

Thermal gradient driven domain wall dynamics

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The issue of whether a thermal gradient acts like a magnetic field or an electric current in the domain wall (DW) dynamics is investigated. Broadly speaking, magnetization control knobs can be classified as energy-driving or angular-momentum driving forces. DW propagation driven by a static magnetic field is the best known example of the former in which the DW speed is proportional to the energy dissipation rate, and the current-driven DW motion is an example of the latter. Here we show that DW propagation speed driven by a thermal gradient can be fully explained as the angular momentum transfer between thermally generated spin current and DW. We found DW-plane rotation speed increases as DW width decreases. Both DW propagation speed along the wire and DW-plane rotation speed around the wire decrease with the Gilbert damping. These facts are consistent with the angular momentum transfer mechanism, but are distinct from the energy dissipation mechanism. We further show that magnonic spin-transfer torque (STT) generated by a thermal gradient has both damping-like and field-like components. By analyzing DW propagation speed and DW-plane rotational speed, the coefficient (β) of the field-like STT arising from the non-adiabatic process, is obtained. It is found that β does not depend on the thermal gradient; increases with uniaxial anisotropy K_{\parallel} (thinner DW); and decreases with the damping, in agreement with the physical picture that a larger damping or a thicker DW leads to a better alignment between the spin-current polarization and the local magnetization, or a better adiabaticity.

I. INTRODUCTION

Manipulating domain walls (DW) in magnetic nanostructures has attracted much attention because of its potential applications in data storage technology [1] and logic gates [2]. The traditional DW control knobs, namely magnetic fields and spin-polarized currents, have certain drawbacks in applications. In the magnetic-field-driven DW motion, energy dissipation is the main cause of DW propagation whose speed is proportional to the energy dissipation rate [3, 4], and the magnetic field tends to destroy unfavorable domains and DWs, instead of driving a series of DWs synchronously [5–7]. An electrical current drives a DW to move mainly through the angular momentum transfer so that it pushes multiple DWs [8–11] in the same direction. To achieve a useful DW speed, it requires high electrical current densities that result in a Joule heating problem [12–14]. To avoid these problems, spin-wave spin current has been proposed as a more energy-efficient control parameter [15–18]. Thermal gradient, a way to generate spin-wave spin current, is an alternative control knob of the DW motion. The investigation on thermal-gradient-driven domain wall motion is meaningful not only for conventional applications, but also for the understanding of spin wave and domain wall dynamics [16, 17, 20–23], as well as for possible recycling of waste heat [19, 24].

To understand the mechanism behind thermal-gradient-driven DW dynamics, there are microscopic theories [15–17, 25, 26] and macroscopic thermodynamic theories [21, 22]. Briefly speaking, the microscopic theories suggest that magnons populated in the hotter region diffuses to the colder region to form a magnon spin current. The magnon spin current passes through a DW and exerts a torque on the DW by transferring spin angular momentum to the DW. Thus, magnons drive the DW propagating toward the hotter region of the nanowire, opposite to the magnon current direction [15, 16, 18]. The thermodynamic theories anticipate that a thermal gradient generates an entropy force which always drives the DW towards the hotter region in order to minimize the system free energy. The macroscopic theories do not provide any microscopic picture about DW dynamics although a thermal gradient is often considered as an effective magnetic field to estimate DW speed [21, 22] from field-driven DW theories. Thus, one interesting issue is whether a thermal gradient in DW dynamics acts like a magnetic field or an electric current. DW propagation speed should be sensitive to both DW width and types of a DW (transverse DW) under an energy-driving force while the speed should be insensitive to the DW and DW structure in the angular-momentum-driving force. This is the focus of the current work.

