

1 **On the importance of tunica intima in the aging aorta: a**
2 **three-layered in silico model for computing wall stresses in**
3 **abdominal aortic aneurysms**

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17 **ABSTRACT**

18 Layer-specific experimental data for human aortic tissue suggest that, in aged arter-
19 ies and arteries with non-atherosclerotic intimal thickening, the innermost layer of
20 the aorta increases significantly its stiffness and thickness, becoming load-bearing.
21 However, there are very few computational studies of abdominal aortic aneurysms
22 (AAAs) that take into account the mechanical contribution of the three layers that
23 comprise the aneurysmal tissue. In this paper, a three-layered finite element model
24 is proposed from the simplest uniaxial stress state to geometrically parametrized
25 models of AAAs with different asymmetry values. Comparisons are made between
26 a three-layered artery wall and a mono-layered intact artery, which represents the
27 complex behavior of the aggregate adventitia-media-intima in a single layer with
28 averaged mechanical properties. Likewise, the response of our idealized geometries
29 is compared with similar experimental and numerical models. Finally, the mechani-
30 cal contributions of adventitia, media and intima are analyzed for the three-layered
31 aneurysms through the evaluation of the mean stress absorption percentage. Re-
32 sults show the relevance and necessity of considering the inclusion of tunica intima
33 in multi-layered models of AAAs for getting accurate results in terms of peak wall
34 stresses and displacements.

35 **KEYWORDS**

36 Arterial wall mechanics ; Abdominal aortic aneurysm; Tunica intima

37 **1. Introduction**

38 An abdominal aortic aneurysm (AAA) is a balloon-like, localized enlargement of the
39 aorta that bulges out beyond the normal diameter of the blood vessel. AAAs affect
40 about 3% of the world population over the age of 50 (LeFevre 2014). Associated risk

41 factors are mostly lifestyle-related (smoking, dyslipidemia, high blood pressure), al-
42 though a heritable component can also play a role. They usually remain asymptomatic
43 until rupture, which can lead to life-threatening internal bleeding with an in-hospital
44 mortality of about 40% and a pre-hospitalization overall mortality of 80% (Kühnl
45 et al. 2017). Repair of an AAA may be done either by open surgery or endovascular
46 aneurysm repair (EVAR). Open repair, as any surgical procedure, may associate with
47 a non-negligible rate of complications such as bleeding during or after surgery, myocar-
48 dial infarction, respiratory impairment or graft infection. On the other hand, EVAR
49 is a minimally invasive technique that only requires small incisions in the groin, but
50 requires a more strict postoperative surveillance over time.

51 The current approach to assess the risk of rupture and to determine whether the
52 patient should undergo surgical repair or not is a dimensional criterion based on the
53 maximum diameter of the lesion (aortic size). If the maximum diameter increases more
54 than 0.5–1 cm in one year, or if it reaches 5.0 cm in women or 5.5 cm in men, surgical
55 repair will be necessary (Lederle et al. 2002; Hans et al. 2005; Grootenboer et al. 2009).
56 Nonetheless, about 13% of AAAs with an aortic size of less than 5 cm rupture, whereas
57 54% of those over 7 cm may not rupture over long periods. Therefore, a more reliable
58 parameter is needed for the assessment of the risk of AAA rupture. Peak wall stresses
59 are suggested by many studies (Rodríguez et al. 2009; Vorp et al. 1998; Raghavan
60 et al. 1996) as a more suitable parameter than the current diameter criterion. However,
61 peak wall stresses cannot be measured in complex geometries just by applying simple
62 analytic techniques, hence, numerical modeling must be used. In this respect, the finite
63 element analysis provides a convenient numerical tool to calculate approximate wall
64 stresses that facilitates the evaluation of the rupture potential of AAAs.

65 From the biomechanical point of view, the aortic wall consist of three layers: adven-
66 titia, the outermost layer; tunica media, which is the medial layer, and tunica intima,
67 which is the innermost layer. In young human arteries and arteries of laboratory an-
68 imals, only the adventitia and media are load-bearing layers and the intima is just a
69 thin layer made up mostly of endothelial cells. However, in aged arteries, the intima
70 attains a significant thickness and the three layers become load-bearing. This is caused
71 by diffuse intimal thickening or intimal hyperplasia, which is considered to be the pre-
72 cursor of atherosclerosis and produces the collagenization of the intima (Movat et al.
73 1958). Some studies explain the thickening as a compensatory response to the wall
74 shear reduction, so that the artery decreases the luminal diameter in response to a
75 reduced blood flow in order to restore shear stress (Glagov and Zarins 1989). Never-
76 theless, there are very few computational studies of AAAs that take into account the
77 mechanical contribution of the intima as an individual layer with its own mechanical
78 properties.

79 Many researchers have extensively studied the layer-specific mechanical properties of
80 human thoracic and abdominal aortas. Weisbecker et al. (2012) tested 14 thoracic aor-
81 tas and 9 abdominal aortas from patients aged between 55 to 77 years with acute non-
82 atherosclerotic intimal thickening, obtaining the material parameters for each layer.
83 Kobielarz et al. (2017) analyzed 27 thoracic aortas from young patients with a mean age
84 of 26 years and early atherosclerotic lesions, concluding that the intima is load-bearing.
85 Amabili et al. (2019) characterized the layer-specific hyperelastic and viscoelastic be-
86 haviour of 12 healthy descending thoracic aortas from patients with an average age of
87 49 years. Akyildiz et al. (2014) studied the mechanical properties of the intimal layer
88 in the presence of atherosclerotic plaques, showing a great dispersion in the tensile
89 and compressive properties of the plaque. Barrett et al. (2019) made a review on the
90 imaging techniques, the experimental tests and the computational methods used to

91 obtain calcified plaque tissue properties, concluding that it is necessary to carry out
92 experimental tests down in scale, towards micron and submicron scales, to understand
93 the calcified plaque mechanical behaviour. In the case of aortic aneurysms, one of
94 the first works to characterize the mechanical properties of each layer was performed
95 by Sokolis et al. (2012), studying layer heterogeneity in 8 ascending thoracic aortic
96 aneurysms from patients aged between 60 to 80 years. Sassani et al. (2015) determined
97 layer dependent tissue properties in abdominal aortic aneurysms from 15 patients aged
98 between 58 and 85 years. Deveja et al. (2018) analyzed the mechanical properties of
99 each layer in thoracic aortic aneurysms and non-aneurysmal aortas from 17 patients.

100 Even from the development of multi-layer constitutive relations for arterial walls by
101 Gasser et al. (2006), and the obtaining of the layer-specific material parameters by
102 several researchers, the intima has been excluded from numerical studies due to its
103 small thickness in young arteries (Alastrué et al. 2007). Prior studies have performed
104 isotropic finite element simulations considering elastic or hyperelastic constitutive laws
105 in mono-layered arterial walls, like Scotti et al. (2005) and Raghavan and Vorp (2000).
106 Other authors carried out more advanced computational models implementing the
107 anisotropy of the arterial wall in patient-specific geometries like Xenos et al. (2010) for
108 a mono-layered AAA wall, Rodríguez et al. (2008) for different idealized mono-layered
109 AAAs, or Alastrué et al. (2007) for a two-layered iliac artery, in which only adventitia
110 and media were taken into account. Further research on hyperelastic constitutive laws
111 also include the numerical implementation of residual stresses, like Ahamed et al.
112 (2016) for evaluating wall stresses using mono-layered patient-specific geometries, or
113 Labrosse et al. (2013), where residual stresses are obtained by experimental testing
114 on pressurized ascending, thoracic and abdominal cylindrical samples. However, none
115 of them consider the increase in stiffening and thickness of the innermost layer of the
116 aorta.

117 Regarding experimental studies on residual stresses and their spatial distribution
118 in the aortic wall, the next contributions could be highlighted. Sokolis et al. (2017)
119 performed a detailed experimental identification of the spatial distribution of circum-
120 ferential residual strains in human aorta considering age and gender; they also studied
121 the regional and interlayer distribution of residual deformations and opening angles
122 in porcine aortas in (Sokolis 2019). Amabili et al. (2019) measured opening angles
123 in ascending human aortas to identify the circumferential residual stresses and axial
124 stretches. Finally, among the few studies of residual strains in aneurysms, the work
125 by Sokolis (2015) on ascending thoracic arteries, considering the variation of residual
126 deformations in the different aortic layers; and the work by Sassani et al. (2015) ob-
127 taining layer dependent residual stretch measurements in abdominal aortic aneurysms,
128 must be featured.

129 The first three-layered models assumed an isotropic linear elastic response for all the
130 layers, like Gao et al. (2006, 2008) for three-layered aneurysmal and non-aneurysmal
131 aortic arches, where the Young's modulus of the medial layer was assumed to be three
132 times larger than that of the intimal and adventitial layer. Gao et al. also performed
133 FSI analyses on two dimensional (2D) axisymmetric geometric models of stented three-
134 layered aneurysms (Gao et al. 2013). Simsek and Kwon (2015) and Gholipour et al.
135 (2018) evaluated the rupture potential of three-layered idealized aneurysmal and non-
136 aneurysmal geometries assuming different hyperelastic isotropic material properties for
137 each layer. Recent studies analyzed the inclusion of residual stresses in three-layered
138 aneurysms, like Pierce et al. (2015) for a patient-specific geometry. Other researchers
139 like Strbac et al. (2017) even studied how to improve the finite element codes for com-
140 puting faster, and more accurate solutions in three-layered patient-specific geometries.

141 Nonetheless, the structural role played by tunica intima during the development of
142 atherosclerosis, and its through-the-thickness stress distribution has not been clarified
143 yet.

