The optimal convergence rate of monotone schemes for conservation laws in the Wasserstein distance

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July 9, 2019

Abstract

In 1994, Nessyahu, Tadmor and Tassa studied convergence rates of monotone finite volume approximations of conservation laws. For compactly supported, Lip⁺-bounded initial data they showed a first-order convergence rate in the Wasserstein distance. Our main result is to prove that this rate is optimal. We further provide numerical evidence indicating that the rate in the case of Lip⁺-unbounded initial data is worse than first-order.

1 Introduction

In their 1994 paper, Nessyahu, Tadmor and Tassa [8] showed that a large class of monotone finite volume methods converge to the entropy solution of the hyperbolic conservation law

$$u_t + f(u)_x = 0, x \in \mathbb{R}, t > 0,$$

 $u(x,0) = u_0(x),$
(1.1)

at a rate of $O(\Delta x)$ in the 1-Wasserstein distance W_1 (using the different name Lip') under the assumption that f is strictly convex ($f'' \ge \alpha > 0$) and the initial datum u_0 is compactly supported and Lip⁺-bounded, i.e.

$$\frac{u_0(x+z) - u_0(x)}{z} \leqslant C, \qquad \forall \ x, z \in \mathbb{R}, \quad z \neq 0.$$
 (1.2)

Recently, Fjordholm and Solem [2] showed a convergence rate of $O(\Delta x^2)$ in W_1 for initial data consisting of finitely many shocks. This raises the question whether the first-order rate in W_1 of [8] can be improved. In this paper we show that this is not possible. We construct a compactly supported and Lip⁺-bounded initial datum for which the convergence rate in W_1 is no better than first-order. In other words, the rate $O(\Delta x)$ in [8] is optimal.

1.1 The Wasserstein distance

In one dimension, the W_1 -distance between two functions u and v takes on the simple form

$$W_1(u,v) = \int_{\mathbb{R}} \left| \int_{-\infty}^x (u(y) - v(y)) \, dy \right| dx,$$

see [15] for more details. In higher dimensions it is the dual of the Lip-norm, referred to as Lip' in [9]. This distance was first utilized in the context of conservation laws in a series of papers by Nessyahu, Tadmor and Tassa [12, 7, 8] where they, among other things, prove convergence rates for a large class of approximations to the solution of the conservation law (1.1) in the W_1 -distance.

Heuristically, one can think of the W_1 -distance as measuring the minimal amount of work needed to move mass from one place to another. In the case of increasing initial data, a monotone scheme provides an approximation of the type shown in Figure 1 after some time has elapsed. Given that the L^1 error

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A. M. Ruf has received funding from the European Unions Framework Programme for Research and Innovation Horizon 2020 (2014-2020) under the Marie Skłodowska-Curie Grant Agreement No. 642768. E. Sande was supported by the European Research Council under the European Unions Seventh Framework Programme (FP7/2007-2013) / ERC grant agreement 339643.

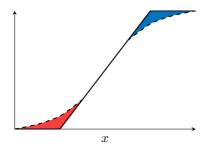


Figure 1: Exact and approximate solution of (1.1).

is $O(\Delta x)$, the surplus of mass on the left (red area) needed to be moved is $O(\Delta x)$ and it needs to be moved a distance of O(1) to the shortage of mass on the right (blue area), see [11] for a similar argument. Therefore, we expect the W_1 -error to be no better than $O(\Delta x) \cdot O(1) = O(\Delta x)$ in this case. The goal of this paper is to make this heuristic argument rigorous.

1.2 Convergence rates for monotone schemes in L^1 and W_1

The L^1 convergence rate for monotone schemes is restricted to $O(\Delta x)$ as these schemes are at most first-order accurate (see Harten, Hyman, and Lax [5]). However, the generic result on convergence rates of monotone schemes for the conservation law (1.1) is the $O(\Delta x^{1/2})$ rate in L^1 , due to Kuznetsov [6] which dates back to 1976. By constructing a (pathological) initial datum, Şabac showed in 1997 that the $O(\Delta x^{1/2})$ rate for monotone methods is, in fact, optimal and cannot be improved without further assumptions on the class of initial data [10] (see [13] for the linear advection equation). For that, Şabac assumed strict convexity of the flux $(f'' \ge \alpha > 0)$, which is the setting considered in the present paper.

assumed strict convexity of the flux $(f'' \ge \alpha > 0)$, which is the setting considered in the present paper. Although the convergence rate $O(\Delta x^{1/2})$ is optimal, in some special cases higher convergence rates for monotone schemes have been shown. For example Harabetian [4] proved that monotone schemes for centered rarefaction waves converge at a rate of $O(\Delta x |\log \Delta x|)$ in L^1 , which is claimed to be optimal in [16]. Before that, Bakhvalov [1] proved the same rate for an upwind scheme in a weaker norm related to W_1 . Wang [16] showed that the rate $O(\Delta x |\log \Delta x|)$ in L^1 also appears close to the critical time of shock formation in certain special cases. Furthermore, Teng and Zhang [14] proved that monotone schemes converge at the optimal rate of $O(\Delta x)$ in L^1 provided the initial datum is piecewise constant with a finite number of discontinuities that only allow for shocks at later times (in the case of a convex flux this means only downward jumps). Later, this result was extended to the convergence rate $O(\Delta x^2)$ in W_1 by Fjordholm and Solem [2].