In this paper, we investigate DW motion along a uniaxial wire with the easy axis along the wire direction under a thermal gradient. We found that the DW always propagates to the hotter region with an accompanied DW-plane rotation. DW propagation speed and

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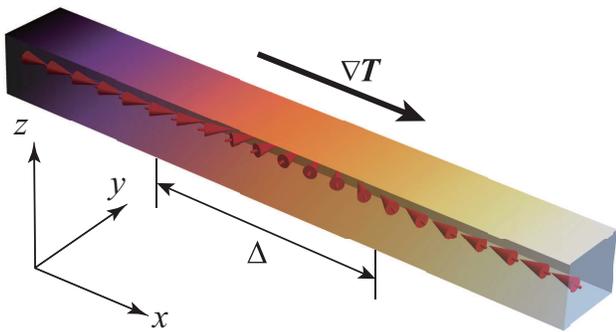


FIG. 1. Schematic diagram of a uniaxial magnetic nanowire with a head-to-head DW at the center under a thermal gradient ∇T . Black (white) color represents colder (hotter) end of the sample.

DW-plane rotation speed increases as the magnetic easy-axis anisotropy and damping decreases. We show that DW motion can be attributed to the angular momentum transfer between magnonic spin current and the DW. Thus, we conclude that a thermal gradient interacts with DW through angular-momentum transfer rather than through energy dissipation. Similar to an electric current [27], a thermal gradient can generate both damping-like (or adiabatic) STT and field-like (or non-adiabatic) STT. From the damping-dependence and anisotropy-dependence of the average DW velocity and DW-plane rotation angular velocity, we extract field-like STT coefficient (β). It is found that β is independent of thermal gradient; is bigger for a thinner DW; and decreases with the damping coefficient. We also show that in the presence of a weak hard-axis anisotropy perpendicular to the wire, the DW still undergoes a rotating motion. The DW propagation speed increases slightly while the DW-plane rotation speed decreases with the strength of the hard-axis anisotropy.

II. MODEL AND METHOD

We consider a uniaxial nanowire of length L_x and cross-section $L_y \times L_z$ along the x -axis (easy axis) with a head-to-head DW at the center, as shown in Fig. 1. L_y, L_z is much smaller than the DW width Δ , and Δ is much smaller than L_x . A thermal gradient is applied along the wire. The highest temperature is far below the Curie temperature T_c . The magnetization dynamics is governed by the stochastic Landau-Lifshitz-Gilbert (LLG) equation [28, 29],

$$\frac{d\mathbf{m}}{dt} = -\gamma \mathbf{m} \times (\mathbf{H}_{\text{eff}} + \mathbf{h}_{\text{th}}) + \alpha \mathbf{m} \times \frac{\partial \mathbf{m}}{\partial t}, \quad (1)$$

where $\mathbf{m} = \mathbf{M}/M_s$ and M_s are respectively the magnetization direction and the saturation magnetization. α is the Gilbert damping constant and γ is the gyromag-

netic ratio. $\mathbf{H}_{\text{eff}} = \frac{2A}{\mu_0 M_s} \sum_{\sigma} \frac{\partial^2 \mathbf{m}}{\partial x_{\sigma}^2} + \frac{2K_{\parallel}}{\mu_0 M_s} m_x \hat{\mathbf{x}} + \mathbf{h}_{\text{dipole}}$ is the effective field, where A is the exchange constant, x_{σ} ($\sigma = 1, 2, 3$) denote Cartesian coordinates x, y, z , K_{\parallel} is the easy-axis anisotropy, and $\mathbf{h}_{\text{dipole}}$ is the dipolar field. \mathbf{h}_{th} is the stochastic thermal field.