144 This work proposes a three-layered model that allows to study the influence of inti-
145 mal thickening from a mechanical point of view on different parametrized geometrical
146 models of AAAs. The calibration of the material model, which is considered hypere-
147 lastic anisotropic, is done through finite element simulations of uniaxial tests of aorta
148 strips cut in the circumferential and axial direction, and the inflation of plane strain
149 aorta rings subjected to systolic blood pressure. Then, peak wall stresses and dis-
150 placements are computed in three different idealized AAA geometries considering a
151 three-layered wall, in which each layer is modeled separately in a continuum mesh
152 using different material parameters, and an intact monolayered human aorta wall. As
153 loading conditions, we apply a static internal pressure of 16 kPa (120 mmHg) to simu-
154 late the luminal pressure at the end-systolic state. No residual stresses are considered.
155 Additionally, comparisons between the three-layered and the intact wall are made,
156 as well as between different material models (elastic and hyperelastic isotropic) from
157 other studies. Finally the stiffness of each layer that make up the aneurysmal tissue is
158 evaluated and compared through its mean stress absorption percentage.

159 2. Methods

160 2.1. Constitutive behavior of arterial tissue

161 Constitutive modeling of arterial tissue has undergone a significative evolution over the
162 past decade. Early-modeled aneurysmal tissue was characterized as a single layer linear
163 elastic material (Martino et al. 2001; Li and Kleinstreuer 2007; Georgakarakos et al.
164 2010; Wang and Li 2011). As a consequence of the uniaxial testing of aortic tissue
165 specimens carried out by Raghavan et al. (2000), the nonlinear elastic behaviour is
166 incorporated in material models, where the mechanical behavior of the arterial wall was,
167 for the first time, modeled as hyperelastic, with a constitutive law based on a simplified
168 criterion derived from the Mooney-Rivlin strain energy function. Thereafter, the vast
169 majority of the computational studies of fully developed aneurysms assumed isotropy
170 (Wang et al. 2002; Chandra et al. 2013; Li et al. 2008; Maier et al. 2010). A high degree
171 of anisotropy was subsequently noticed by Geest et al. (2006) after performing biaxial
172 testing to characterize the mechanical properties of aortic tissue in the longitudinal and
173 circumferential direction. Then, the obtained experimental data would be fitted to a
174 four parameter exponential strain function proposed by Vito and Hickey (1980). Later
175 on, the understanding of the arterial histology by means of extensive experimental data
176 has led to new and more accurate constitutive models that make it possible to analyze
177 the multi-layered nature of the arterial wall as an anisotropic fiber-reinforced material
178 (Holzapfel et al. 2000). The aforementioned continuum approach was considered in
179 this study by means of the constitutive model developed by Holzapfel et al. (2000) and
180 Gasser et al. (2006). This model asserts that each artery layer may be understood as a
181 composite reinforced material constituted by two families of collagen fibers embedded
182 in a soft incompressible matrix, which is mostly made up of elastin. The collagen fibers
183 are arranged in spirals and symmetrically oriented with respect to the circumferential
184 direction. The strain energy function used to model each layer of the artery wall is

185 given by

$$\Psi = \Psi_{\text{iso}} + \Psi_{\text{aniso}}. \quad (1)$$

186 Ψ can be divided in an isotropic part, Ψ_{iso} , which represents the energy stored in the
 187 non-collagenous soft matrix, and anisotropic part, Ψ_{aniso} , which provides the energy
 188 stored in the collagen fibers

$$\Psi_{\text{iso}} = \frac{\mu}{2} (\bar{I}_1 - 3) \quad (2)$$

189

$$\Psi_{\text{aniso}} = \frac{k_1}{2k_2} \sum_{i=1}^{N=2} [\exp(k_2 \bar{E}_i^2) - 1]. \quad (3)$$

190 \bar{E}_i , which stands for the Green-Lagrange strain-like quantity, can be expressed as

$$\bar{E}_i = \kappa \bar{I}_1 + (1 - 3\kappa) (\bar{I}_{4i} - 1) \quad (4)$$

191 where

$$\bar{I}_{4i} = \mathbf{a}_{0i} \otimes \mathbf{a}_{0i} : \bar{\mathbf{C}}. \quad (5)$$

192 The non-collagenous soft matrix is modeled as an incompressible isotropic neo-
 193 Hookean material, with $\mu > 0$ as the shear modulus in the undeformed configuration,
 194 and \bar{I}_1 as the first strain invariant of a modified right Cauchy-Green tensor, $\bar{\mathbf{C}} = \bar{\mathbf{F}}^T \bar{\mathbf{F}}$.
 195 $\bar{\mathbf{F}}$ represents the isochoric part of the deformation gradient and comes from a multi-
 196 plicative decomposition of the deformation gradient $\mathbf{F} = \left(J^{\frac{1}{3}} \mathbf{I}\right) \bar{\mathbf{F}}$, where $J^{\frac{1}{3}}$ and $\bar{\mathbf{F}}$
 197 represent the volumetric and isochoric part of the deformation gradient, respectively,
 198 and \mathbf{I} is a second-order unit tensor. In equation (3) the strain energy stored in the
 199 collagen fibers is defined as an exponential function, where N is the number of fiber
 200 families of each layer. In accordance with Schriefel et al. (2012), a two-fiber family is
 201 considered for all the layers in this study. $k_1 > 0$ is a stress-like parameter, while $k_2 > 0$
 202 is a dimensionless parameter, and both are determined from mechanical tests of the
 203 tissue. \bar{E}_i represents the strain in the direction defined by the mean orientation of each
 204 fiber family, which is in turn denoted by the vector \mathbf{a}_{0i} . The parameter $\kappa \in [0, 1/3]$
 205 is also unitless and describes the level of dispersion of the fiber directions. According
 206 to the value of κ , collagen fibers may be perfectly aligned ($\kappa = 0$), which means that
 207 there is no dispersion, or randomly distributed ($\kappa = 1/3$), which corresponds with a
 208 spherical distribution of the density function and the material becomes isotropic. κ and
 209 \mathbf{a}_{0i} are determined from histological data. Finally, \bar{I}_{4i} is the pseudo-invariant of $\bar{\mathbf{C}}$.

210 The material model presented above was used in all our simulations. It is based
 211 on experimental tests and histological analysis performed on non-aneurysmal aortas
 212 considering two families of collagen fibers. Other authors like Gasser et al. (2012) and
 213 Sassani et al. (2015), have made improvements to this model by identifying the spatial
 214 organization of the collagen fiber network. In these works, the spatial distribution of
 215 collagen in each layer of the aortic wall, which determines its strength and stiffness,
 216 is reproduced more precisely using three families of collagen fibers instead of two, one

217 circumferential and two diagonals, allowing a better characterization of the risk of
218 aortic rupture.

219 *2.2. Finite element modeling*

220 The simulations presented in this article were conducted by using the FEM commercial
221 software Abaqus/Standard 6.14, in which the constitutive model explained previously
222 is built-in (201 2014b). To check the viability of the proposed three-layered model,
223 different geometries were considered, and in all of them we used the thicknesses and
224 average properties of each layer obtained experimentally by Weisbecker et al. (2012)
225 for the abdominal aorta. From the simplest to the most complex one, we developed
226 finite element models of uniaxial tests performed on rectangular aorta strips cut in the
227 axial and circumferential direction, human aorta plane strain rings, and finally three
228 different parametrized geometric aneurysms with intimal thickening. Residual stresses
229 are not included in the finite element models developed.

230 In the three types of finite element models developed, experimental mechanical prop-
231 erties of non-aneurysmal abdominal aortic tissues are used, so that all models share
232 the same properties and the results could be comparable.

233 **3. Finite element models of uniaxial test of aorta strips**

234 Based on the work developed in Gasser et al. (2006), finite element computations of
235 uniaxial tension tests were performed on rectangular intact and layer-separated aorta
236 strips with non-atherosclerotic intimal thickening cut in the axial and circumferential
237 direction. The specimens are loaded in the longitudinal direction with a force F that
238 produces an elongation δ and are assumed to be stress free in the undeformed con-
239 figuration. The definition of axial and circumferential specimens as well as the model
240 configuration are illustrated schematically in Figure 1.

241 The referential dimensions of the strips were 20 mm for the length, 6 mm for the
242 width and a total thickness of 2.69 mm. In the three-layered separated strips, the
243 thickness of the different layers that comprise the tissue are: 0.68 mm for the intima,
244 0.94 mm for the media and 1.07 mm for the adventitia. The in-plane dimensions are
245 based on the ones provided in the Abaqus Benchmarks Guide (201 2014a), where a
246 rectangular 10 x 3 x 0.5 mm adventitial strip is analyzed under uniaxial tension. Since
247 we had a maximum thickness of 1.07 mm and we wanted to keep a similar in-plane
248 aspect ratio not to distort the results, the final dimensions of our strips had to be bigger
249 by a factor of two compared to the benchmark model. Regarding the thicknesses, they
250 are in accordance with the median thicknesses of the intima, media and adventitia
251 determined by Weisbecker et al. (2012). The different material constants as well as
252 the orientations of the two families of fibers considered for the layer-separated and the
253 intact artery wall, which considers the aggregate adventitia-media-intima in a single
254 layer with averaged mechanical properties, are summarized in Table 1.

255 Exploiting the symmetry of the problem, only one half of the geometry was modeled.
256 To model the incompressible deformation of the arterial tissue with sufficient precision,
257 a total of 12,000 eight node linear solid hybrid elements (C3D8H) were used for the
258 adventitia, 12,100 for the media and 36,000 for the intima, with a minimum of three
259 elements through-the-thickness, whereas 60,000 elements were required for the intact
260 layer models. Regarding the type of element used, it is important to consider the

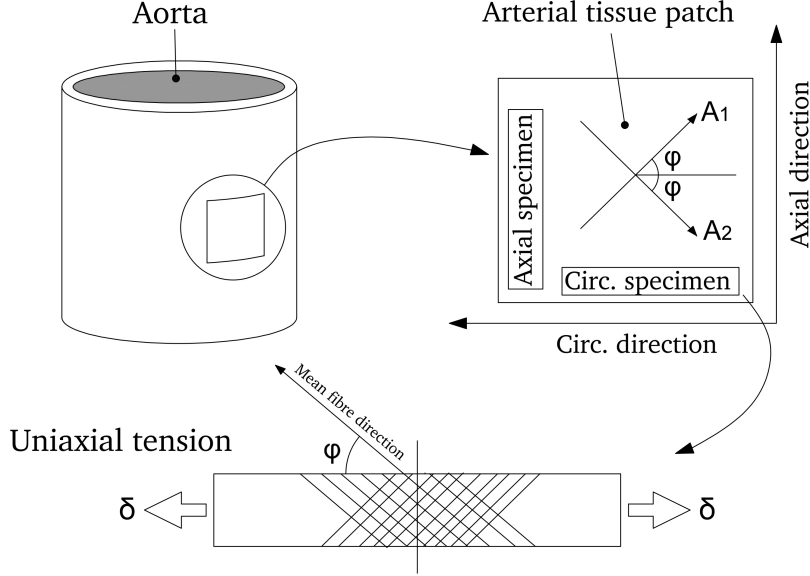


Figure 1. Definition of axial and circumferential specimens and uniaxial tension test configuration (adapted from Gasser et al. (2006)). A_1 and A_2 represent the mean direction of each family of fibers.

