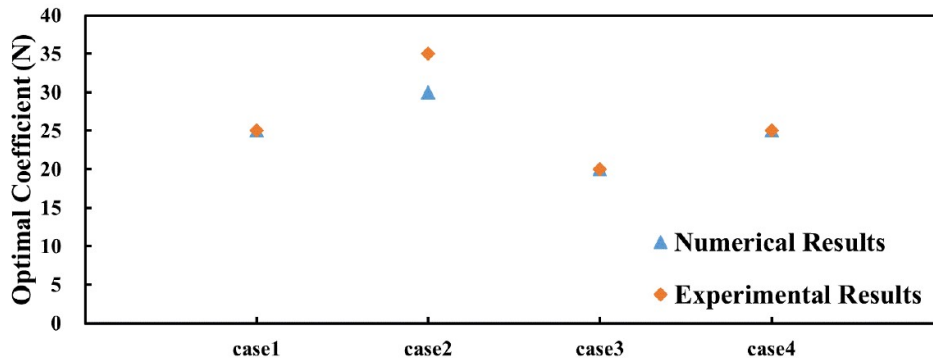
Fig. 10. (f) CWRs of experimental and numerical results under case 1.



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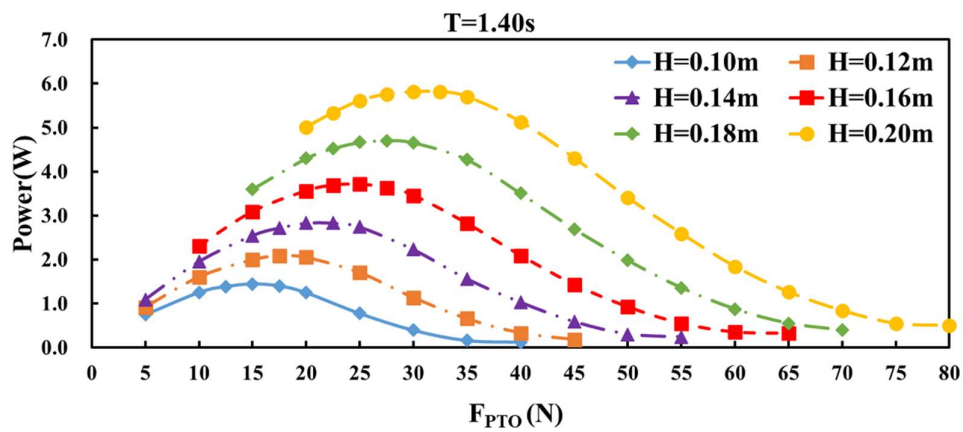
Fig. 10. (g) CWRs of experimental and numerical results under case 2.



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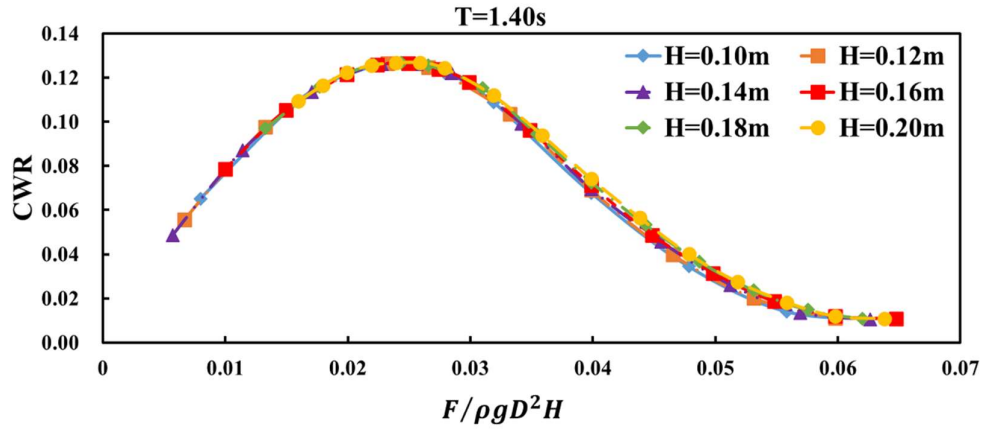
Fig. 10. (h) Optimal coefficient.



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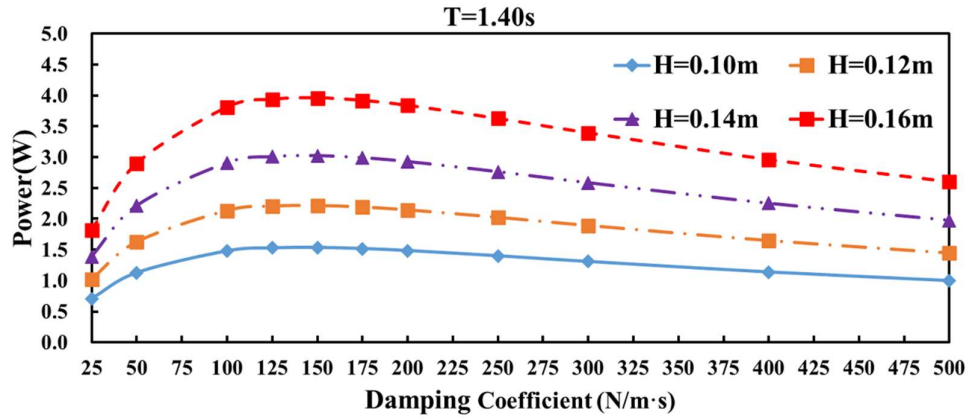
Fig. 11. (a) Mean power as a function of constant PTO damping force.