The stochastic LLG equation is solved numerically by MUMAX3 package [30] in which we use adaptive Heun solver. To balance stability and efficiency, we choose the time step 10^{-14} s with the cell size $(2 \times 2 \times 2)$ nm³. Magnetic charges at the two ends of the wire are removed to avoid their attraction to the DW. The saturation magnetization $M_s = 8 \times 10^5$ A/m and exchange constant $A = 13 \times 10^{-12}$ J/m are used to mimic permalloy in our simulations. The thermal field follows the Gaussian process characterized by following statistics [31]

$$\begin{aligned} \langle h_{\text{th},ip}(t) \rangle &= 0, \\ \langle h_{\text{th},ip}(t) h_{\text{th},jq}(t + \Delta t) \rangle &= \frac{2k_B T_i \alpha_i}{\gamma \mu_0 M_s a^3} \delta_{ij} \delta_{pq} \delta(\Delta t), \end{aligned} \quad (2)$$

where i and j denote the micromagnetic cells, and p, q represent the Cartesian components of the thermal field. T_i and α_i are respectively temperature and the Gilbert damping at cell i , and a is the cell size. k_B is the Boltzmann constant [28]. The numerical results presented in this study are averaged over 15 random configurations (for DW velocity) and 4000-5000 random configurations (for spin current).

Under the thermal gradient $\nabla_x T$, magnetization at different positions deviate from their equilibrium directions differently and small transverse components m_y and m_z are generated. The transverse components vary spatial-temporally and depend on the local temperature. This variation generates a magnonic spin current [16]. This magnonic spin current can interact with spin textures such as DWs. In the absence of damping (the thermal field also vanishes), the spin current along the x direction can be defined from the spin continuity equation derived from Eq. (1) as follows [15],

$$\frac{\partial \mathbf{m}}{\partial t} = -\frac{1}{1 + \alpha^2} \mathbf{m} \times \hat{\mathbf{x}} m_x K_{\parallel} - \frac{\partial \mathbf{J}}{\partial x}, \quad (3)$$

where

$$\mathbf{J}(x) = \frac{2\gamma A}{\mu_0 M_s} \mathbf{m} \times \frac{\partial \mathbf{m}}{\partial x}, \quad (4)$$

is the spin current density along x -direction due to the exchange interaction. $\mathbf{J}(x)$ can be numerically calculated [15, 23]. In the presence of damping as well as the thermal field, the contribution of the damping term and the thermal term is proportional to α , which is relatively small. More importantly, according to the fluctuation-dissipation theorem [28], the damping term and the thermal term should cancel each other after average over a long time. Since the time scale of DW dynamics is much longer than the thermal fluctuation, the combined contribution of damping and thermal terms should be very small.

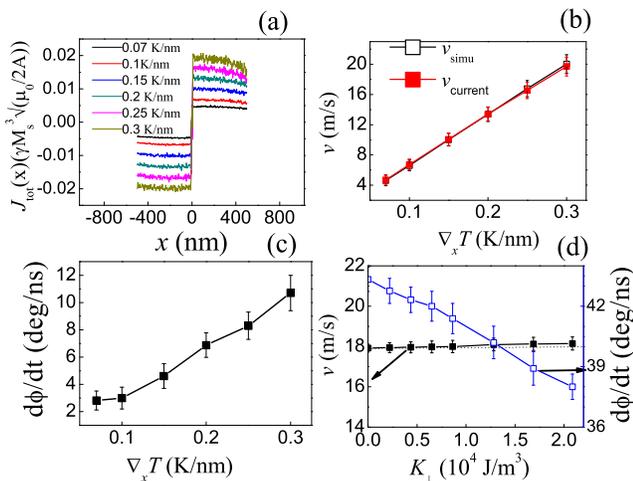


FIG. 2. (a) The spatial dependence of spin current densities $J_{\text{tot}}(x)$ for various $\nabla_x T$. The DW center is chosen as $x = 0$. (b) Thermal gradient dependence of DW velocity v_{simu} from micromagnetic simulations (open squares) and v_{current} computed from total spin current (solid squares). (c) Thermal gradient dependence of DW-plane rotation angular velocity (squares). In (a)(b)(c) model parameters are $L_x = 2048$ nm, $L_y = L_z = 4$ nm, $\alpha = 0.004$ and $K_{\parallel} = 5 \times 10^5$ J/m³. (d) v_{simu} (solid squares) and $d\phi/dt$ (open squares) as a function of K_{\perp} for $L_x = 1024$ nm and $\nabla_x T = 0.5$ K/nm.