Table 1. Constitutive parameters for the layer-separated specimens and the intact (three-layer composite) wall of the human abdominal aorta (taken from Weisbecker et al. (2012)).

Layer	μ (MPa)	k_1 (MPa)	k_2 (-)	φ ($^\circ$)	κ (-)
Intima	0.044	10.14	0.00	40.5	0.25
Media	0.028	0.81	12.42	39.1	0.18
Adventitia	0.010	0.38	3.35	40.59	0.11
Intact wall	0.019	5.15	8.64	38.8	0.24

261 fact that, the bulk modulus of an incompressible material is much greater than its
 262 shear modulus. Due to this, a displacement-based element is not suitable since a pure
 263 hydrostatic stress state would not produce changes in the displacement field. Therefore,
 264 a mixed formulation, using not only displacement but stress variables, is required to
 265 solve the equilibrium equations. For that purpose, hybrid elements are used in our
 266 simulations to model the incompressible behavior of soft tissue, which is a realistic
 267 assumption since it is mostly made up of water. Values of the Cauchy stresses and
 268 strains were computed for each integration point in the tensile direction. The results
 269 are compared in terms of stress vs strain curves with the experimental results obtained
 270 by Weisbecker et al. (2012).

271 *3.1. FEM analysis results for the uniaxial tests of aorta strips*

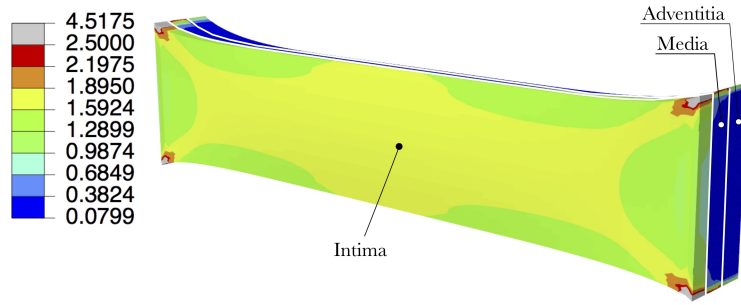
272 Figure 2 shows the computed Cauchy stress in the tensile direction for the circumfer-
 273 ential and axial three-layered patches at a total displacement of 2.5 mm. In agreement
 274 with the results obtained by Gasser et al. in (Gasser et al. 2006) no significant change
 275 is observed in the thickness of the specimens, while the width decreases in the mid-
 276 dle part of the strips due to the incompressibility constraint. Despite the similarity
 277 of the transition zones at the end of strips, all the specimens show a stiffer response

278 in the circumferential direction. Tunica intima exhibits the maximum stress values of
 279 the layer-separated specimens, reaching stresses of 2.5 MPa when it is cut in the cir-
 280 cumferential direction. The adventitial and medial strip, with maximum values of 1
 281 and 1.2 MPa respectively, present a softer behavior than the intimal strip. One of the
 282 main reasons for this is the degree of dispersion of the collagen fibers, κ , which is much
 283 higher in the intima ($\kappa = 0.25$) than in any other layer. κ controls the start of the
 284 stiffening effect produced by the alignment of the collagen fibers in the direction of the
 285 applied load, therefore, higher values of κ provide a stiffer response at equal stretches.
 286 Concerning the intact and the three-layered patch, we observe a parallel structural
 287 response for the first one compared to the above analyzed separated intima layer: the
 288 high dispersion of the collagen fibers for this case, $\kappa = 0.24$, which is in fact very similar
 289 to the value of the intima layer, $\kappa = 0.25$, leads to a macroscopic stiffer behaviour,
 290 where fibers do not need to rotate before carrying load and just a small reduction of the
 291 width of the specimen is noticed. On the other hand, the three-layered patch shows a
 292 dissimilar mechanical behaviour in which we observe noteworthy stress discontinuities
 293 between the layers where the intima is absorbing the largest amount of stress.

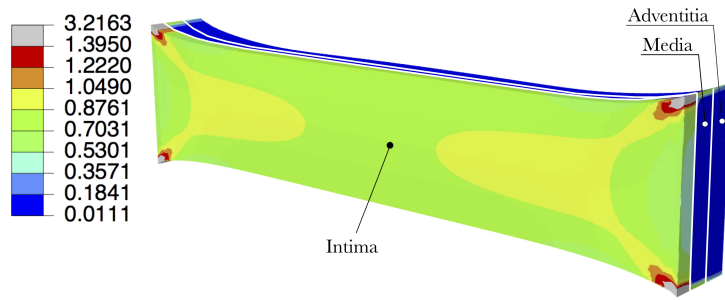
294 Figure 3 shows the stress versus stretch response in the direction of the applied load
 295 for the circumferential and axial specimens. The Cauchy stress was computed as $\sigma =$
 296 $F\lambda/(TW)$, where F stands for the applied force, T for the thickness of the specimen,
 297 W for the width (both in the undeformed configuration), and $\lambda = l/L$ represents the
 298 stretch in the loading direction, where l and L are the lengths of the specimen in
 299 the deformed and reference configuration, respectively. The qualitative stress-stretch
 300 response of the three-layered patch is similar to the one reported by Holzapfel et al.
 301 (2005) for coronary arteries, and Weisbecker et al. (2012) for the abdominal aorta. As
 302 it can be seen, the intima manifests an early exponential stiffening at low stretches in
 303 both circumferential and axial directions. This stress-stretch response is closely related
 304 to the high degree of dispersion in the collagen fiber directions previously commented,
 305 which is in turn, associated with the collagenization of the innermost hyperelastic
 306 layer during the development of the diffuse intimal thickening of the aorta (Movat et al.
 307 1958). Media and adventitia curves show a softer behavior in both directions, where the
 308 exponential stiffening produced by the anisotropic contribution of the collagen fibers
 309 to the strain energy function is delayed in comparison to the intimal layer. For a total
 310 Cauchy stress of 0.7 MPa, the axial specimen of the intimal layer reaches a maximum
 311 stretch of 1.24, while adventitial and medial strips have maximum stretches of 2.5 and
 312 1.5, respectively. As for the three-layered tissue, despite being made up of intima, media
 313 and adventitia, its mechanical reaction is somewhat less stiff in comparison with the
 314 intima and the intact wall, probably due to the loss of strain energy produced during
 315 the discontinuous stress migration from tunica intima to the other two layers.

316 4. Human aorta plane strain rings

317 Before assessing the effects of the intimal thickening in an AAA geometry, a simpler case
 318 is studied. To test the feasibility of the three-layered model proposed, a human aorta
 319 plane geometry was modeled with the configuration shown in figure 4. The dimensions
 320 of the rings were 10 mm for inner radius and a different thickness depending on the
 321 layer modeled. For the layer-separated rings the thickness of each layer is the same
 322 as the ones used previously for the uniaxial test simulations: 0.68, 0.94 and 1.07 mm
 323 for intima, media and adventitia, respectively. Plane strain boundary conditions were
 324 applied for all the models.



(a) Three-layered patch cut in the circumferential direction.



(b) Three-layered patch cut in the axial direction.

Figure 2. Finite element computations of the Cauchy stress in the tensile direction at a displacement of 2.5 mm (MPa). The grey zones are a result of edge effects caused by the stress concentrations due to the displacement constraint applied on the lateral face of the specimen.

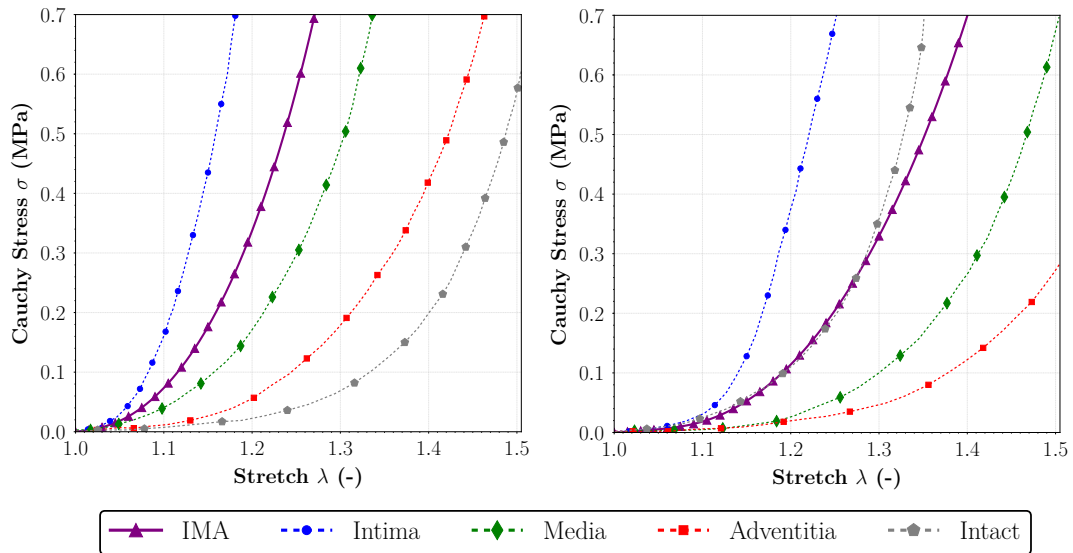


Figure 3. Computed Cauchy stress vs stretch curves of the circumferential (left) and axial specimens (right) in solid purple curve for the IMA (intima-media-adventitia) three-layered tissue patch and the intact tissue patch. Experimental results by Weisbecker et al. (2012) in dashed curves for intima, media and adventitia tissues.