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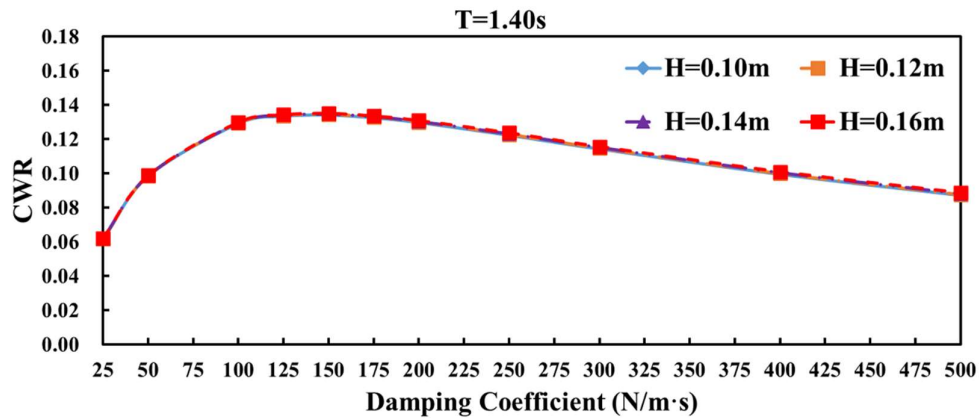
Fig. 11. (b) CWR as a function of normalized constant PTO damping force.



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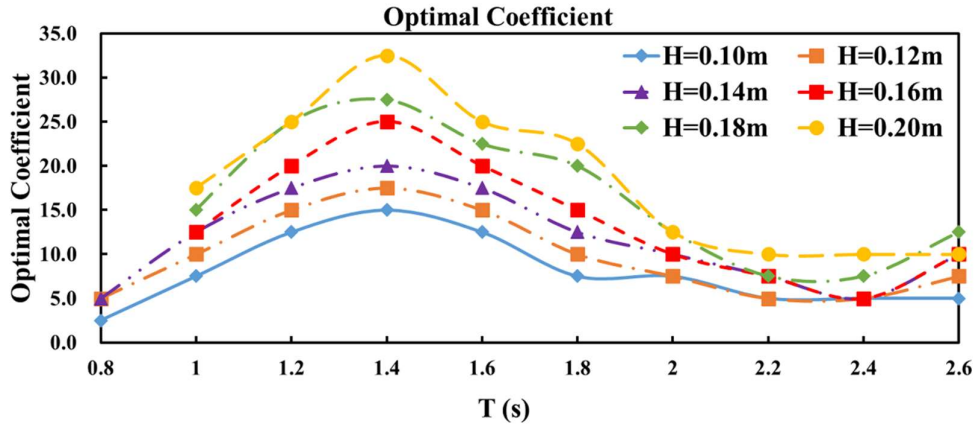
Fig. 12. (a) Mean power as a function of linear PTO damping force.



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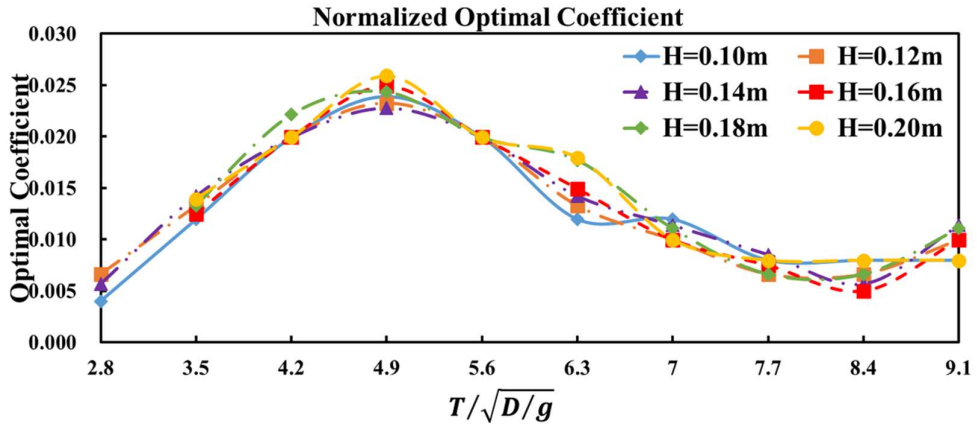
Fig. 12. (b) CWR as a function of linear PTO damping force.



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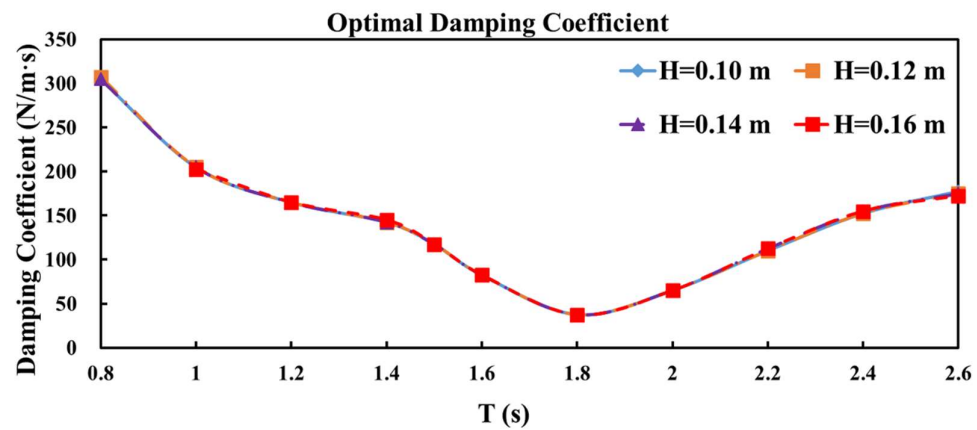
Fig. 13. (a) Optimal coefficient under constant PTO force.