The seminal work on convergence rates in the Wasserstein distance is Nessyahu and Tadmor's 1992 paper [7]. Using the dual equation studied by Tadmor in [12], the authors showed that conservative, Lip^+ -stable and Lip' -consistent schemes converge at a rate of $O(\Delta x)$ in the Wasserstein distance, for Lip^+ -bounded (i.e., rarefaction-free), compactly supported initial data. Examples of schemes that satisfy these assumptions are the Lax–Friedrichs, Engquist–Osher, and Godunov scheme. Nessyahu, Tadmor and Tassa later used that framework to prove the same convergence rate for so-called Godunov type schemes. In addition to the aforementioned schemes, a subset of (formally) second-order MUSCL schemes also falls into this class. Notably, Nessyahu and Tassa [9] also covered the case of Lip^+ -unbounded initial data and showed a convergence rate of $O(\varepsilon|\log \varepsilon|)$ in W_1 for viscous regularizations of (1.1).

Table 1 provides an overview of the results concerning convergence rates for monotone schemes in both L^1 and the Wasserstein distance.

Remark 1.1. As remarked in [7] and [2] one can recover the well-known half-order rate in L^1 from the first-order rate in W_1 by utilizing the Sobolev interpolation inequality $\|Dg\|_{L^1(\mathbb{R})} \leqslant C\|D^2g\|_{L^1(\mathbb{R})}^{1/2} \|g\|_{L^1(\mathbb{R})}^{1/2}$, see e.g. [3, Thm. 9.3], as follows. Let u be the solution of the conservation law (1.1) and $u_{\Delta x}$ a monotone approximation to it. Then let Dg be a suitable approximation of the error, $u - u_{\Delta x}$, such that $\|D^2g\|_{L^1}$ is bounded by the total variation of u_0 . Then it follows that $\|u - u_{\Delta x}\|_{L^1(\mathbb{R})} \leqslant C \operatorname{TV}(u_0)^{1/2}W_1(u,u_{\Delta x})^{1/2}$. Note that this inequality can also be used in the other direction: The optimality of the convergence rate $O(\Delta x)$ in L^1 for piecewise constant, decreasing initial data implies the optimality of the convergence rate $O(\Delta x^2)$ in W_1 . Moreover, optimality of the convergence rate $O(\Delta x^{1/2})$ in L^1 for general $L^1 \cap BV$ initial data implies that the convergence rate in W_1 cannot be better than $O(\Delta x)$ in the general case.

Case considered	L^1 rate	Optimal	W_1 rate	Optimal
General $L^1 \cap BV$ initial data [6]	$\Delta x^{1/2}$	[10]	_	_
Lip^+ -bounded, compactly supported initial data $[6, 7, 8]$	$\Delta x^{1/2}$	_	Δx	Thm. 3.1
Rarefaction solutions [4]	$\Delta x \log \Delta x $	[4, 16]	_	_
Decreasing data (before shocks) [16]	$\Delta x \log \Delta x $	[16]	_	_
Decreasing piecewise constant initial datum (finitely many pieces) [14, 2]	Δx	[14]	Δx^2	Rem. 1.1

Table 1: Short overview of results regarding rates of convergence for monotone schemes.

2 Preliminaries

In this section we present the class of numerical methods and the observations needed to prove the optimal rate.

2.1 Monotone schemes and first-order convergence in W_1

Let $x_{i-1/2}$, $i \in \mathbb{Z}$, be equidistant points, Δx apart, and let $C_i = [x_{i-1/2}, x_{i+1/2})$, $t^n = n\Delta t$, $n \in \mathbb{N}$, and $\lambda = \Delta t/\Delta x$. Then we consider schemes of the form

$$u_i^{n+1} = u_i^n - \lambda \left(F(u_i^n, u_{i+1}^n) - F(u_{i-1}^n, u_i^n) \right),$$

$$u_i^0 = \frac{1}{\Delta x} \int_{C_i} u_0(x) \, dx,$$
(i \in \mathbb{Z})
(2.1)

where the numerical flux function $F(\cdot,\cdot)$ is consistent with the flux f, i.e. F(u,u)=f(u). The scheme is monotone if and only if $F(\cdot,\cdot)$ is nondecreasing in the first argument and nonincreasing in the second and a certain CFL condition is satisfied.