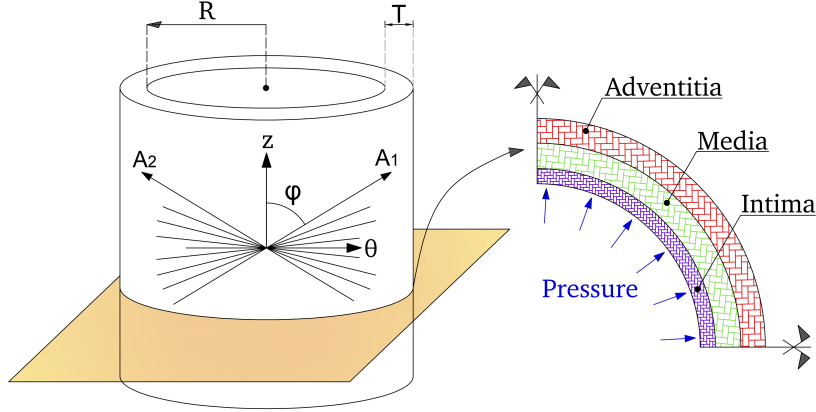


Figure 4. Scheme of the model configuration for the human aorta plane strain rings (adapted from Gasser et al. (2006)).

325 In order to simulate the end-systolic state, in which the artery undergoes the largest
 326 wall stresses, an internal pressure of 16 kPa (120 mmHg) was implemented. Material
 327 constants for the layer-separated and intact three-layer composite artery wall are col-
 328 lected in Table 1. Regarding the layer-separated models, 1,215 eight-node solid hybrid
 329 elements (C3D8H) were used for intima, 486 elements for the media and 748 elements
 330 for the adventitia, with a minimum of two elements through-the-thickness of each
 331 layer. In the intact aorta rings a total of 2500 elements were required to obtain pre-
 332 cise through-the-thickness stress distributions. Circumferential stresses and stretches were
 333 computed at each integration point across the thickness of the artery wall.

334 *4.1. Results for aorta plane strain rings*

335 The computed circumferential stresses produced by an internal pressure of 16 kPa
 336 (120 mmHg) are depicted in Figure 5. The absence of residual stresses leads to a
 337 pure tension state through the whole thickness in both the layer-separated and three-
 338 layered configurations. Regarding the layer-separated rings, we observe maximal values
 339 at the inner radius of the adventitial ring of about 0.32 MPa, which decrease to 0.24
 340 MPa at the outer, while the stress distributions of the media and intima are quite
 341 similar at the inner radius, reaching values close to 0.32 MPa, but differ from the outer
 342 radius, where the intima shows slightly higher circumferential stresses that go up to
 343 0.28 MPa. Furthermore, we notice big differences between the through-the-thickness
 344 circumferential stresses of the intact artery, the two-layered and the three-layered rings,
 345 which are depicted in Figure 6. As it can be seen, the intact artery shows an analogous
 346 non-linear stress distribution to the one observed previously for the layer-separated
 347 cylinders, with a range of stress values that goes from 0.017 MPa at the inner surface
 348 to 0.012 at the outer. On the other hand, in agreement with Alastrué et al. (2007), we
 349 observe again "the discontinuities caused by the heterogeneity of the two-layered and
 350 the three-layered wall". As shown in Figure 6, the existing stress value of 0.224 MPa
 351 at the inner part of the intima suddenly drops at the interface with the media, where
 352 it took a value of 0.04 MPa. In the same way, another stress jump is found at the
 353 interface between media and adventitia, but this time not as important as the previous
 354 one, dropping from 0.04 to 0.02 MPa.

355 The results in terms of internal pressure versus circumferential stretch (p_i/λ_θ) are

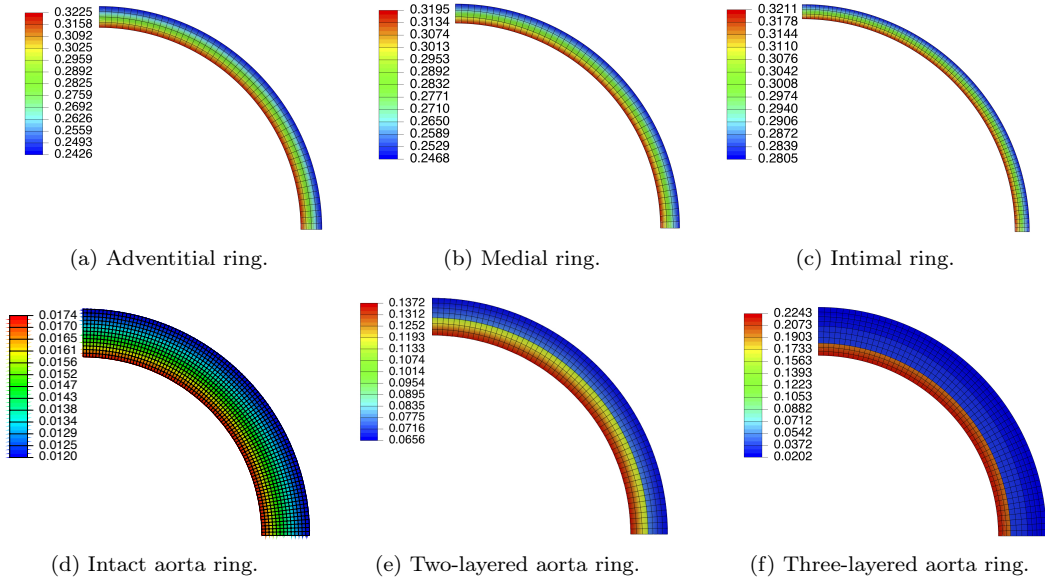


Figure 5. Circumferential stress distributions in the aorta plane strain rings at an internal pressure of 16 kPa. The magnitude of the stresses is given in MPa.

356 illustrated in Figure 7. Once more we can see how the internal pressure/circumferential
 357 stretch response tends to stiffen with increasing κ . As we saw in the uniaxial tests, with
 358 an early exponential stiffening, tunica intima is acting again as the stiffest layer, giving
 359 a total circumferential stretch of 1.11 at an internal pressure of 16 kPa, while the
 360 adventitial layer is the softest with a final stretch of 1.34 for the same internal pressure.
 361 The medial layer shows a delayed structural response that is between the intima and
 362 the adventitia, reaching stretch values of 1.25. The three-layered and the intact rings
 363 have a similar pressure/stretch behavior, even if the former one is much stiffer despite
 364 the large stress discontinuities at the interfaces between the layers which produce a
 365 decrease of the stored strain energy in the collagen fibers.

366 5. Parametrized idealized geometrical models of AAAs

367 Once the effects of intimal thickening have been assessed in simpler geometries, a more
 368 realistic shape is needed to take into account the influence of the typical geometri-
 369 cal non-linearity that characterises fusiform aneurysms, which are the most common
 370 ones. For this purpose, an in-house code (Díaz 2016) was developed. This code uses
 371 the application program interface of the open source CAD/CAE package SALOME
 372 (201 2015) to create the digital model of three-dimensional extruded solid geometries.
 373 The code considers all the geometric and physical variables that characterise an ide-
 374 alized aneurysm, such as length, azimuthal asymmetry, wall thickness, the undilated
 375 diameter at the inlet/outlet sections and the maximum diameter at the midsection of
 376 the AAA. The circular cross sections have the ability to rotate around the three axis,
 377 and the geometries are different in terms of wall heterogeneity and asymmetry, which
 378 are depicted by cross sections perpendicular to the z -axis, hence coinciding with its
 379 centerline. The asymmetry is given by β and is defined as $\beta = r/R$ and schematically
 380 illustrated in Figure 8 as originally proposed by Vorp et al. (1998), where r and R
 381 are the radius measured at the midsection of the AAA cavity from the longitudinal

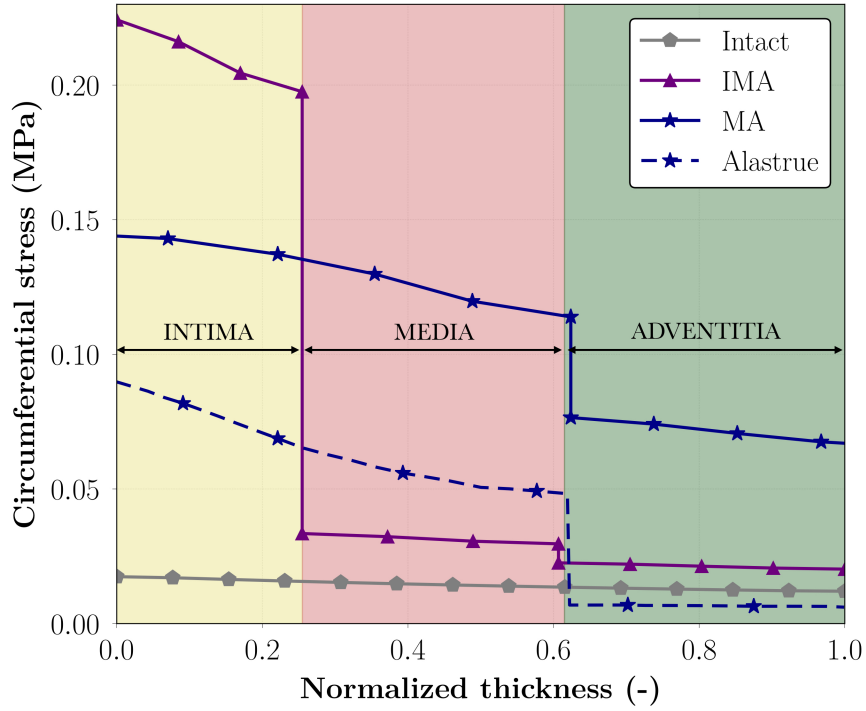


Figure 6. Through-the-thickness circumferential stresses of the three-layered (IMA), two-layered (MA) and intact artery rings. Comparison with the results obtained by Alastrué et al. (2007) for a two-layered human iliac artery plane strain ring when an internal pressure of 16 kPa is applied without residual stresses (MPa).