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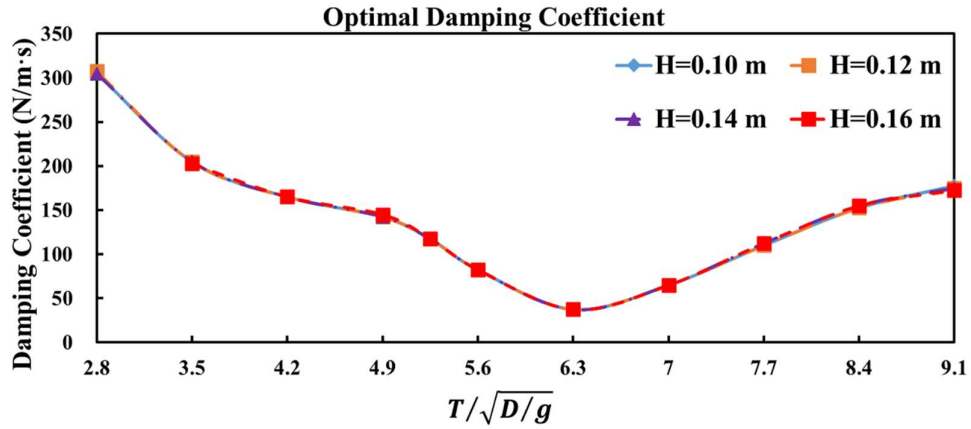
Fig. 13. (b) Normalized optimal coefficient under constant PTO force.



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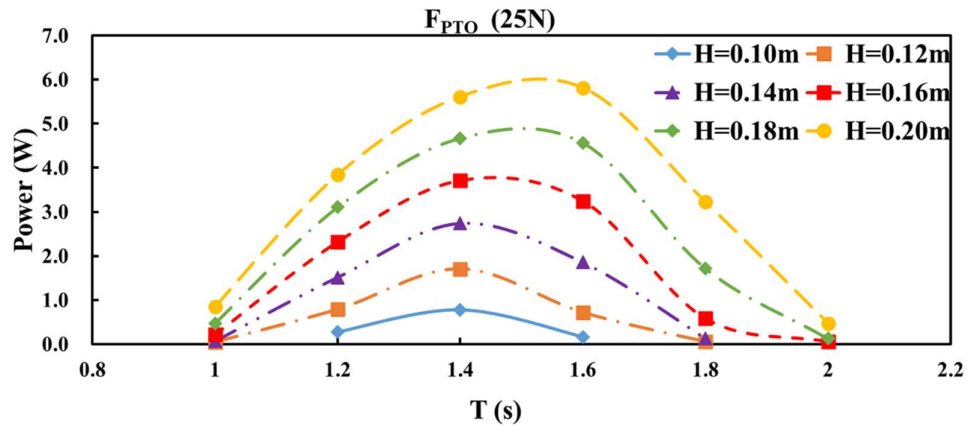
Fig. 13. (c) Optimal coefficient under linear PTO force.



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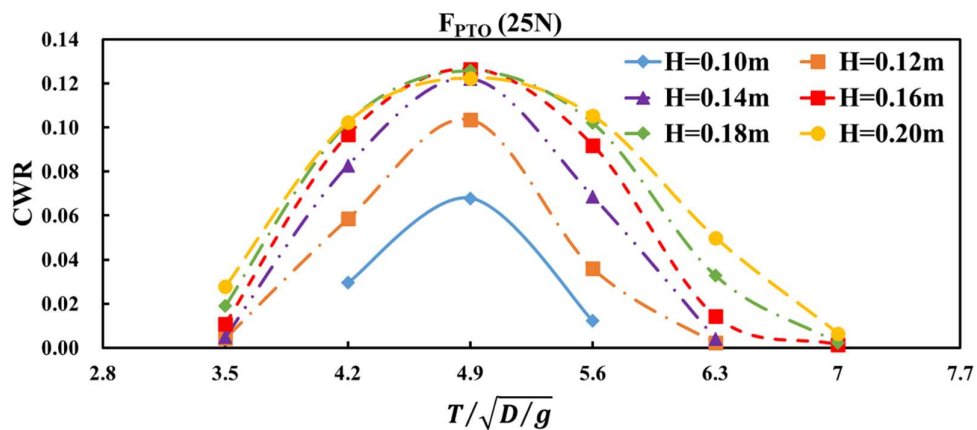
Fig. 13. (d) Optimal coefficient under linear PTO force.



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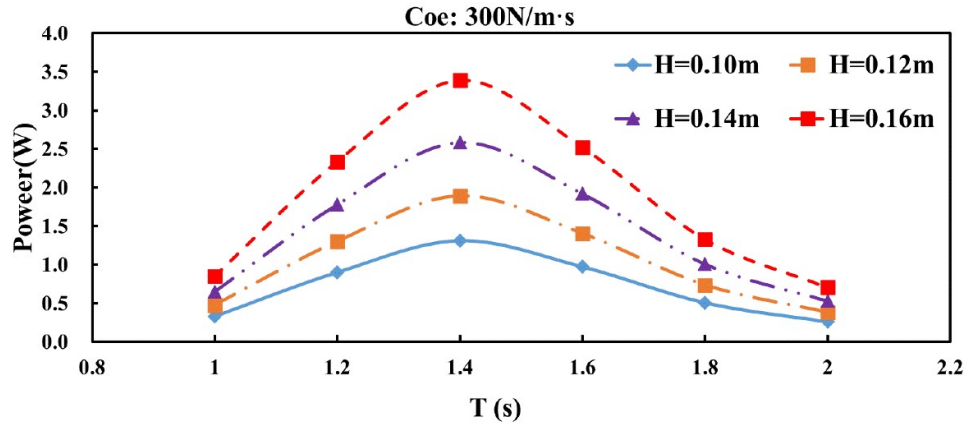
Fig. 14. (a) Capture power under constant PTO damping force.