A numerical approximation $u_{\Delta x}$, defined by $u_{\Delta x}(x,t) = u_i^n$ for $(x,t) \in \mathcal{C}_i \times [t^n, t^{n+1})$, is said to be discrete Lip⁺-stable if

$$||u_{\Delta x}(t)||_{\mathrm{DLip}^+} := \max_{x} \frac{u_{\Delta x}(x + \Delta x, t) - u_{\Delta x}(x, t)}{\Delta x} \leqslant C, \quad \forall t > 0,$$

and a numerical method is called W_1 -consistent in [7, 8] if

$$W_1(u_{\Delta x}(\cdot,0),u_0) + W_1((u_{\Delta x})_t, -f(u_{\Delta x})_x) \leqslant O(\Delta x),$$

where the second term is the Wasserstein distance for functions in space-time. Here, the derivatives involving $u_{\Delta x}$ are defined in [8]. Let \mathcal{S}_t be the exact evolution operator to (1.1), i.e. such that $\mathcal{S}_t u(\cdot, s) = u(\cdot, s+t)$, and let \mathcal{A} be the piecewise constant projection operator,

$$\mathcal{A}u(x) = \frac{1}{\Delta x} \int_{\mathcal{C}_i} u(y) \, dy, \quad x \in \mathcal{C}_i.$$
 (2.2)

Then the method (2.1) is W_1 -consistent if it can be rewritten in a form $u_i^{n+1} = P\mathcal{A}\mathcal{S}u_{\Delta x}(t^n+)$. Here $u_{\Delta x}(t^n+)$ is the numerical approximation calculated with (2.1) at t^n and P is a scheme-dependent projection operator that satisfies

$$W_1(PAv, v) \leq O(\Delta x^2) \operatorname{TV}(v),$$
 (2.3)

see [8, Thm. 2.1]. Assuming that the numerical approximation is Lip^+ -stable and W_1 -consistent, Nessyahu et al. proved the following.

Theorem 2.1 (Nessyahu et al. [8, Thm. 2.3]). Assume that u_0 is compactly supported and Lip^+ -bounded, see (1.2), and that the numerical approximation $u_{\Delta x}$ is discrete Lip^+ -stable and W_1 -consistent. Then for any T > 0,

$$W_1(u(t), u_{\Delta x}(t)) \leqslant C_T \Delta x,$$

for all $0 \le t \le T$.

Note that this theorem includes all monotone schemes and even some (formally) higher-order schemes, as long as they are discrete Lip⁺-stable and satisfy (2.3). Examples of monotone numerical schemes that satisfy these assumptions are the Godunov, Lax-Friedrichs and Enquist-Osher schemes, see [8].

The Godunov scheme is one (or the) example of a monotone scheme. It consists of piecewise constant projection and exact evolution in time. Using the operators S and A, it can be written in the simple form

$$u_{\Delta x}(x,t) = \left(S_{t-t^n} \mathcal{A} \left(S_{\Delta t} \mathcal{A}\right)^n u_0\right)(x,t), \qquad t \in [t^n, t^{n+1}). \tag{2.4}$$

2.2 The error equation

Let $E(x,t) = \int_{-\infty}^{x} (u(y,t) - v(y,t)) dy$, where u and v are two solutions to (1.1), possibly with different initial data u_0 and v_0 . Then

$$W_1(u(t), v(t)) = ||E(t)||_{L^1(\mathbb{R})}.$$

Due to the L^1 -contraction property of solutions to (1.1), E is Lipschitz continuous in both time and space. Now let

$$a(u,v) = \int_0^1 f'(\gamma u + (1-\gamma)v) \, d\gamma.$$

Then, as u and v are solutions of (1.1), E satisfies the transport equation

$$E_t + a(u, v)E_x = 0,$$

$$E(x, 0) = \int_{-\infty}^x (u(y, 0) - v(y, 0)) dy.$$
(2.5)

Hence, if $E(0) \ge 0$, then $E(t) \ge 0$ at any later time t > 0. Given nondecreasing initial data and a strictly convex flux, u and v are continuous for t > 0 and thus a(u, v) is as well. It follows that (2.5) is well-defined for any t > 0.

After an integration by parts, we can see that the time-derivative of the Wasserstein distance between two solutions of (1.1) satisfies

$$\frac{d}{dt} \|E(t)\|_{L^{1}(\mathbb{R})} = -\int_{\mathbb{R}} a(u,v)(t) \partial_{x} |E(t)| dx = \int_{\mathbb{R}} D_{x} a(u,v)(t) |E(t)| dx,$$

where $D_x a(u, v)$ is to be understood as the distributional derivative of a(u, v). Note that since the flux is strictly convex $(f'' \ge \alpha > 0)$ we have

$$D_x a(u,v) \geqslant \frac{\alpha}{2}(u_x + v_x).$$

It follows that $D_x a(u, v)$, and consequently $\frac{d}{dt} ||E(t)||_{L^1(\mathbb{R})}$, is nonnegative, if u and v are nondecreasing.