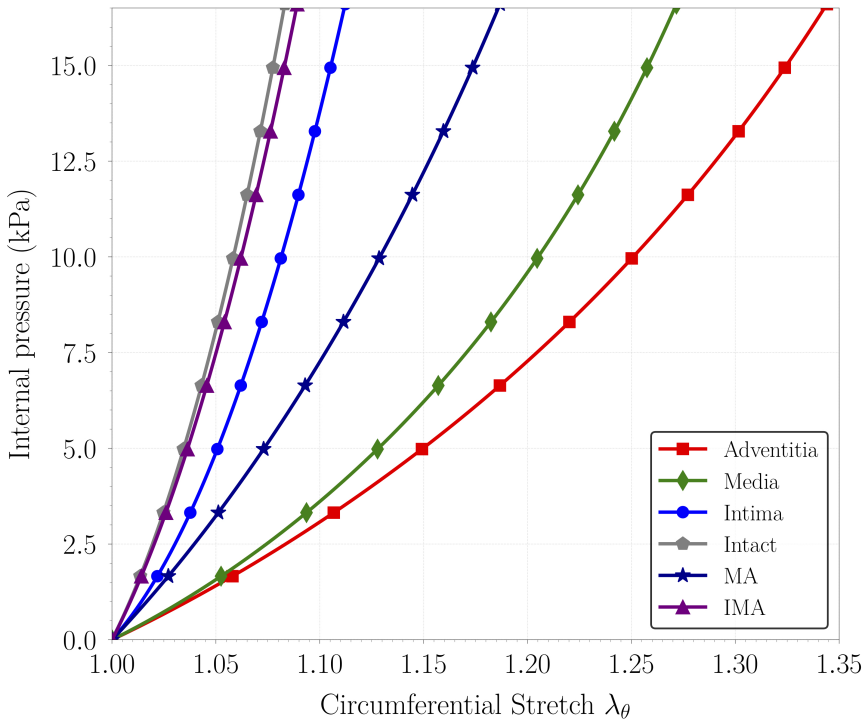


Figure 7. Computed internal pressure versus circumferential stretch of the aorta rings at an internal pressure of 16 kPa.

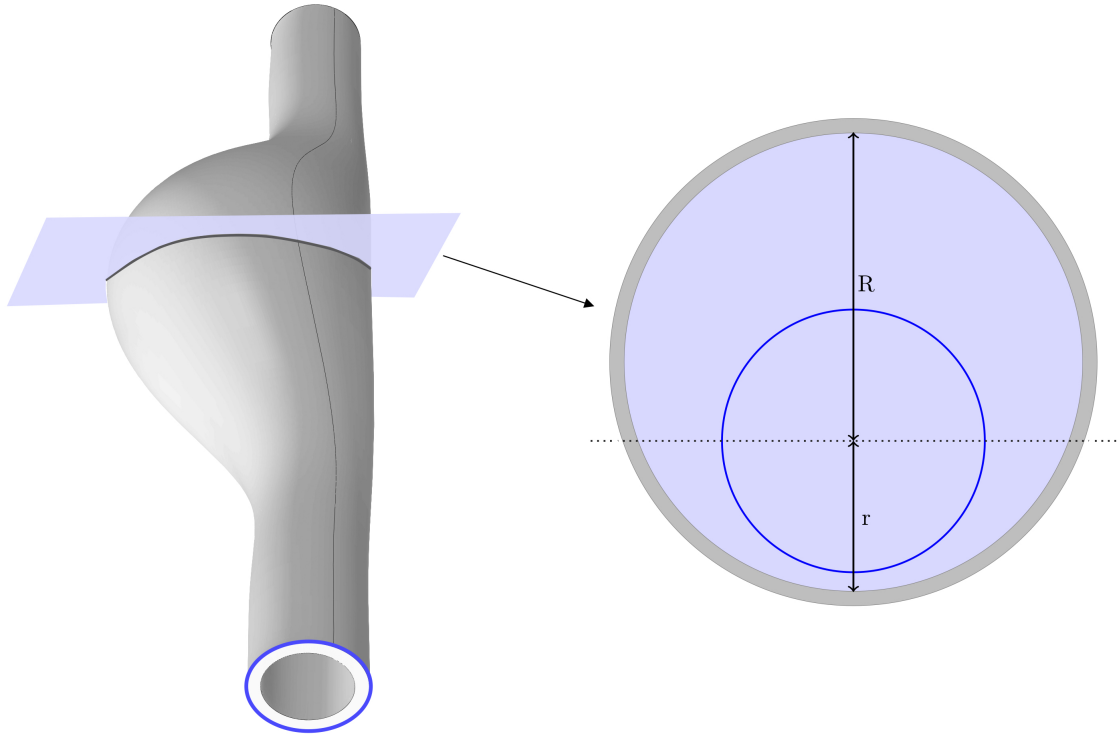


Figure 8. Graphical description of the azimuthal asymmetry.

382 z-axis to the posterior and anterior walls, respectively. An aneurysm for which only the
 383 anterior wall is dilated whereas the posterior wall is approximately flat, corresponds to
 384 a value of $\beta = 0.2$. A value of $\beta = 1.0$ corresponds to azimuthal symmetry.

385 Following the aforementioned procedure, three different geometries of AAA models
 386 with a total length of 23 cm were generated, varying the value of the asymmetry
 387 parameter between $\beta = 1.0$ (azimuthal symmetry) and $\beta = 0.2$ (only the anterior
 388 wall is dilated), with a medium value of $\beta = 0.6$. A value of $d = 2$ cm was adopted
 389 for the undilated diameter at the inlet and outlet sections, and a maximum diameter
 390 of 6 cm was considered at the midsection of the AAA sac. The common value used
 391 from a clinical outlook to recommend surgical repair or endovascular intervention is
 392 AAA transverse diameter between 5 and 6 cm (Galland et al. 1998). Consequently,
 393 a maximum diameter of 6 cm was chosen for this study, since it is comparable to
 394 the largest transverse dimension for assessment of rupture potential. Considering that
 395 this is not a patient-specific study, the uniform wall thickness assumption seems to be
 396 reasonable. In this manner, a total constant wall thickness of 2.69 mm has been adopted
 397 in all the geometries. For the layer-separated models, the thicknesses for intima, media
 398 and adventitia remain the same as the ones considered previously (0.68, 0.94 and 1.07
 399 mm). The resultant geometries are depicted in Figure 9.

400 The effect of the luminal pressure at the end-systolic state was simulated once again
 401 by applying an internal static pressure of 16 kPa (120 mmHg) on the inner surface of
 402 the aneurysm. We do not use a dynamic loading since no change can be observed for the
 403 hyperelastic model with respect to simple static loading, if neither viscoelastic model
 404 nor Fluid Structure Interaction analysis is considered. Residual stresses have been
 405 generally neglected in this study. These simplification is typically used when dealing

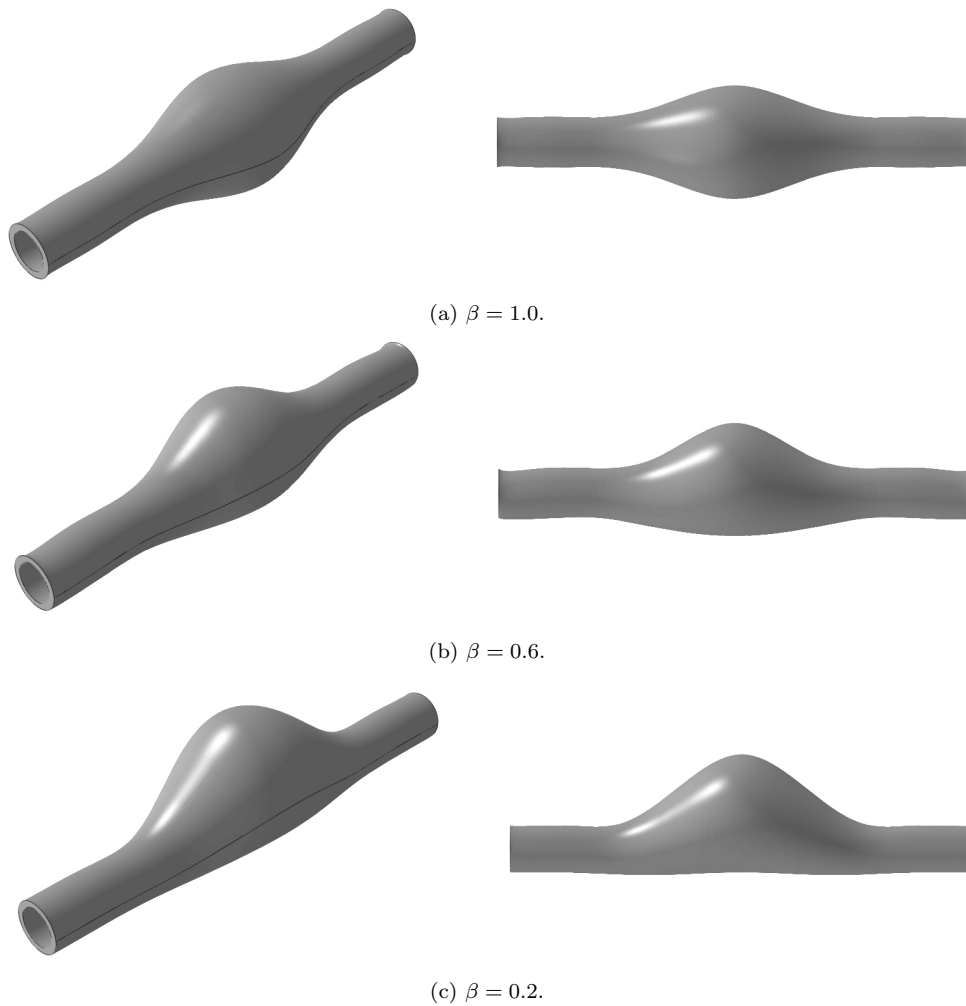


Figure 9. Idealized geometries of the AAA models considered in the study.