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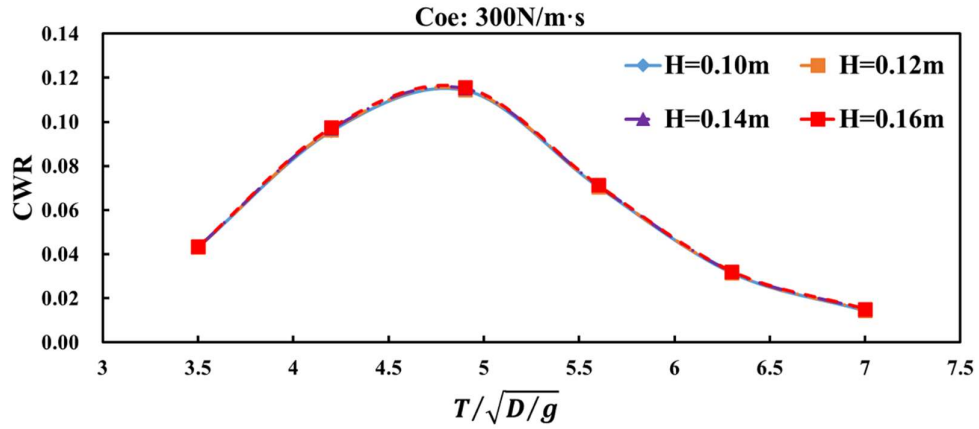
Fig. 14. (b) CWR under constant PTO damping force.



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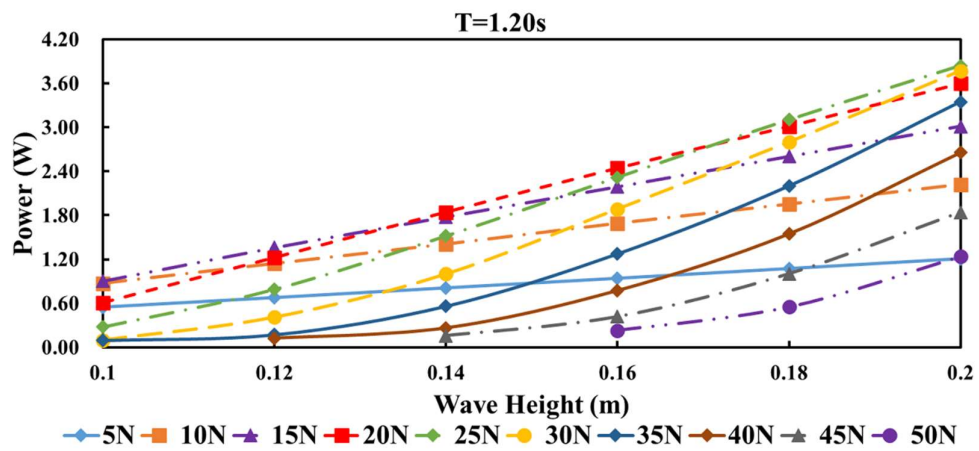
Fig. 14. (c) Capture power under linear PTO damping force.



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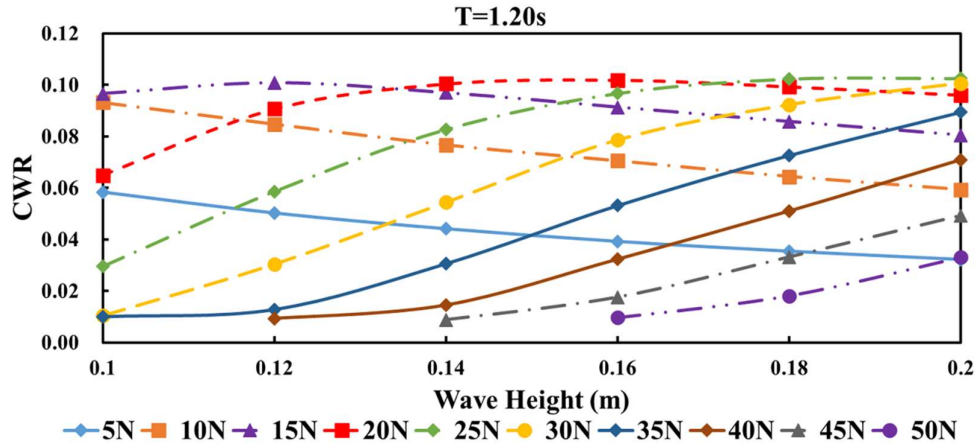
Fig. 14. (d) CWR under linear PTO damping force.



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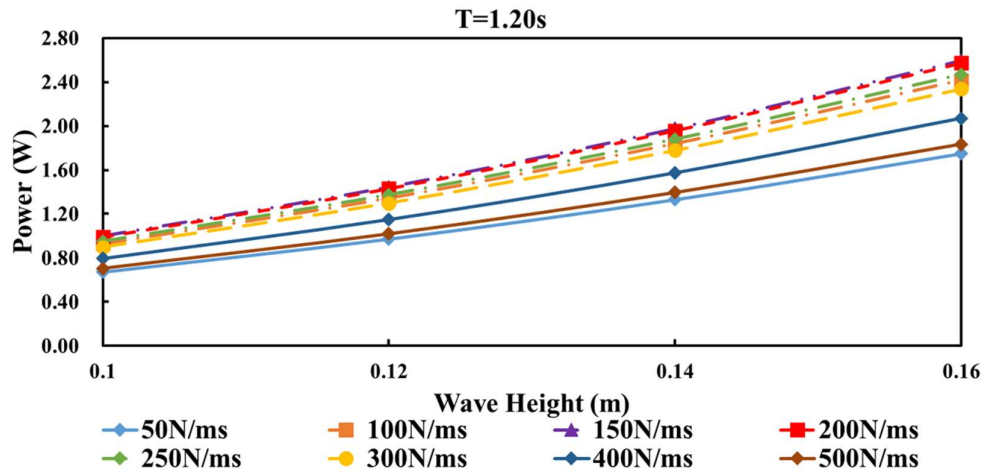
Fig. 15. (a) Capture power under constant PTO damping force.