2.3 The projections

This section contains two useful lemmas on the projection operator \mathcal{A} in the case of nondecreasing functions. The first one shows that the primitive of the projection error is nonnegative.

Lemma 2.2. For a nondecreasing function v, the projection operator A defined in (2.2) satisfies

$$\int_{-\infty}^{x} (Av - v)(y) \, dy \geqslant 0 \tag{2.6}$$

for all $x \in \mathbb{R}$.

Proof. For $x \in C_i$ we find that

$$\int_{-\infty}^{x} (\mathcal{A}v - v)(y) \, dy = \int_{-\infty}^{x_{i-1/2}} (\mathcal{A}v - v)(y) \, dy + \int_{x_{i-1/2}}^{x} (\mathcal{A}v - v)(y) \, dy,$$

where the first term vanishes due to conservation of mass of A. As Av is constant in C_i and is the average of the function v which is nondecreasing,

$$\int_{x_{i-1/2}}^{x} (\mathcal{A}v - v)(y) \, dy \geqslant 0,$$

and we can conclude that (2.6) holds.

The second lemma states that the projection operator \mathcal{A} preserves positivity of the difference between the primitives.

Lemma 2.3. Let u and v be two nondecreasing functions. Then

$$\int_{-\infty}^{x} (v - u)(y) \, dy \geqslant 0 \quad \forall \ x \in \mathbb{R} \quad \Rightarrow \quad \int_{-\infty}^{x} (\mathcal{A}v - u)(y) \, dy \geqslant 0 \quad \forall \ x \in \mathbb{R}. \tag{2.7}$$

Proof. Assume that the inequality to the left in (2.7) holds. Then, using Lemma 2.2,

$$\int_{-\infty}^{x} (\mathcal{A}v - u)(y) \, dy = \int_{-\infty}^{x} (v - u)(y) \, dy + \int_{-\infty}^{x} (\mathcal{A}v - v)(y) \, dy \geqslant 0.$$

3 Optimality

With the observations in the previous section, we can prove that the first-order rate in Theorem 2.1 is optimal in W_1 :

Theorem 3.1. Let f be strictly convex and let u_0 be compactly supported and Lip^+ -bounded, i.e., satisfy (1.2). Then the optimal convergence rate in the Wasserstein distance of monotone finite volume schemes (2.1) satisfying (2.3), is $O(\Delta x)$.

We postpone the (short) proof to the end of the section.

The initial datum u_0 has to be Lip⁺-bounded and compactly supported for Theorem 2.1 to hold. We will therefore consider compactly supported initial data u_0 consisting of one increasing, Lip⁺-bounded part, increasing from 0 to M and one decreasing part, decreasing from M to 0. One realization of a suitable initial datum is

$$u_{0}(x) = \begin{cases} 0, & x < x_{s}, \\ M \frac{x - x_{s}}{x_{0} - x_{s}}, & x_{s} \leq x < x_{0}, \\ M, & x_{0} \leq x < x_{M}, \\ M \left(1 - \frac{x - x_{M}}{x_{e} - x_{M}}\right), & x_{M} \leq x < x_{e}, \\ 0, & x \geqslant x_{e}, \end{cases}$$
(3.1)

where $x_s < x_0 < x_M < x_e$, and $[x_s, x_e]$ is the support of u_0 , see Figure 2.

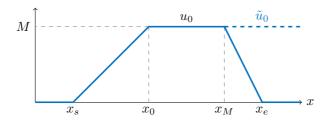


Figure 2: The initial datum used to show optimality of the convergence rate

Since we only expect the increasing part of u_0 to contribute to the reduced convergence rate $O(\Delta x)$, we will simplify our calculations by separating the increasing parts of u and $u_{\Delta x}$ from the decreasing parts. To that end, for any fixed time T > 0, without restrictions, we will assume that $x_M - x_0$ is big enough such that there exists an x^* satisfying $x_0 < x^* < x_M$ where $u_{\Delta x}(x^* + f'(M)t, t) = M$ for all $0 \le t \le T$. Using x^* we introduce the increasing auxiliary functions

$$\tilde{u}_{0}(x) = \begin{cases} u_{0}(x), & x < x^{*} \\ M, & x \geqslant x^{*}, \end{cases}
\tilde{u}(x,t) = \begin{cases} u(x,t), & x < x^{*} + f'(M)t, \\ M, & x \geqslant x^{*} + f'(M)t, \end{cases}
\tilde{u}_{\Delta x}(x,t) = \begin{cases} u_{\Delta x}(x,t), & x < x^{*} + f'(M)t \\ M, & x \geqslant x^{*} + f'(M)t. \end{cases}$$
(3.2)

Then, the assumption above implies

$$\int_{-\infty}^{x^* + f'(M)t} (\tilde{u} - \tilde{u}_{\Delta x}) \ dy = \int_{-\infty}^{x^* + f'(M)t} (u - u_{\Delta x}) \ dy = 0,$$

and therefore

$$W_1(u(t), u_{\Delta x}(t)) \geqslant W_1(\tilde{u}(t), \tilde{u}_{\Delta x}(t))$$
.