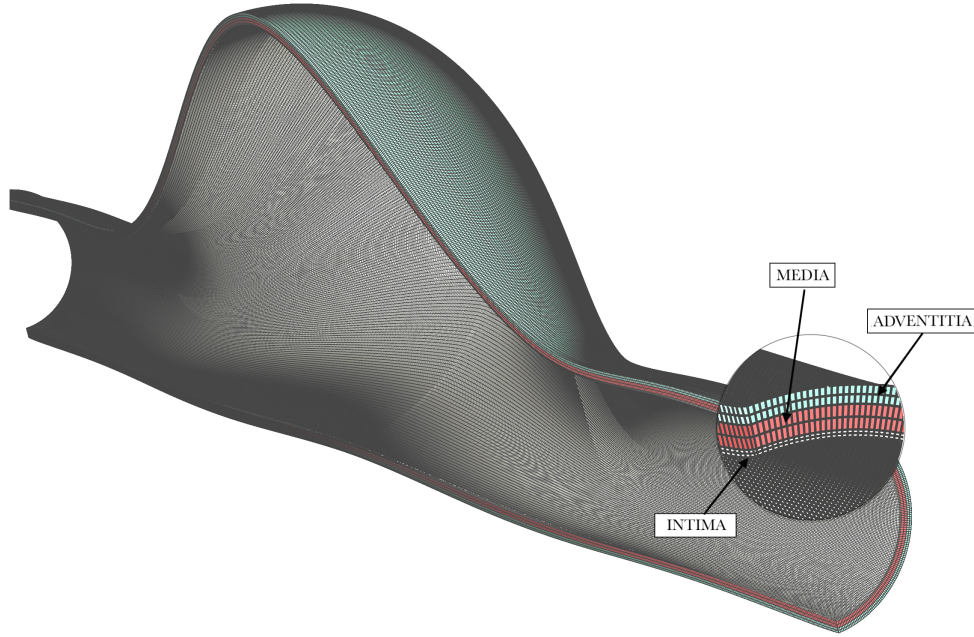


Figure 10. Sagittal view of the typical mesh used for the simulations of the three-layered aneurysms. This geometry corresponds to $\beta = 0.2$.

406 with complicated 3D AAA geometries (Li et al. 2010; Gee et al. 2010; Humphrey
 407 and Holzapfel 2012; Strbac et al. 2017). In our case, given that we are considering
 408 a multilayered wall, it would be even more complicated to quantify the value of the
 409 residual stresses for each layer. Therefore, the implementation of residual stresses is
 410 out of the scope of this paper and it will be prepared for a forthcoming publication.
 411 Nevertheless, our study is relevant in elucidating the limits and uncertainties introduced
 412 by this assumption.

413 Applying proper boundary conditions referred to a cylindrical coordinate system,
 414 the constraining effect caused by the iliac and renal arteries was simulated by im-
 415 posing zero longitudinal displacement at both ends of the undilated sections (Vorp
 416 et al. 1998). Even though this type of boundary conditions smooths the numerical re-
 417 sponse (Rodríguez et al. 2008), the length of the AAA must be enough not to produce
 418 stiffening effects along the geometry and stress concentrations at the proximal and dis-
 419 tal parts. The three-dimensional AAA geometries were meshed using *ABAQUS/CAE*
 420 preprocessor with a minimum of two linear solid hexaedral hybrid elements (C3D8H)
 421 across the thickness of each layer, so detailed results in terms of peak wall stresses can
 422 be obtained. The element sizes were the same as the ones used previously for the plane
 423 strain aorta rings and the uniaxial specimens, but extended into the third dimension
 424 maintaining a proper aspect ratio. Figure 10 shows an example of the typical mesh
 425 used for the simulations. Table 2 shows a quantitative summary of the meshes with
 426 the total number of elements and nodes used for each AAA model.

427 *5.1. Results for the idealized aneurysm models*

428 Distributions of the circumferential stresses, as well as displacement fields in end-
 429 systolic conditions for three different values of β are depicted in Figure 11 and Figure 13,
 430 respectively (only one-half of the geometry cut by a sagittal plane is shown for clarity).

Table 2. Number of elements and nodes used in the three different parametrized geometrical models of AAAs studied.

AAA model	Number of elements		Number of nodes	
	Three-layered	Intact	Three-layered	Intact
$\beta = 1.0$	514,080	466,480	1,091,910	1,005,040
$\beta = 0.6$	481,500	440,608	1,017,288	949,560
$\beta = 0.2$	649,000	548,544	1,370,611	1,202,011

431 First of all, it is important to mention that we have taken the circumferential stress as
432 the prevailing stress, since the maximum principal stresses are almost perfectly aligned
433 with the circumferential direction. This is in agreement with some data on aneurysms
434 that identify normal stresses as a more reliable indicator than Von Mises stress, a
435 yield criterion developed for ductile metals which is not a suitable measure in this case
436 because of the absence of shear stress (Raghavan et al. 2006, 2011). As it can be seen
437 in Figure 11, both the intact and the three-layered artery present a stress gradient
438 through-the-thickness of the aneurysmal wall, in which the inner surface absorbs the
439 maximum circumferential stresses. As shown in Figure 12, this through-the-thickness
440 stress variation is fairly flat for the intact artery, with maximum stress differences that
441 go from 0.20 to 0.23 MPa in the $\beta = 0.2$ geometrical model. However, as previously
442 noticed in the plane strain rings, the three-layered AAA models show a remarkable
443 discontinuous gradient that is manifested in huge stress jumps at the interface between
444 the layers, where the major stress drop is found at the interface between the intima
445 and the media in all the models, with a maximum value of 0.64 MPa in the most
446 asymmetric aneurysm ($\beta = 0.2$).

447 That said, and in good agreement with Rodríguez et al. (2008), it is worth pointing
448 out that the degree of asymmetry is rather considerable: for aneurysms with the same
449 length, wall thickness and diameter of the undilated sections, the peak wall stresses
450 increase by 32% from the symmetric ($\beta = 1.0$) to the most asymmetric geometry
451 ($\beta = 0.2$). Thus, we can say that the geometry, and more specifically the asymmetry of
452 the sac is a determining factor to rupture potential since the strongest stress gradients
453 are always located at inflection points of the curvature. For $\beta = 1.0$ the maximum
454 stress is distributed uniformly around the sac, as well as the highest displacements are,
455 which is logical due to the azimuthal symmetry. In case of the $\beta = 0.6$ and $\beta = 0.2$ ge-
456 ometries, notable stress concentrations occur for both the three-layered and the intact
457 wall at the superolateral part of the sac. By contrast, the maximum displacements are
458 found in the inferior part, which is fairly flat. This phenomena responds to the prin-
459 ciples of the membrane theory of shells: because the artery wall can be considered as
460 a structural element with a small thickness compared to the other dimensions, we can
461 say that the stiffening at the inflection points is due to the combination between mem-
462 brane and bending forces produced by the curvature, while the flatness of the inferior
463 part only generates bending forces which leads to a softer response with larger displace-
464 ments. Table 3 summarises the peak wall stresses obtained for each model and establish
465 a comparison, in terms of stresses and displacements, between the three-layered and
466 the intact models with respect to the former one. Regarding the peak stress values, we
467 observe an overall stress difference of about 30% that slightly increases with asymme-
468 try, reaching a maximum $\Delta\sigma_{\max}$ of 35.9% for the $\beta = 0.2$ aneurysm. Contrastingly,
469 variations in displacements decrease with the asymmetry from a noteworthy ΔU_{\max}
470 about 54% for $\beta = 1.0$ to a insignificant difference of 0.75 % for the $\beta = 0.2$ geometry.

471 To assess the effects of the heterogeneity of the aneurysmatic wall and the material

472 anisotropy, we have compared our results for the hyperelastic three-layered anisotropic
 473 (H3A) AAA wall with the results obtained by Scotti et al. (2005, 2008), where also
 474 peak wall stresses and displacements are analyzed in parametrized aneurysms. We
 475 have chosen this study to establish a comparison, given that the parameterization of
 476 the AAA geometries is the same, considering $\beta = r/R$ to define the asymmetry of sac,
 477 and a similar systolic pressure of 15.7 kPa (118 mmHg). Comparisons are made between
 478 our H3A wall, in which each layer work independently, a elastic isotropic mono-layered
 479 (EIM) and a hyperelastic isotropic mono-layered wall (HIM), based on the Mooney-
 480 Rivlin constitutive model. The results of the comparison are collected in Table 4. In
 481 terms of stresses, first of all we observe how the percentage difference increases with the
 482 asymmetry: for the elastic wall (EIM) the range of difference is between 43% and 54%,
 483 while for the HIM, given that the hyperelastic wall can undergo larger deformations
 484 than the elastic one, and therefore develops higher stresses, the differences are between
 485 38% and 51%, both maximum differences associated with the most asymmetric AAA
 486 ($\beta = 0.2$). Regarding the displacements, for the EIM wall the differences become
 487 greater as the asymmetry increases, reaching a maximum variation of 48%. Strikingly,
 488 for asymmetry values of $\beta = 1.0$ and $\beta = 0.6$ the HIM wall undergo larger deformations
 489 than the H3A wall, with a maximum difference of 9.43%, probably due to the stiffening
 490 effect produced by tunica intima in the inner surface of the artery. The displacements of
 491 the most asymmetric AAA geometries, $\beta = 0.2$, are quite similar with a small variation
 492 of just 0.75%.

493 One of the main advantages of the layer-separated models is that we can easily
 494 isolate the layers to see the maximum stresses of each one. The circumferential stress
 495 distributions of adventitia, media and intima during systole are depicted in Figure 14
 496 for the three different AAA models. As it can be seen, the patterns of circumferential
 497 stresses remain unchanged from the anterior to the posterior wall of the AAA in all
 498 cases, with a uniform distribution around the sac for the symmetric model, and stress
 499 concentrations at the inflection points of the curvature of the sac in the asymmetric
 500 aneurysms as previously commented, which means that, despite the significant stress
 501 jumps found in Figure 12, there is a strong stress transmission from the inner to the
 502 outer wall of the sac that is damped by the tunica intima, which acts as a natural
 503 stiffener for the artery. Table 5 summarizes the percentage of stress absorbed by each
 504 layer with respect to the total circumferential stress. This percentage is measured by
 505 what we have called mean stress absorption. As shown, the intima is the stiffest layer,
 506 absorbing a minimum of 0.443 MPa and maximum of 0.645 MPa during peak systolic,
 507 which leads to stress absorptions of 80.49% and 80.42% for the $\beta = 1.0$ and $\beta = 0.2$
 508 respectively and a mean absorption of 78.33%. The adventitia is the softest layer, with
 509 a range of values between 0.193 and 0.251 MPa, and a mean stress absorption of 7.68%,
 510 while the media is a bit stiffer with a 11.05%. These results are in accordance with the
 511 previously analyzed uniaxially loaded aorta strips and the inflated plane strain rings,
 512 where the early stiffening effect of the intima due to the high dispersion of the collagen
 513 fibers was predicted.