Integrating the x -component of Eq. (3) over a space enclosed the DW in the center and noticing the absence of the first term on the right, we have

$$v_{\text{current}} = \frac{1}{2} \int_{-L_x/2}^{L_x/2} \frac{\partial m_x}{\partial t} dx = -\frac{2\gamma A}{\mu_0 M_s} \left[\frac{1}{2} (J_x|_{\text{left}} - J_x|_{\text{right}}) \right]. \quad (5)$$

where we have assumed the fluctuations in the domains are small and the DW is not far from a symmetric one. $J_x|_{\text{left}}$, $J_x|_{\text{right}}$ mean the x -components of the total spin current on the left and right sides of the DW. The equation clearly shows that the DW propagates opposite to the spin current. This is the theoretical DW velocity under the assumption of angular momentum conservation, and it will be compared with the directly simulated DW velocity below.

III. RESULTS

A. Average spin current and DW velocity

To substantiate our assertion that DW propagation under a thermal gradient is through angular-momentum effect instead of energy effect, we would like to compare the DW velocity obtained from micromagnetic simulations and that obtained from total spin current based on

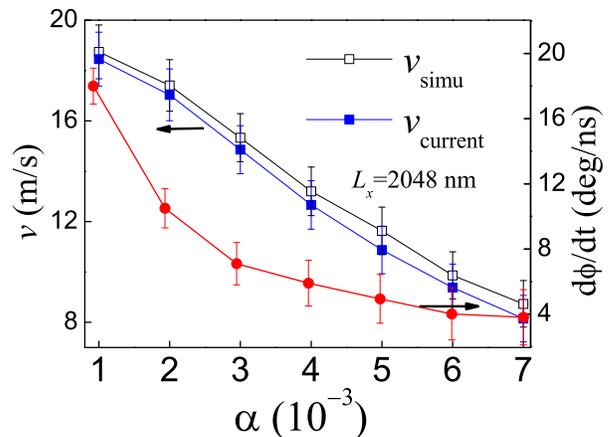


FIG. 3. Damping α dependence of the DW dynamics: v_{simu} (Open squares); v_{current} (solid squares); and $d\phi/dt$ (solid circles). Model parameters are $\nabla_x T = 0.2$ K/nm, $K_{\parallel} = 5 \times 10^5$ J/m³, $L_x = 2048$ nm and $L_y = L_z = 4$ nm.

Eq. (5). Eq. (4) is used to calculate $J_x(x)$. Fig. 2(a) is spatial distribution of the ensemble averaged $J_x(x)$ with DW at $x = 0$ for various thermal gradients. The sudden sign change of $J_x(x)$ at the DW center is a clear evidence of strong angular-momentum transfer from spin current to the DW. Technically, magnetization of the two domains separated by the DW point to the opposite directions, thus the spin current polarization changes its sign. In calculating DW velocity v_{current} from Eq. (5), the spin currents before entering DW and after passing DW are the averages of $J_x(x)$ over $x \in [-2\Delta, -\Delta]$ and $x \in [\Delta, 2\Delta]$, where Δ is the DW width which is 16 nm in the current case. The thermal gradient dependence of v_{current} is shown in Fig. 2(b) (solid squares). v_{current} compares well with the velocity v_{simu} (open squares) obtained directly from simulations by extracting the speed of the DW center along x -direction. The DW velocity is linearly proportional to the temperature gradient $v = C\nabla_x T$, with the thermal mobility $C = 6.66 \times 10^{-8}$ m²s⁻¹K⁻¹ for v_{simu} or $C = 6.59 \times 10^{-8}$ m²s⁻¹K⁻¹ for v_{current} . It is noted that v_{current} almost coincides with v_{simu} except a small discrepancy at very high thermal gradient when the nonlinear effects is strong. The small discrepancy may be attributed to the large fluctuations as well as the contribution from the damping, the dipolar and stochastic fields. These observations are consistent with magnonic STT [15, 16, 25, 26]. It is observed that the DW-plane rotates around the x -axis counter-clockwise for head-to-head DW and clockwise for tail-to-tail DW during DW propagation. DW rotation speed $d\phi/dt$ (squares) is shown in Fig. 2 (c) as a function of $\nabla_x T$.