We now estimate the W_1 -error between \tilde{u} and $\tilde{u}_{\Delta x}$ from below.

Proposition 3.2. Let u_0 be as described above and $u_{\Delta x}$ the numerical approximation (2.4). Then for 0 < t < T,

$$W_1\left(u(t), u_{\Delta x}(t)\right) \geqslant \sum_{n=0}^{N} W_1\left(\mathcal{A}\tilde{u}_{\Delta x}(t^n-), \tilde{u}_{\Delta x}(t^n-)\right),$$

where N is such that $t \in [t^N, t^{N+1})$ and $u_{\Delta x}(t^n)$ is the numerical approximation right before averaging.

Proof. Let $t \in [t^N, t^{N+1})$ and $E_{\Delta x}(x, t) = \int_{-\infty}^x (\tilde{u}_{\Delta x} - \tilde{u})(y, t) dy$. Because of Lemma 2.2

$$E_{\Delta x}(x,0) = \int_{-\infty}^{x} (\tilde{u}_{\Delta x} - \tilde{u})(y,0) \, dy = \int_{-\infty}^{x} (\mathcal{A}\tilde{u}_{0} - \tilde{u}_{0})(y) \, dy \geqslant 0$$

for all $x \in \mathbb{R}$. The fact that $E_{\Delta x}$ satisfies the transport equation (2.5) and Lemma 2.3 imply that $E_{\Delta x}$ is nonnegative for all $t \ge 0$. Hence,

$$W_1(u(t), u_{\Delta x}(t)) \geqslant W_1(\tilde{u}(t), \tilde{u}_{\Delta x}(t)) = \int_{\mathbb{R}} E_{\Delta x}(x, t) dx = \int_{\mathbb{R}} \int_{-\infty}^{x} (\tilde{u}_{\Delta x} - \tilde{u})(y, t) dy dx.$$

As \tilde{u}_0 is nondecreasing and the conservation law (1.1) and the scheme (2.1) are monotonicity preserving, \tilde{u} and $\tilde{u}_{\Delta x}$ will be nondecreasing at any later time. It follows from the argument in Section 2.2 that for $t \in [t^N, t^{N+1})$ the W_1 error between \tilde{u} and $\tilde{u}_{\Delta x}$ will be nondecreasing. Hence,

$$\begin{split} W_1\left(\tilde{u}(t),\tilde{u}_{\Delta x}(t)\right) &\geqslant W_1\left(\tilde{u}(t^N),\mathcal{A}\tilde{u}_{\Delta x}(t^N)\right) \\ &= \int_{\mathbb{R}} \int_{-\infty}^x \left(\mathcal{A}\tilde{u}_{\Delta x} - \tilde{u}\right)(y,t^N) \, dy dx \\ &= \int_{\mathbb{R}} \int_{-\infty}^x \left(\mathcal{A}\tilde{u}_{\Delta x} - \tilde{u}_{\Delta x}\right)(y,t^N -) \, dy dx + \int_{\mathbb{R}} \int_{-\infty}^x \left(\tilde{u}_{\Delta x} - \tilde{u}\right)(y,t^N -) \, dy dx, \end{split}$$

where, in the last line, we have added and subtracted $\tilde{u}_{\Delta x}(t^N)$. We can now continue the same procedure on the last term in the above N times, and we end up with

$$W_1\left(u(t), u_{\Delta x}(t)\right) \geqslant \sum_{n=0}^{N} \int_{\mathbb{R}} \int_{-\infty}^{x} \left(\mathcal{A}\tilde{u}_{\Delta x} - \tilde{u}_{\Delta x}\right)(y, t^n -) \, dy dx = \sum_{n=0}^{N} W_1\left(\mathcal{A}\tilde{u}_{\Delta x}(t^n -), \tilde{u}_{\Delta x}(t^n -)\right),$$

which is what we wanted to prove.

In order to conclude that the $O(\Delta x)$ rate is optimal in W_1 , we need to show that for the increasing part of u_0 the projection error $W_1\left(\mathcal{A}\tilde{u}_{\Delta x}(t^n-),\tilde{u}_{\Delta x}(t^n-)\right)$ is bounded from below by $C\Delta x\Delta t$ for any $0 \leq t^n < T$.