514 The computed circumferential stresses in our idealized AAAs are similar to the
 515 circumferential Cauchy failure stresses reported by Sassani et al. (2015): 0.51 MPa for
 516 the intima, 1.09 MPa for the media and 1.73 MPa for the adventitia. According to these
 517 failure stress values, we can conclude that the intima layer ruptures for asymmetry
 518 values of $\beta = 0.6$ and $\beta = 0.2$. It is only in the case of azimuthal symmetry, with
 519 $\beta = 1.0$, that the intima withstands the load exerted by the luminal pressure. This
 520 conclusion is consistent with two statements that reinforce the importance of including
 521 the intima in multi-layered models of AAAs: the rupture begins in the intima; and the

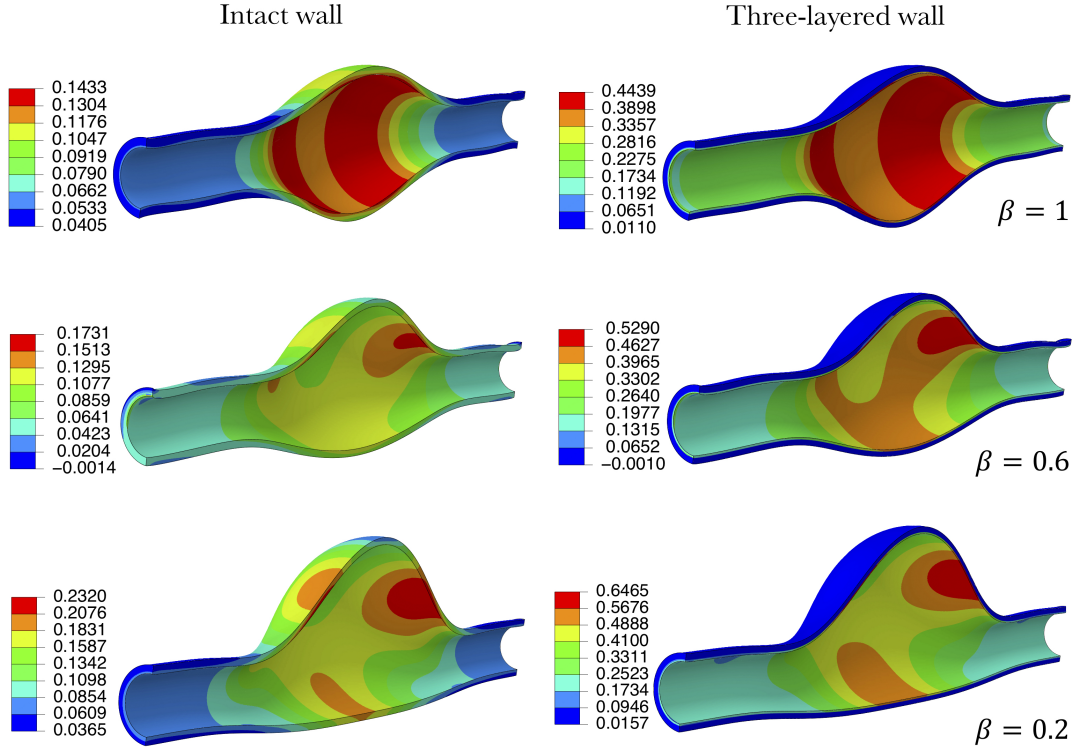


Figure 11. Contour plots of the circumferential stresses in the intact and in the three-layered aneurysmatic wall for asymmetry values of $\beta = 1.0$, $\beta = 0.6$ and $\beta = 0.2$ during peak systolic (MPa).

522 asymmetry of the sac increases the probability of aneurysm rupture. Even though the
 523 peak wall stresses will be smoothed if we include residual stresses, the results obtained
 524 by our idealized AAAS are accurate enough.

525 6. Conclusions

526 This investigation attempts to demonstrate the importance of considering the mechan-
 527 ical contribution of the three layers that make up aortic tissue during the development
 528 of intimal hyperplasia. To do this, finite element analyses were performed on three dif-
 529 ferent idealized geometries of AAA models subjected to realistic loading and boundary
 530 conditions. These simulations were calibrated considering the structural response of
 531 the aneurysmal tissue through uniaxial tests of aorta strips cut in the circumferential

Table 3. Maximum circumferential wall stresses σ_{\max} and displacements U_{\max} in the differ- ent asymmetric AAA models and comparison between the three-layered and the intact artery wall. $\Delta\sigma_{\max}$ and ΔU_{\max} show the % difference of the stress and displacement obtained with the three-layered and intact AAA models with respect to the baseline three-layered method.

AAA model	σ_{\max} (MPa)		U_{\max} (mm)		$\Delta\sigma_{\max}\%$	$\Delta U_{\max}\%$
	Three-layered	Intact	Three-layered	Intact		
$\beta = 1.0$	0.44	0.14	2.65	1.22	31.8	53.96
$\beta = 0.6$	0.52	0.17	3.76	3.24	32.7	13.82
$\beta = 0.2$	0.64	0.23	6.65	6.60	35.9	0.75

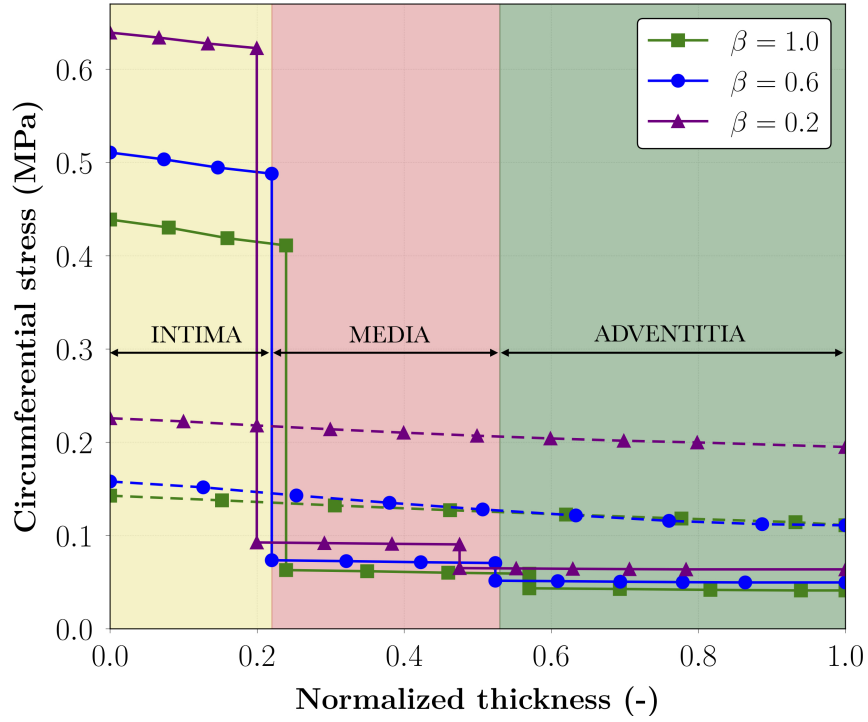


Figure 12. Through-the-thickness circumferential stresses in the three-layered (solid curves) and intact (dashed curves) AAA wall for asymmetry values of $\beta = 1.0$, $\beta = 0.6$ and $\beta = 0.2$.

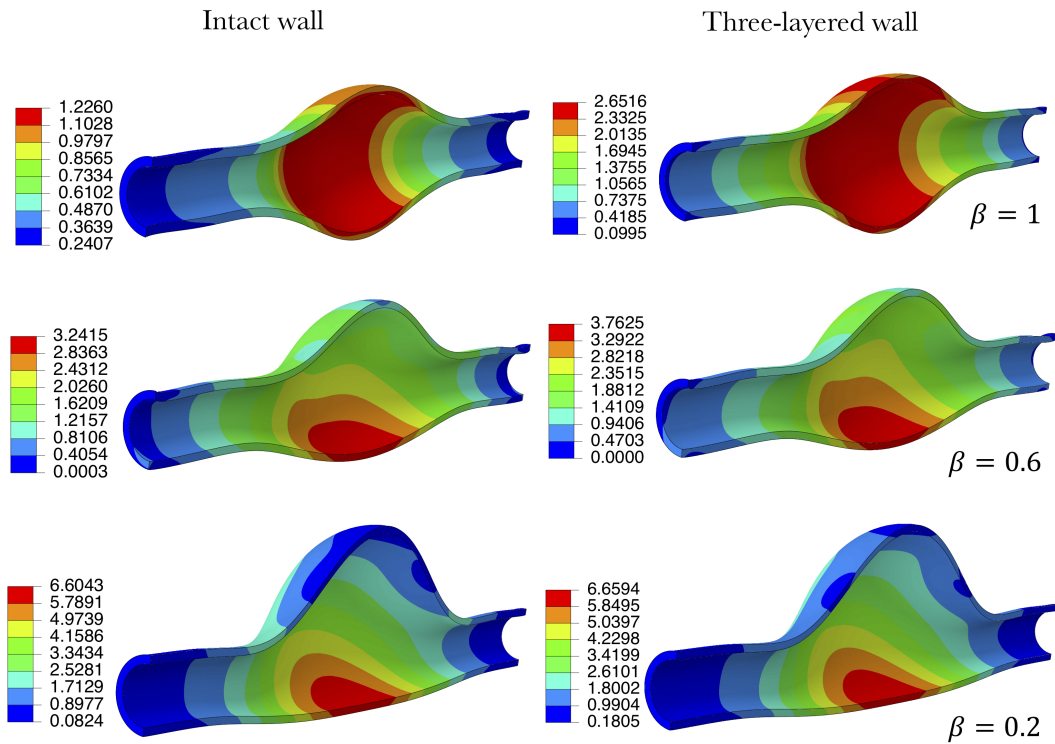


Figure 13. Displacement fields in the intact and in the three-layered aneurysmal wall for asymmetry values of $\beta = 1.0$, $\beta = 0.6$ and $\beta = 0.2$ during peak systolic, in mm.