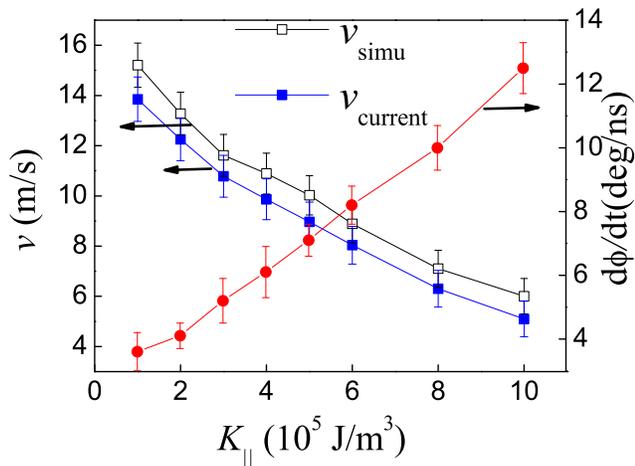


FIG. 4. Anisotropy K_{\parallel} dependence of the DW dynamics: v_{simu} (open squares); v_{current} (solid squares); and $d\phi/dt$ (solid circles). Model parameters are $L_x = 2048$ nm, $L_y = L_z = 4$ nm, $\alpha = 0.004$ and $\nabla_x T = 0.2$ K/nm.

B. Damping and anisotropy dependence of DW dynamics

An energy-effect and angular-momentum-effect have different damping-dependence and anisotropy-dependence of DW dynamics. To distinguish the roles of energy and the angular-momentum in thermal-gradient driven DW dynamics, it would be useful to probe how the DW dynamics depends on α and K_{\parallel} . Damping have two effects on the spin currents: one is the decay of spin current during its propagation so that the amount of spin angular momentum deposited on a DW should decrease with the increase of the damping coefficient. As a result, the DW propagation speed and DW-plane rotation speed should also be smaller for a larger α . Indeed, this is what we observed in our simulations as shown in Fig. 3(a) for DW speed and DW-plane rotation speed (open squares for v_{simu} , solid circles for v_{current} , and stars for $d\phi/dt$). The model parameters are $L_x = 2048$, $L_y = L_z = 4$ nm, $\nabla_x T = 0.2$ K/nm and $K_{\parallel} = 5 \times 10^5$ J/m³. The second damping effect is that the larger α helps the spin current polarization to align with the local spin. This second effect enhances the adiabatic process that is important for non-adiabatic STT or field-like torque discussed in the next subsection. Therefore, α -dependence of DW dynamics supports the origin of thermal driven DW dynamics to be the angular-momentum effect, not the energy effect that would lead to a larger v_{simu} and $d\phi/dt$ for a larger α [3, 4, 33–35] instead of a decrease observed here.

Here we would like to see how the DW dynamics depends on uniaxial anisotropy K_{\parallel} . Fig. 4 shows both v_{simu} (open squares), v_{current} (filled squares) and $d\phi/dt$ (circles) for $L_x = 2048$ nm, $\alpha = 0.004$ and $\nabla_x T = 0.2$. The DW propagation speed, v_{simu} decreases with K_{\parallel} while