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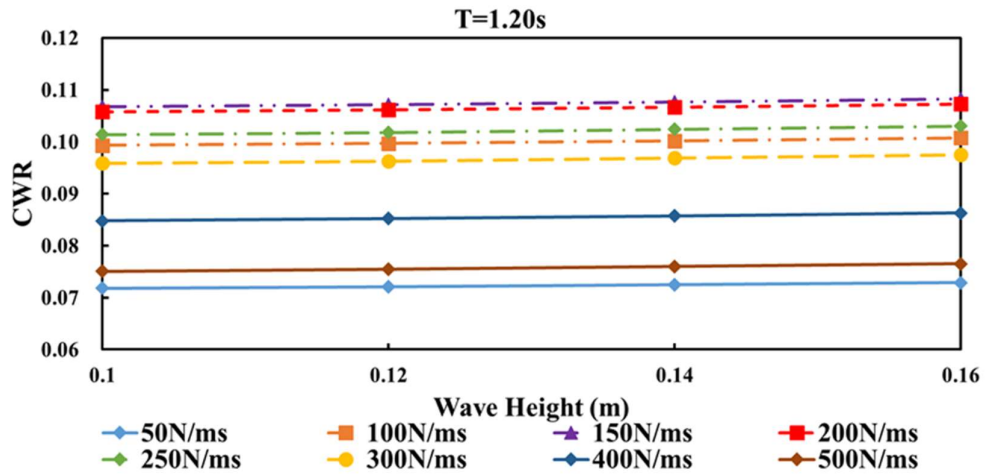
Fig. 15. (b) CWR under constant PTO damping force.



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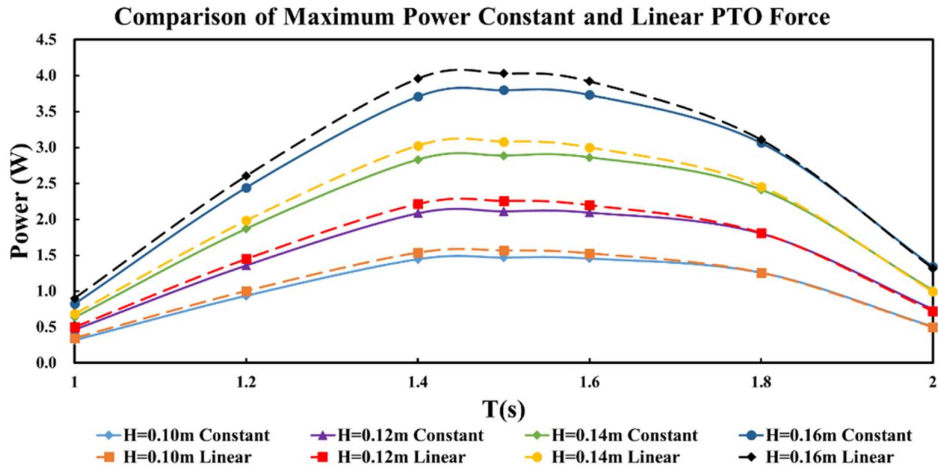
Fig. 15. (c) Capture power under linear PTO damping force.



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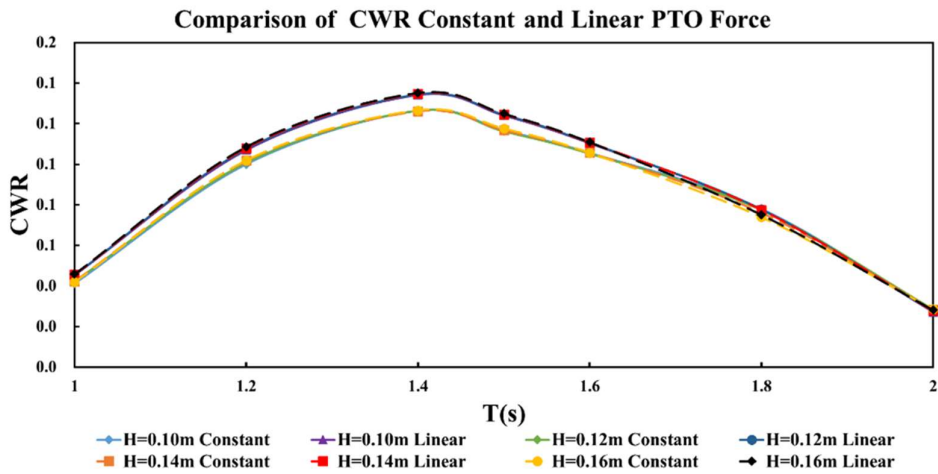
Fig. 15. (d) CWR under linear PTO damping force.