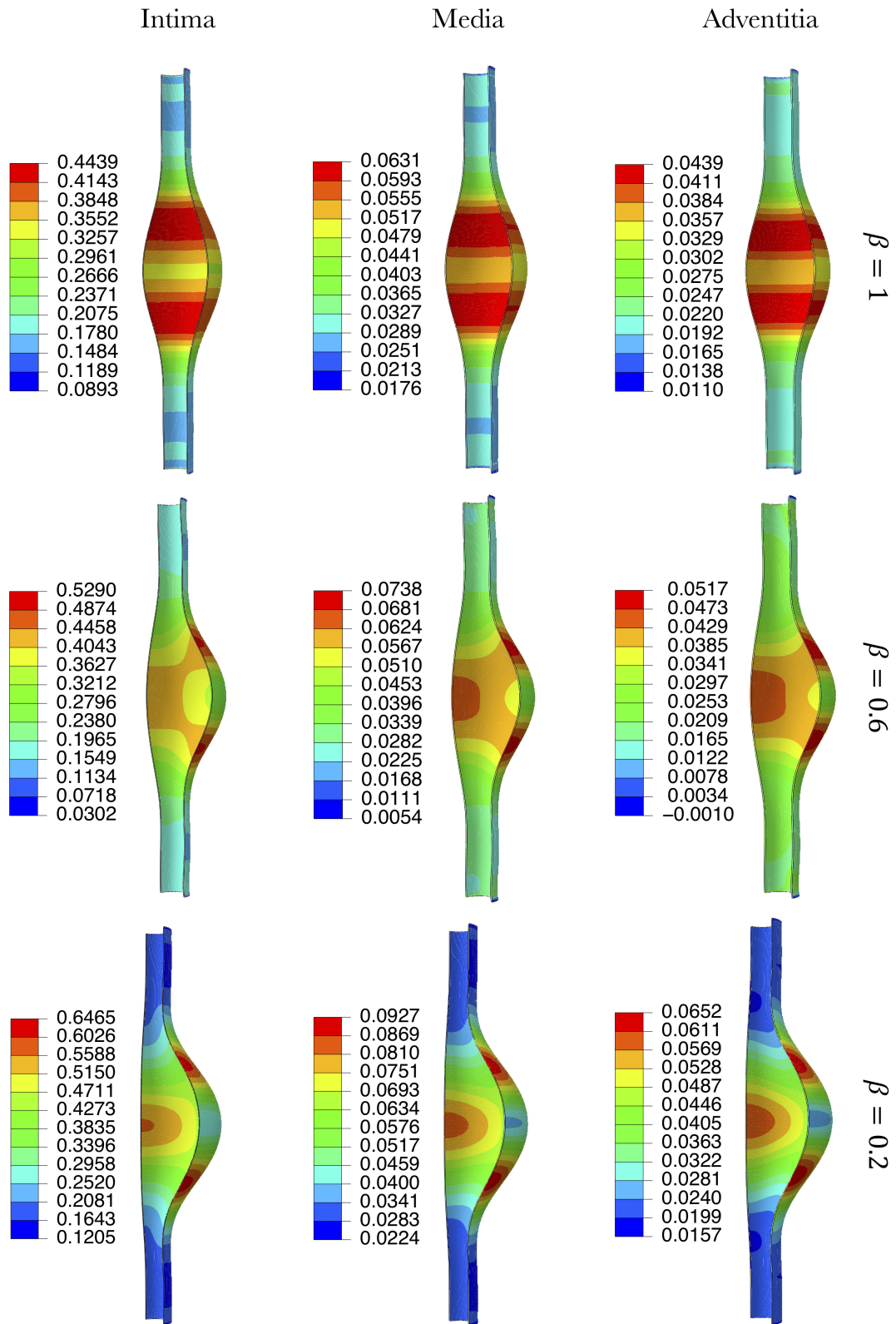


Figure 14. Circumferential stress distributions in adventitia, media and intima layers for $\beta = 1.0$, $\beta = 0.6$ and $\beta = 0.2$ models. The magnitude of the stress is given in MPa.

Table 4. Maximum wall stresses σ_{\max} and displacements U_{\max} in the different asymmetric hyperelastic anisotropic three-layered (H3A) AAA models at a peak systolic pressure of 16 kPa (120 mmHg) and comparison with the results obtained by Scotti et al. (2005, 2008) for the same geometries and a similar systolic blood pressure of 15.7 kPa for an elastic isotropic mono-layer (EIM) and a hyperelastic isotropic (Mooney-Rivlin) mono-layer AAA wall. The parenthesis show the % differences of the EIM and the HIM with respect to the H3A.

AAA model	σ_{\max} (MPa)			U_{\max} (mm)		
	EIM	HIM	H3A	EIM	HIM	H3A
$\beta = 1.0$	0.25(-43.18%)	0.27(-38.63%)	0.44	1.71(-35.47%)	2.90(+9.43%)	2.65
$\beta = 0.6$	0.26(-50.00%)	0.28(-46.15%)	0.52	2.50(-33.51%)	4.30(+14.36%)	3.76
$\beta = 0.2$	0.29(-54.68%)	0.31(-51.56%)	0.64	3.40(-48.87%)	6.60(-0.75%)	6.65

Table 5. Maximum circumferential wall stresses (MPa) in adventitia, media and intima and mean percentage of peak wall stress absorbed by each layer for the three different AAA models. The parenthesis show the percentage of stress absorbed by each layer with respect to the total circumferential stress.

AAA model	σ_{\max} (MPa)		
	Adventitia	Media	Intima
$\beta = 1.0$	0.043(7.83%)	0.063(11.49%)	0.443(80.69%)
$\beta = 0.6$	0.051(7.12%)	0.073(10.19%)	0.529(73.88%)
$\beta = 0.2$	0.065(8.10%)	0.092(11.47%)	0.645(80.42%)
Mean stress absorption	7.68%	11.05%	78.33%

532 and axial directions and plane strain human aorta rings under systolic blood pressure.
533 Resultant stresses and displacements were obtained for an intact (mono-layered)
534 artery wall, which represents the adventitia-intima-media in a single layer with av-
535 eraged properties, and a three-layered wall, in which each layer has been modeled
536 separately with its own material properties using a continuum mesh. We observed
537 differences of about 30% concerning the stresses and a maximum of 53% in displace-
538 ments. This comparison was also carried out with other studies performed on idealized
539 AAA geometries with the same parameterization, but with different constitutive mod-
540 els. We found out maximum differences of 54% in terms of stresses and 48% in terms
541 of displacements for an elastic isotropic mono-layered wall (EIM), and 51% and 14%
542 for a hyperelastic isotropic mono-layer wall (HIM), with respect to a three-layered
543 hyperelastic anisotropic wall (H3A).

544 Regarding the idealized geometries used, our results corroborate that the stress dis-
545 tribution is strongly dependent on the asymmetry of the sac. Symmetric AAAs showed
546 a uniform stress distribution around the sac, while the most asymmetric geometries
547 presented noteworthy stress concentrations at the inflection points of the curvature of
548 the sac, leading to greater peak wall stresses, and therefore higher rupture potential.

549 The obtained results show an early exponential stiffening of tunica intima that makes
550 it definitely load-bearing when it becomes thickened because of intimal hyperplasia.
551 Intimal hyperplasia may be caused by two factors. The first one is the collageniza-
552 tion; The diffuse thickening of the innermost layer of the abdominal aorta has been
553 associated by many studies with collagenization of the elastic and hyper-plastic layers,
554 which increases the dispersion in the families of collagen fibers and stiffens up the in-
555 timal layer. The second factor could be related to the proliferation of smooth muscle
556 cells between the endothelium and the internal elastic lamina. In any case, the intimal
557 layer shows the highest percentages of stress absorption.

558 Our results suggest that tunica intima acts as a stiffener when thickened due to

559 hyperplasia. The intima shows a mean stress absorption ratio of 78%, while the ad-
560 ventitial and medial layers only absorb 7% and 11% of the total circumferential stress
561 on the AAA, respectively. The peak wall stresses in the three-layered AAA models are
562 comparable to the experimental results obtained by Kobielarz et al. (2017) for young
563 arteries in the initial stages of atherosclerotic process, and the investigations carried
564 out by Akyildiz et al. (2014) for human atherosclerotic intima tissue. These facts point
565 out the necessity of including tunica intima in multi-layered models of AAAs to obtain
566 accurate peak wall stresses, and to improve the rupture risk assessment.

567 The models developed in this investigation have three important limitations. First,
568 the material properties are obtained from non-aneurysmal aortas and then applied
569 to idealized AAA geometries. Second, the constitutive model used only considers two
570 families of collagen fibers. And last but not least, residual stresses are not included.

571 The first two limitations will be bypassed in forthcoming publications by considering
572 the mechanical properties of aneurysmal abdominal aortas as in (Kobielarz et al. 2017),
573 and by incorporating a third family of collagen fibers as in (Sassani et al. 2015). Even
574 though these two new assumptions will make our models more precise in predicting
575 the risk of rupture, we do not expect substantial changes in the overall mechanical
576 response of the multilayered wall.

577 If residual stresses are included, we expect a decrease in the compressive stresses
578 in the intima and the tensile stresses in the media and adventitia. In agreement with
579 Sokolis (2015), residual stresses will lessen the stress jumps at the interface between the
580 intima and the media. Nevertheless, we believe that the thickened intima will remain
581 as the layer with the higher stress absorption ratio.

582 This research is a first step in the development of a three-layered model to simulate
583 the mechanical behavior of the abdominal aorta in a more precise way. The model
584 should be improved in future works by incorporating residual stresses and layer-specific
585 properties obtained from aneurysmal aortas.

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