DW-plane rotational speed increases with K_{\parallel} . These results seem follow partially the behavior of magnetic-field induced DW motion, in which DW propagation speed is proportional to DW width ($\Delta \sim \sqrt{\frac{A}{K_{\parallel}}}$) or decrease with K_{\parallel} , and partially electric current driven DW motion, in which DW-plane rotational speed increases with K_{\parallel} . Thus, one may tend to conclude that a thermal gradient behaves more like a magnetic field rather than an electric current from the DW width dependence of DW propagation speed, opposite to our claim of the angular-momentum effects of the thermal gradient. It turns out, this is not true. The reason is that magnon spectrum, $\omega_{\mathbf{k}} = \frac{2\gamma}{\mu_0 M_s} (Ak^2 + K_{\parallel})$, has a gap in a system with magnetic anisotropy. The larger K_{\parallel} is, the bigger the energy gap will be. Thus, it becomes harder to thermally excite magnon. As a result, the spin current decreases as K_{\parallel} increases. To see whether the thermal-gradient driven DW motion is due to the angular-momentum transfer or not, one should compare whether v_{simu} and v_{current} maintain a good agreement with each other as K_{\parallel} varies. Indeed, a good agreement between v_{simu} and v_{current} is shown in Fig. 4. This conclusion is also consistent with existing magnonic STT theories [33–35].

C. Separation of adiabatic and non-adiabatic torques

We have already demonstrated that a thermal gradient interacts with DW through magnonic STT rather than through energy dissipation. It is then interesting to know what kind of STTs a thermal gradient can generate. Specifically, whether a magnonic spin current generates damping-like (adiabatic), or field-like (Non-adiabatic) torques, or both just like an electric current

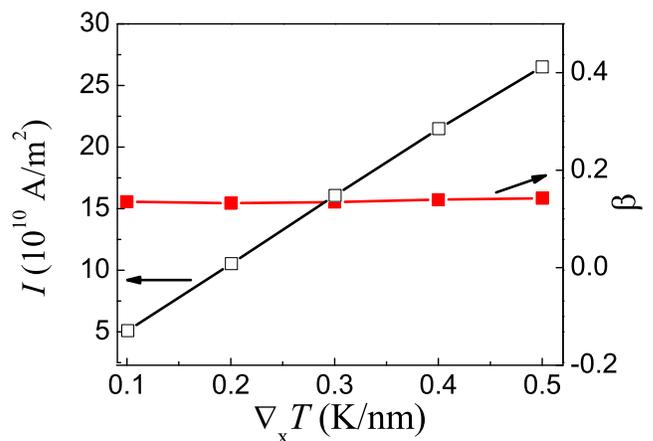


FIG. 5. Model parameters are $K_{\parallel} = 5 \times 10^5$ J/m³, $\alpha = 0.004$, $L_x = 1024$ and $L_y = L_z = 4$ nm. Effective electric current density I (open squares) and β (solid squares) are plotted as functions of $\nabla_x T$.

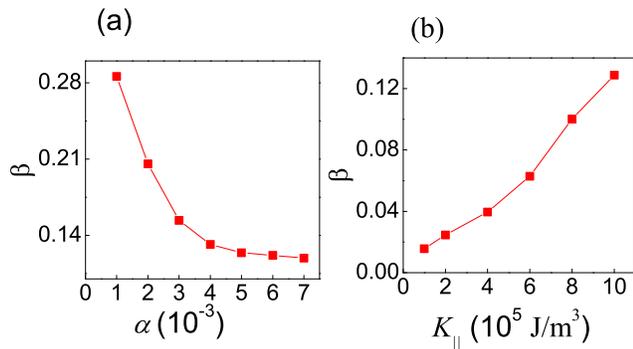


FIG. 6. Model parameters are $\nabla_x T = 0.5$ K/nm, $L_x = 1024$ nm and $L_y = L_z = 4$ nm. (a) α -dependence of β for $K_{\parallel} = 10^6$ and J/m^3 . (b) K_{\parallel} -dependence of β for $\alpha = 0.004$.