Proposition 3.3. Let u_0 be increasing and assume that $\beta_1 \geqslant f' \geqslant \beta_2 > 0$ on [-M, M], where $M = ||u||_{L^{\infty}(\mathbb{R})}$. Then if

$$\lambda \leqslant \frac{1}{2\beta_1}$$

we have

$$W_1\left(\mathcal{A}u_{\Delta x}(t^n-), u_{\Delta x}(t^n-)\right) \geqslant \frac{\Delta x \Delta t}{2} \left(1 - \beta_1 \lambda\right) \beta_2 \text{TV}\left(u_0\right).$$

Proof. From the positivity of the projection error and the conservation of mass in each cell C_i , we get that

$$W_1\left(\mathcal{A}u_{\Delta x}(t^n-), u_{\Delta x}(t^n-)\right) = \int_{\mathbb{R}} \int_{-\infty}^x \left(\mathcal{A}u_{\Delta x} - u_{\Delta x}\right)(y, t^n-) \, dy dx$$
$$= \sum_i \int_{\mathcal{C}_i} \int_{x_{i-1/2}}^x \left(u_i^n - u_{\Delta x}(y, t^n-)\right) \, dy dx.$$

Each term in the sum can be rewritten in the following way,

$$\int_{\mathcal{C}_{i}} \int_{x_{i-1/2}}^{x} (u_{i}^{n} - u_{\Delta x}(y, t^{n} -)) \, dy dx$$

$$= \int_{\mathcal{C}_{i}} (x - x_{i-1/2}) \left(u_{i}^{n} - \frac{1}{(x - x_{i-1/2})} \int_{x_{i-1/2}}^{x} u_{\Delta x}(y, t^{n} -) \, dy \right) dx$$

$$= \int_{\mathcal{C}_{i}} \frac{1}{2} (x - x_{i-1/2})^{2} \left(\frac{u_{\Delta x}(x, t^{n} -)}{x - x_{i-1/2}} - \frac{1}{(x - x_{i-1/2})^{2}} \int_{x_{i-1/2}}^{x} u_{\Delta x}(y, t^{n} -) \, dy \right) dx$$

$$= \frac{1}{2} \int_{\mathcal{C}_{i}} \left(u_{\Delta x}(x, t^{n} -)(x - x_{i-1/2}) - \int_{x_{i-1/2}}^{x} u_{\Delta x}(y, t^{n} -) \, dy \right) dx$$

$$= \frac{1}{2} \int_{\mathcal{C}_{i}} \int_{x_{i-1/2}}^{x} (u_{\Delta x}(x, t^{n} -) - u_{\Delta x}(y, t^{n} -)) \, dy dx, \tag{3.3}$$

after an integration by parts in the variable x from the first to the second line (which is justified by the fact that $u_{\Delta x}$ is continuous for t > 0 as $f' \ge \beta_2 > 0$ and $f'' \ge \alpha > 0$). By integrating only over

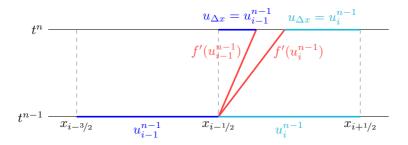


Figure 3: Transportation of the cell averages u_{i-1}^{n-1} and u_i^{n-1} when calculating $u_{\Delta x}$ in C_i .

the part of C_i where $u_{\Delta x}$ is constant and ignoring the Riemann fan between $x_{i-1/2} + f'(u_{i-1}^{n-1})\Delta t$ and $x_{i-1/2} + f'(u_i^{n-1})\Delta t$ (see Figure 3) we can bound the last term (3.3) from below as follows:

$$\begin{split} \frac{1}{2} \int_{\mathcal{C}_i} \int_{x_{i-1/2}}^x \left(u_{\Delta x}(x,t^n-) - u_{\Delta x}(y,t^n-) \right) dy dx \\ &\geqslant \frac{1}{2} \int_{x_{i-1/2} + f'(u_i^{n-1}) \Delta t}^x \int_{x_{i-1/2}}^x \left(u_{\Delta x}(x,t^n-) - u_{\Delta x}(y,t^n-) \right) dy dx \\ &\geqslant \frac{1}{2} \int_{x_{i-1/2} + f'(u_i^{n-1}) \Delta t}^{x_{i+1/2}} \int_{x_{i-1/2}}^{x_{i-1/2} + f'(u_{i-1}^{n-1}) \Delta t} \left(u_{\Delta x}(x,t^n-) - u_{\Delta x}(y,t^n-) \right) dy dx \\ &= \frac{1}{2} f'(u_{i-1}^{n-1}) \Delta t \int_{x_{i-1/2} + f'(u_i^{n-1}) \Delta t}^{x_{i+1/2}} \left(u_{\Delta x}(x,t^n-) - u_{i-1}^{n-1} \right) dx \\ &= \frac{1}{2} \left(\Delta x - f'(u_i^{n-1}) \Delta t \right) f'(u_{i-1}^{n-1}) \Delta t \left(u_i^{n-1} - u_{i-1}^{n-1} \right) \end{split}$$

Then, summing up,

$$\int_{\mathcal{C}_i} \int_{x_{i-1/2}}^x \left(u_{\Delta x}(x, t^n -) - u_{\Delta x}(y, t^n -) \right) \, dy dx \geqslant \frac{\Delta x \Delta t}{2} \left(1 - f'(u_i^{n-1}) \lambda \right) f'(u_{i-1}^{n-1}) \left(u_i^{n-1} - u_{i-1}^{n-1} \right)$$

$$\geqslant \frac{\Delta x \Delta t}{2} (1 - \beta_1 \lambda) \beta_2 \left(u_i^{n-1} - u_{i-1}^{n-1} \right).$$

Summing over all i, the result follows.