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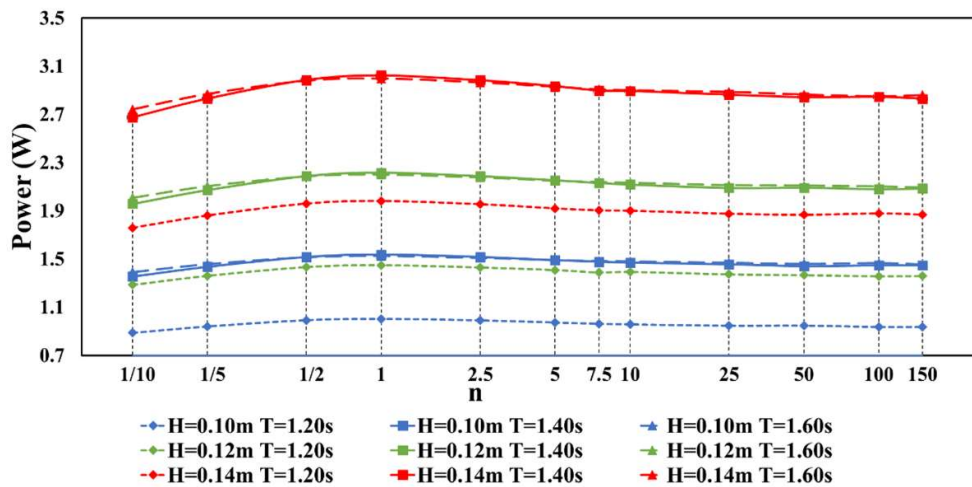
Fig. 16. (a) Comparison of capture power between constant and linear PTO force.



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Fig. 16. (b) Comparison of CWR between constant and linear PTO force.



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Fig. 17. (a) Capture power as a function of the PTO nonlinear factor.

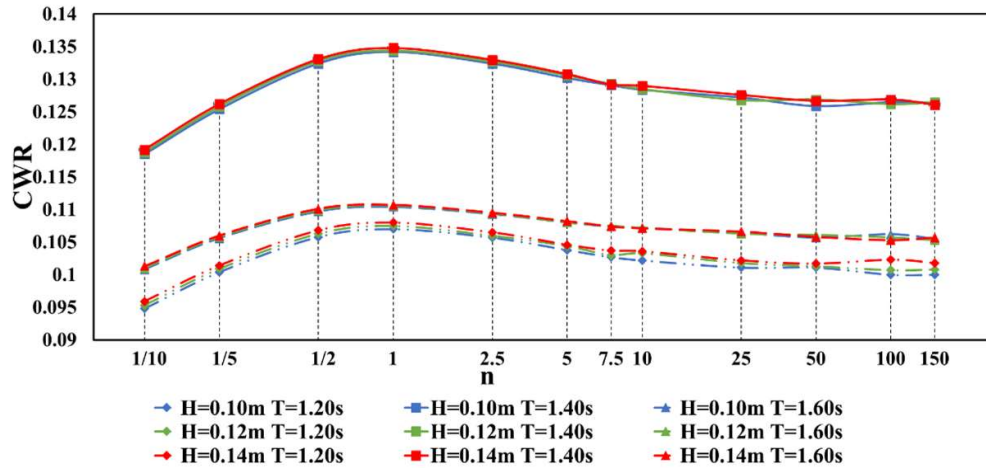


Fig. 17. (b) CWR as a function of the PTO nonlinear factor.

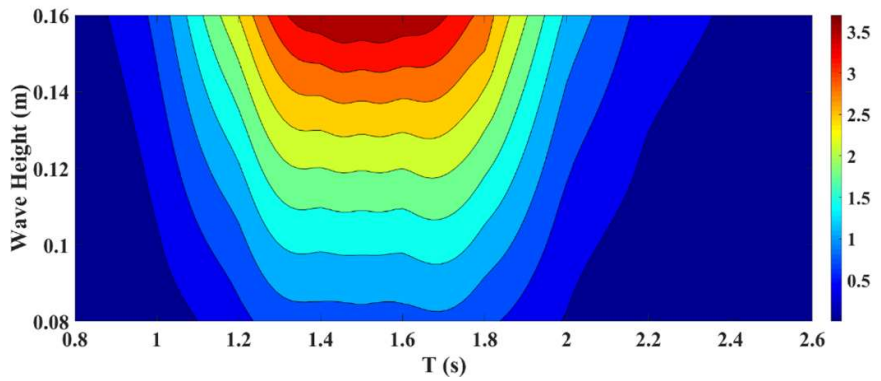


Fig. 18. (a) Capture Power contours with constant PTO force.

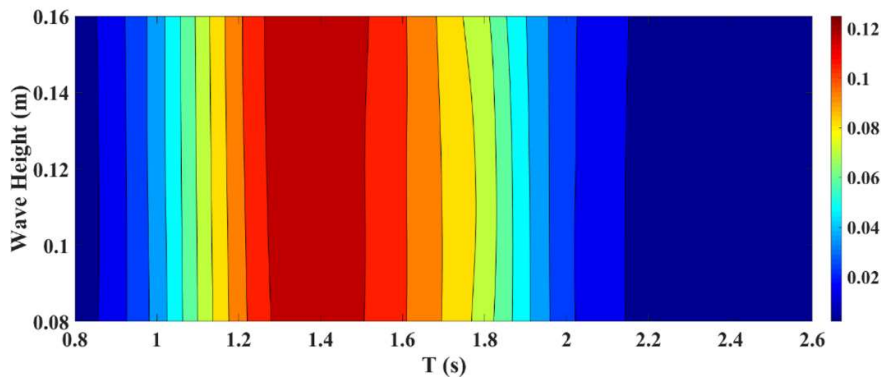
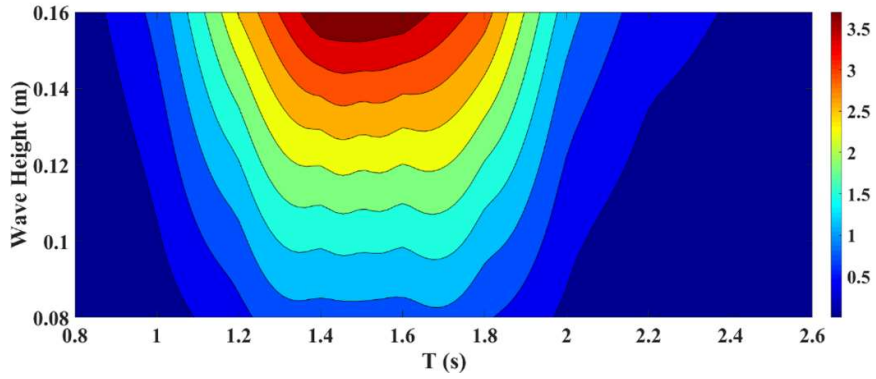