[27] does. To extract the STT generated from a thermal gradient, we approximate DW dynamics by the motion of its collective modes of DW center X and the tilted angle ϕ of DW-plane. Subject to both damping-like and field-like torques, using the travelling-wave ansatz [33–35], $\tan(\theta/2) = \exp[(x - X)/\Delta]$ where $\Delta \sim \sqrt{A/K_{\parallel}}$, one can derive the equations for X and ϕ ,

$$\frac{\alpha}{\Delta} \frac{dX}{dt} + \frac{d\phi}{dt} = \frac{\beta}{\alpha} u, \quad \frac{1}{\Delta} \frac{dX}{dt} - \frac{\alpha d\phi}{dt} = \frac{u}{\alpha}. \quad (6)$$

From the above two equations, one can straightforwardly find DW propagating speed and DW-plane rotation speed,

$$v = \frac{(1 + \alpha\beta)}{(1 + \alpha^2)} u, \quad \dot{\phi} = \frac{(\beta - \alpha)}{(1 + \alpha^2)} u. \quad (7)$$

One can extract β and equivalent electric current density $I = (2eM_s u)/g\mu_B P$ from v and $d\phi/dt$ obtained in simulations. For $\alpha = 0.004$, $K_{\parallel} = 10^6 \text{ J/m}^3$, the I and β are obtained and plotted in Fig. 5 as a function of $\nabla_x T$. It is evident that I linearly increases with $\nabla_x T$ and β is independent of $\nabla_x T$ as it should be. We then fixed $\nabla_x T = 0.5$ K/nm, and repeat simulations and analysis mentioned above for various α and K_{\parallel} . Fig. 6 (a) and (b) shows β as a function of α and K_{\parallel} . From the figure, it is evident that β decreases with α . This is because the larger damping favors the alignment of spin current polarization with the local spin so that the non-adiabatic effect, β , becomes smaller. β increases with K_{\parallel} for the similar reason: Larger K_{\parallel} means a thinner DW so that it is much harder for the spin current polarization to reverse its direction after passing through the thinner DW, i.e. a stronger non-adiabatic effect.

In some experiments, the temperature gradient is generated by a laser spot [36]. The laser spot will induce a Gaussian distribution of the temperature over the space [36, 37]. In Fig. 7, we show the DW motion in a Gaussian temperature profile $T(x) = T_0 \exp\left(-\frac{(x-x_L)^2}{2\sigma^2}\right)$ by plotting the DW position against the time. Here we use the

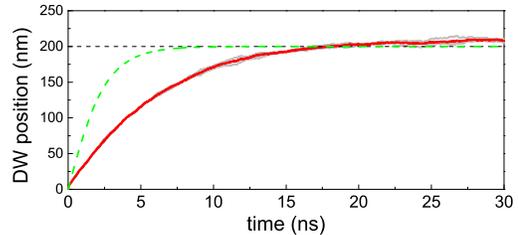


FIG. 7. Domain wall position versus time in a Gaussian temperature profile. The gray lines are raw data for different random seeds and the red line is the averaged result. The green dashed line is theoretical result using the thermal mobility $C = 6.66 \times 10^{-8} \text{ m}^2\text{s}^{-1}\text{K}^{-1}$ obtained from Fig. 2(b).

same parameters as those in Fig. 2(b), except a longer wire $L_x = 2048$ nm, and $T_0 = 400$ K, $\sigma = 200$ nm, $x_L = 200$ nm. Theoretically, if the instantaneous DW speed under a Gaussian temperature is the same as that in the constant thermal-gradient case, we should expect $\frac{dx}{dt} = C \frac{dT}{dx}$, where the thermal mobility C is the same as that in Fig. 2(b). Using $C = 6.66 \times 10^{-8} \text{ m}^2\text{s}^{-1}\text{K}^{-1}$, the above differential equation for $x(t)$ can be numerically solved with initial condition $x(0) = 0$. The result is plotted in Fig. 7 in green dashed line. The simulated speed is smaller than this theoretical result. This is probably because, for the constant thermal-gradient, we focus on the steady-state DW motion speed. In a Gaussian temperature, the DW cannot immediately follow the local temperature gradient. Before the DW can reach the steady-state speed corresponding to the local temperature, it already moves to a position of smaller temperature gradient. More details about DW motion in Gaussian temperature profile may be an issue of future studies.