We are now ready to prove Theorem 3.1.

Proof of Theorem 3.1. Let f be strictly convex. Without loss of generality we can assume $\beta_1 \geqslant f' \geqslant \beta_2 > 0$ on [-M, M] (otherwise we consider $\hat{u} = u + C$ and $\hat{u}_{\Delta x} = u_{\Delta x} + C$ for some suitable constant C, which will not affect the Wasserstein distance). Combining Proposition 3.3 with Proposition 3.2, we find that

$$W_1(u(t), u_{\Delta x}(t)) \geqslant \frac{\beta_2}{2} (1 - \beta_1 \lambda) t^N \text{TV}(u_0) \Delta x,$$

for $t \in [t^N, t^{N+1})$, which concludes the proof of Theorem 3.1.

4 Numerical experiments

To illustrate our result, we consider two numerical experiments using Burgers' equation,

$$u_t + \left(\frac{u^2}{2}\right)_x = 0,$$

on the interval [-1,1] with the initial data

$$u_0^1(x) = \begin{cases} 0, & x < -0.75, \\ 2x + 1.5 & -0.75 \leqslant x < -0.25, \\ 1, & -0.25 \leqslant x < 0.25, \\ -4x + 2 & 0.25 \leqslant x < 0.5, \\ 0, & x \geqslant 0.5, \end{cases} \quad \text{and} \quad u_0^2(x) = \begin{cases} 0, & x < 0, \\ 1, & x \geqslant 0. \end{cases}$$

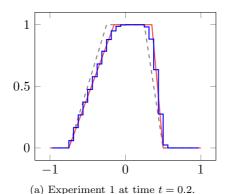
The first initial datum, u_0^1 , is an example of the compactly supported Lip⁺-bounded u_0 in (3.1), and therefore fits into the context considered in this paper. The second initial datum, u_0^2 , on the other hand is Lip⁺-unbounded. For both experiments we use the Godunov scheme, i.e., the monotone scheme (2.1) with the numerical flux function

$$F(a,b) = \frac{1}{2} \max(\max(a,0)^2, \min(b,0)^2),$$

and $\frac{\Delta t}{\Delta x} = 0.5$. The exact solution for Experiment 1 (for t < 0.25) and Experiment 2 is

$$u^{1}(x,t) = \begin{cases} 0, & x < -0.75, \\ 5x/3 + 5/4 & -0.75 \leqslant x < -0.25 + t, \\ 1, & -0.25 + t \leqslant x < 0.25 + t, \\ -20x/3 + 10/3 & 0.25 + t \leqslant x < 0.5, \\ 0, & x \geqslant 0.5, \end{cases}$$
 and
$$u^{2}(x,t) = \begin{cases} 0, & x < 0, \\ x/t, & 0 \leqslant x < t, \\ 1, & x \geqslant t, \end{cases}$$

respectively. Figures 4a and 4b show the initial data for Experiment 1 and 2 (respectively) in gray (dashed), the exact solutions in red (straight), and the numerical approximations calculated with the Godunov scheme in blue (piecewise constant). Tables 2a and 2b show the observed convergence rates of Experiment 1 and 2, where n is the number of cells in the discretization. The first table clearly shows that the W_1 error is $O(\Delta x)$ in Experiment 1 and therefore numerically illustrates the optimality result of the present paper. The second table indicates that in the case of a single upward jump, i.e., Lip⁺-unbounded initial datum, we can expect a convergence rate of $O(\Delta x | \log \Delta x |)$ not only in L^1 as shown by Harabetian [4], but also in W_1 (see also Figure 5).



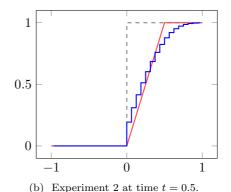


Figure 4: Exact solution and numerical approximation and initial datum.

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	n	L^1 OOC	W_1 OOC
	32	0.822	1.196
	64	0.896	1.123
	128	0.861	1.075
	256	0.884	1.046
	512	0.900	1.029

n	L^1 OOC	W_1 OOC
32	0.598	0.764
64	0.641	0.759
128	0.675	0.761
256	0.708	0.769
512	0.739	0.782

(b) Experiment 2 at time t=0.5.

Table 2: Observed order of convergence in L^1 and W_1 .