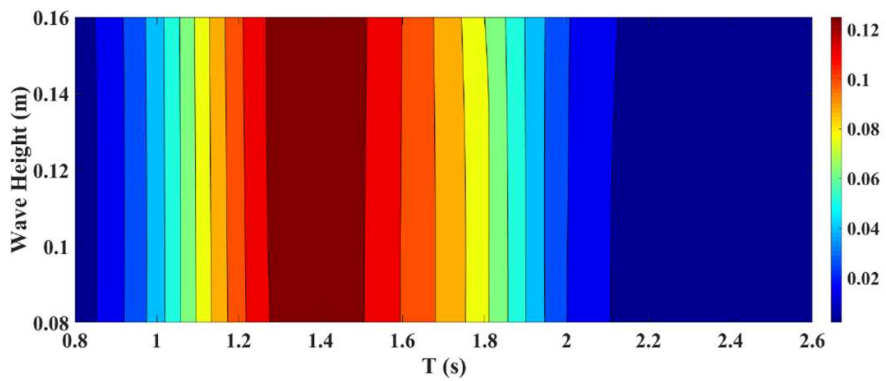
Fig. 18. (b) CWR contours with constant PTO force.



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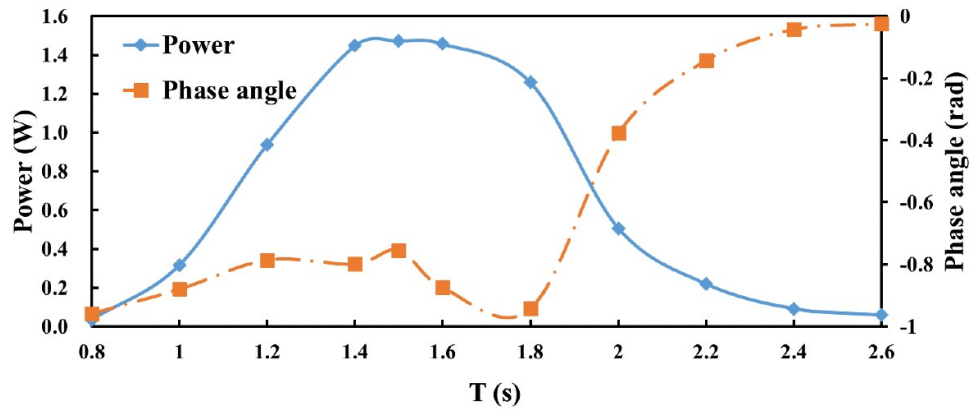
Fig. 18. (c) Capture Power contours with linear PTO force.



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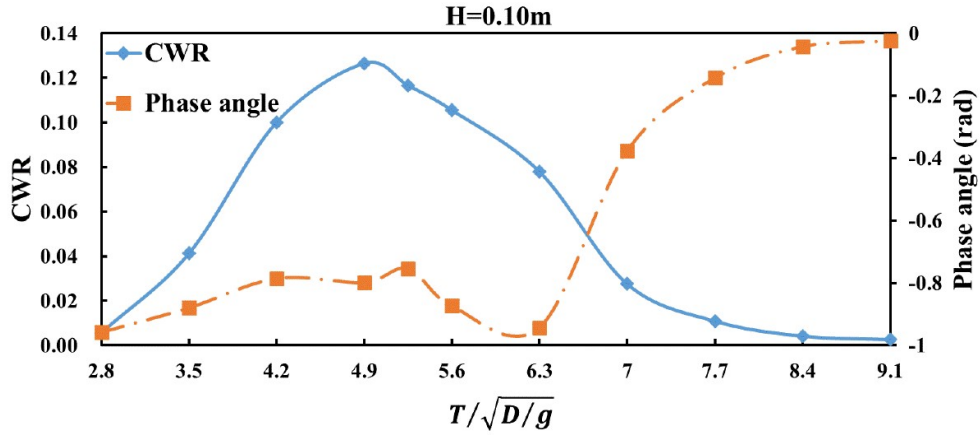
Fig. 18. (d) CWR contours with linear PTO force.



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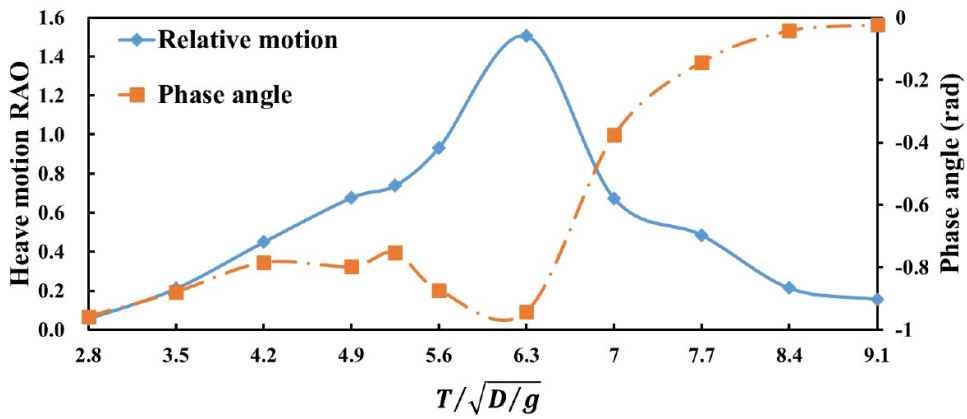
Fig. 19. (a) Power and phase angle with constant PTO force.