IV. DISCUSSION AND SUMMARY

We have studied the thermal gradient-driven DW dynamics in an uniaxial nanowire. In reality, there is always certain hard anisotropy in a wire whose cross-section is not a perfect ellipse. Thus, it is interesting to see how the above results will change in a weak biaxial nanowire with a small hard anisotropy $K_{\perp} = 1/2\mu_0 M_s^2 (\mathcal{N}_z - \mathcal{N}_y)$, say along y -direction. Our simulations show that a DW still propagates towards the higher temperature region in a similar way as that in a uniaxial wire. Interestingly, as shown in Fig. 2(d) for the K_{\perp} -dependence of v_{simu} (solid squares) and $d\phi/dt$ (open squares), DW speed increases

slightly with K_{\perp} . This may be due to the increase of torque along θ -direction [33] since Γ_{θ} is proportional to $(\mathcal{N}_z - \mathcal{N}_y)$. This is also consistent with the early results for the uniaxial wire that v_{simu} (which includes stochastic thermal field and demagnetisation fields) is always larger than v_{current} (where the transverse fields are neglected). At the meanwhile, $d\phi/dt$ decreases with K_{\perp} .

The main purpose of this paper is to study the magnonic effects in thermal-gradient-driven domain wall dynamics. We consider the spin waves explicitly and all the material parameters (exchange constant A , crystalline anisotropy K , saturation magnetization M_s , and Gilbert damping α) are assumed to be constant. Indeed, the atomistic magnetic moments are independent of temperature. At the atomistic level, the exchange constant A originating from the Pauli exclusion principle and the crystalline anisotropy K originating from the spin-orbit coupling only weakly depend on the temperature because of the vibration of atoms [39]. In micromagnetic models, because finite volumes that contains many magnetic moments are considered as unit cells, the parameters A , K , and M_s depend on the temperature. This is because the thermally excited spin waves with wavelengths shorter than the length scale of the unit cells are included in the effective A , K , and M_s by doing an average [16, 38]. Since we use small mesh size $2 \times 2 \times 2 \text{ nm}^3$, only spin waves of very short wavelength affect the parameters A , K , and M_s in our model. Those short-wavelength spin waves possess high energy as well as low density of states, so their contributions to the effective A , K , and M_s are not significant. The Gilbert damping α depends on the temperature non-monotonically [40–43]. The underlying mechanism is still under debate, but for many cases the dependence is not significant in a wide range of temperature.

In summary, our results show that the uniform thermal gradient always drives a DW propagating towards the hotter region and the DW-plane rotates around the easy axis. The DW velocity and DW-plane rotational speed decrease with the damping coefficient. The DW velocity obtained from simulation agrees with the velocity obtained from angular momentum conservation when the magnon current density ($J(x)$) from the simulation is used to estimate the amount of angular momentum transferred from magnon current to the DW. All the above findings lead to the conclusion that the thermal gradient interacts with DW through angular-momentum transfer rather than through energy dissipation. Furthermore, we demonstrated that the magnonic STT generated by a thermal gradient has both damping-like and field-like components. The field-like STT coefficient β is determined from DW speed and DW-plane rotation speed. β does not depend on the thermal gradient as expected, but increases with a decrease of DW width. This behavior can be understood from the expected strong misalignment of magnon spin polarization and the local spin so that non-adiabatic torque (also called field-like torque) is larger. For the same reason, a larger Gilbert damping

results in a better alignment between spin current polarization and the local spin, thus β should decrease with α . The thermal gradient can be a very interesting control knob for nano spintronics devices, especially those made from magnetic insulators.

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