5 Concluding remarks

In this paper we have shown optimality of the convergence rate $O(\Delta x)$ in W_1 for monotone schemes in the case of Lip⁺-bounded initial data with compact support, and where the flux is assumed to be strictly convex. As noted in Table 1 it is an open question whether the corresponding L^1 rate of $O(\Delta x^{1/2})$ is also optimal for this case since Şabac's counter-example is Lip⁺-unbounded. Our numerical experiments (see Table 2a) suggest that the counter-example considered here cannot be used to prove optimality of the rate $O(\Delta x^{1/2})$ in L^1 in this case.

The convergence rate in W_1 for Lip^+ -unbounded initial data is still unknown. Our numerical test indicates that in the case of a rarefaction solution it could be the same as the L^1 rate, $O(\Delta x | \log \Delta x |)$. This is consistent with the rate $O(\varepsilon | \log \varepsilon |)$ proved in [9] for the viscous regularization of conservation laws with Lip^+ -unbounded initial data. Furthermore it can be heuristically explained by the same argument as in Section 1.1 since the L^1 error in this case is $O(\Delta x | \log \Delta x |)$ [4].

Finally, to our knowledge there are currently no results on convergence rates in the Wasserstein distance for schemes for one-dimensional systems or for multidimensional conservation laws, although the W_1 -distance can readily be defined in several dimensions.

References

- [1] N. S. Bakhvalov. Estimation of the error of numerical integration of a first-order quasilinear equation. Zhurnal Vychislitel'noi Matematiki i Matematicheskoi Fiziki, 1(5):771–783, 1961.
- [2] U. S. Fjordholm and S. Solem. Second-order convergence of monotone schemes for conservation laws. SIAM J. Numer. Anal., 54(3):1920–1945, 2016.
- [3] A. Friedman. *Partial differential equations*. Robert E. Krieger Publishing Co., Huntington, N.Y., original edition, 1976.
- [4] E. Harabetian. Rarefactions and large time behavior for parabolic equations and monotone schemes. Comm. Math. Phys., 114(4):527–536, 1988.
- [5] A. Harten, J. M. Hyman, and P. D. Lax. On finite-difference approximations and entropy conditions for shocks. *Comm. Pure Appl. Math.*, 29(3):297–322, 1976. With an appendix by B. Keyfitz.

⁽a) Experiment 1 at time t = 0.2.

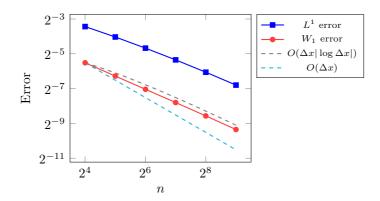


Figure 5: Observed order of convergence in L^1 and W_1 compared with $O(\Delta x |\log \Delta x|)$ and $O(\Delta x)$.

- [6] N. Kuznetsov. Accuracy of some approximate methods for computing the weak solutions of a first-order quasi-linear equation. *USSR Comp. Math. and Math. Physics*, 16(6):105–119, 1976.
- [7] H. Nessyahu and E. Tadmor. The convergence rate of approximate solutions for nonlinear scalar conservation laws. SIAM J. Numer. Anal., 29(6):1505–1519, 1992.
- [8] H. Nessyahu, E. Tadmor, and T. Tassa. The convergence rate of Godunov type schemes. SIAM J. Numer. Anal., 31(1):1–16, 1994.
- [9] H. Nessyahu and T. Tassa. Convergence rate of approximate solutions to conservation laws with initial rarefactions. SIAM J. Numer. Anal., 31(3):628–654, 1994.
- [10] F. Şabac. The optimal convergence rate of monotone finite difference methods for hyperbolic conservation laws. SIAM J. Numer. Anal., 34(6):2306–2318, 1997.
- [11] S. Solem. Convergence rates of the front tracking method for conservation laws in the Wasserstein distances. SIAM J. Numer. Anal., 56(6):3648–3666, 2018.
- [12] E. Tadmor. Local error estimates for discontinuous solutions of nonlinear hyperbolic equations. SIAM J. Numer. Anal., 28(4):891–906, 1991.
- [13] T. Tang and Z. H. Teng. The sharpness of Kuznetsov's $O(\sqrt{\Delta x})$ L^1 -error estimate for monotone difference schemes. *Math. Comp.*, 64(210):581–589, 1995.
- [14] Z.-H. Teng and P. Zhang. Optimal L^1 -rate of convergence for the viscosity method and monotone scheme to piecewise constant solutions with shocks. SIAM J. Numer. Anal., 34(3):959–978, 1997.
- [15] C. Villani. *Topics in optimal transportation*, volume 58 of *Graduate Studies in Mathematics*. American Mathematical Society, Providence, RI, 2003.
- [16] W.-C. Wang. On L^1 convergence rate of viscous and numerical approximate solutions of genuinely nonlinear scalar conservation laws. SIAM J. Math. Anal., 30(1):38–52, 1999.