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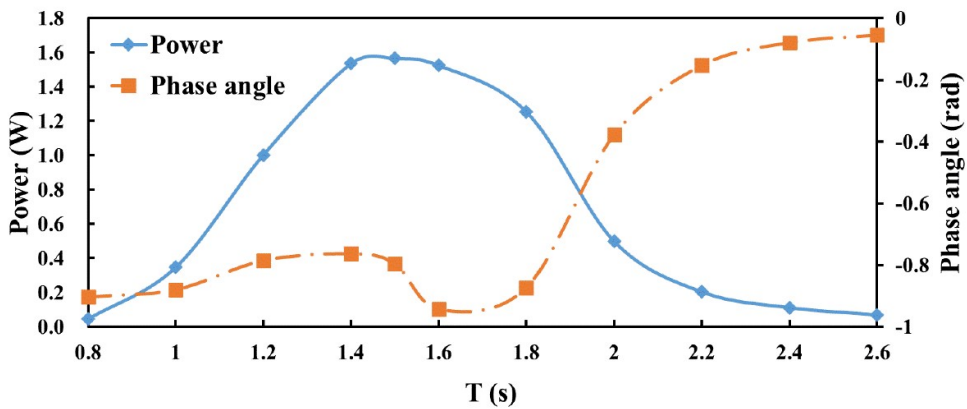
Fig. 19. (b) CWR and phase angle with constant PTO force.



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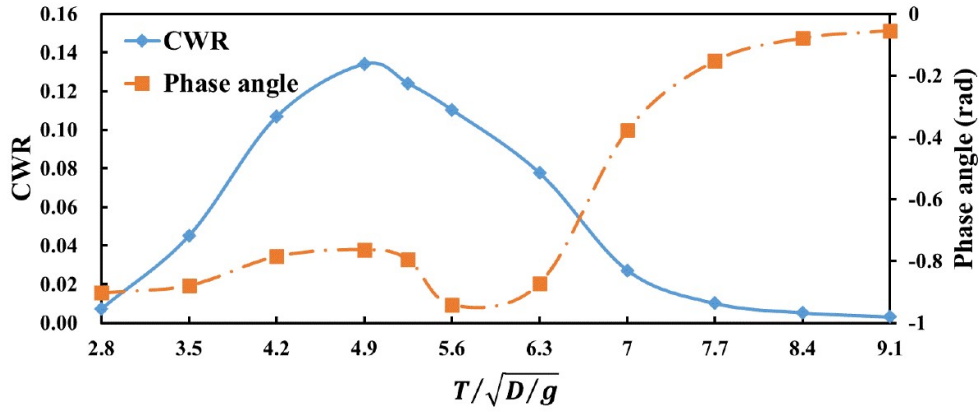
Fig. 19. (c) Relative heave motion RAO and phase angle with constant PTO force.



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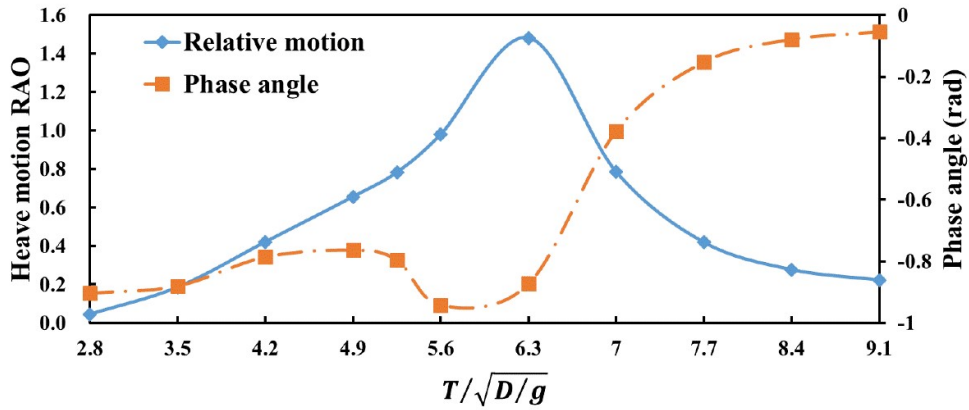
Fig. 19. (d) Power and phase angle with linear PTO force.



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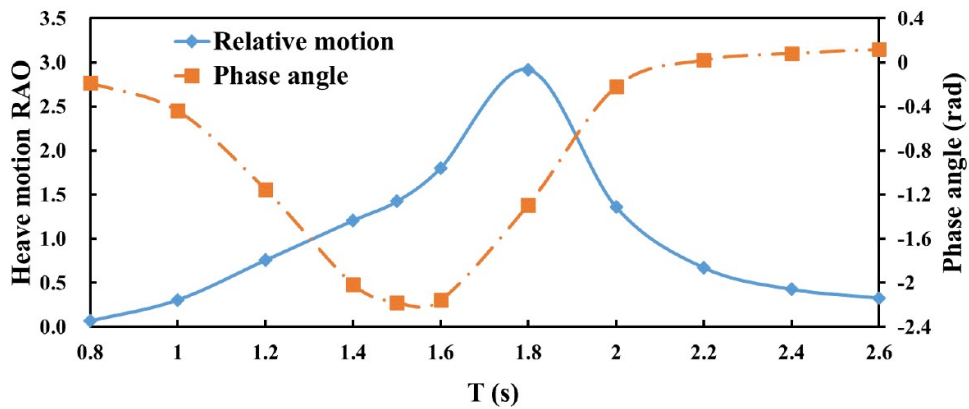
Fig. 19. (e) CWR and phase angle with linear PTO force.



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Fig. 19. (f) Relative heave motion RAO and phase angle with linear PTO force.



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Fig. 19. (g) Relative heave motion RAO and phase angle without PTO force.