

### **3D patient-specific numerical modeling of the soft palate considering adhesion from the tongue**

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# 3D patient-specific numerical modeling of the soft palate considering adhesion from the tongue

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## Abstract

Collapse of the soft palate in the upper airway causes obstructive sleeping problem. A comprehensive numerical simulation of the soft palate contributes to providing useful information to the clinical study. The adhesion effect from the tongue still remains to be investigated, and no cohesive simulation for the surface tension has been presented. In this study, the traction-separation cohesive approach was addressed to describe the adhesion effect from the surface tension of the lining liquid between the soft palate and the tongue. According to pull-off experimental tests of human lining liquid from the oral surface of the soft palate, the corresponding cohesive properties, including the critical normal traction stress and the failure separation displacement, were obtained. The 3D patient-specific soft palate geometry was accounted for, based on one specific patient's computed tomography (CT) images. The calculation results showed that influence of the adhesion from the tongue surface on the global response of the soft palate depends on the length ratio between the cohesive length and the soft palate length. When the length of the cohesive zone was smaller than half of the soft palate length, the adhesion's influence was negligible. When the adhesion length was larger than 70 percent of soft palate length, the adhesion force contributes to preventing the soft palate from collapsing towards to the pharynx wall, i.e. the closing pressure was more negative than in the no adhesion case. These results may provide useful information to the clinical treatment of OSA patients.

*Keywords:* Cohesive approach, surface tension, patient-specific modeling, numerical simulation, obstructive sleep apnea

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## 1. Introduction

For obstructive sleep apnea (OSA) patients, the collapse of the soft palate in the upper airway is one of the main causes of breathing stops. Typically, when patients sleep with their mouths closed, the collapse of the soft palate towards the pharynx wall will block normal airflow in the upper airway (Fig .1). The direct reason for the collapse of the soft palate is attributed to the negative pressure that may be caused by the narrowing upper airway. Biomechanical behavior of the soft palate in response to the negative pressure including the collapse will, in turn, have an influence on the pressure distribution of the upper airway and the air flow velocity. Hence,

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8 a comprehensive biomechanics study of the soft palate in the upper airway contributes to the clinical research  
9 of OSA and provides input to improved understanding of the phenomenon.

10 Numerical modeling of the soft palate has been addressed to investigate the soft palate's response to the  
11 airway pressure. [Berry et al. \(1999\)](#) presented an approximate 2-dimensional (2D) cantilever model of the  
12 soft palate and the collapse shape of the soft palate was obtained. [Malhotra et al. \(2002\)](#) employed a 2D  
13 planar model to investigate the closing pressure of the soft palate. In their finite element (FE) model, based  
14 on clinical results, a fitted Young's modulus value of soft palate was obtained. [Sun et al. \(2007\)](#) presented the  
15 movement of soft palate during breathing with a simplified 3D model. [Wang et al. \(2012\)](#) presented a fluid-solid  
16 interaction numerical simulation for the upper airway, including the soft palate. [Inouye et al. \(2015\)](#) showed a  
17 computational model for the soft palate closure to simulate the cleft palate repair. These research works showed  
18 a basic method for numerical modeling of the soft palate. On the other hand, according to Fig. 1, contact  
19 between the soft palate and the tongue is observed. The surface tension of the upper airway mucosal lining  
20 liquid was shown to produce an influence on the upper airway collapsibility ([Kirkness et al., 2003](#)). The surface  
21 tension generated by the mucosal lining liquid between the soft palate and the tongue produces an adhesion  
22 force to prevent the soft palate debonding from the tongue. To the authors' knowledge, investigation of this  
23 surface tension has not been addressed in the literature. Therefore, a numerical method was developed to assess  
24 the adhesion's influence on the soft palate's behavior. In addition, the adhesion investigation also can contribute  
25 to evaluating the influence of the upper airway humidity (especially the dry throat) on the global response of the  
26 soft palate.

27 Surface tension is due to the elasticity of a fluid surface obtaining the least surface area. Modeling of the  
28 surface tension for fluids has been presented in the literature, such as a continuum method ([Brackbill et al.,  
29 1992](#)) and additive-nonadditive modeling method ([Van Oss et al., 1988](#)). When a solid body, such as a soft  
30 tissue, is in contact with a liquid surface, the adhesion is a result of the corresponding positive adhesion energy  
31 ([Jensen et al., 2015](#)). For the liquid, the surface energy is often described to be equal to the surface tension  
32 ([Shen et al., 2015](#)). Studies of the adhesion between soft materials and liquids are mainly described with  
33 surface indentation models, using a spherical indenter and a flat substrate ([Hui et al., 2015](#); [Watson et al., 2015](#);  
34 [Chakrabarti and Chaudhury, 2013](#)). In addition, contact mechanics theories such as Johnson-Kendall-Roberts  
35 (JKR) ([Johnson et al., 1971](#)) and Derjaguin-Muller-Toporov (DMT) ([Derjaguin et al., 1994](#)) have been used to  
36 describe the relation between the external separation force and the adhesion energy, which corresponds to the  
37 surface tension ([Xu et al., 2014](#); [Jensen et al., 2015](#); [Rasuli et al., 2010](#)).

38 On the other hand, a cohesive zone approach combined with the finite element method has been presented  
39 in the simulations to solve fracture mechanics problems, such as the needle insertion ([Oldfield et al., 2013](#)) and  
40 failure of the brittle rocks ([Gui et al., 2016](#)). In addition, the cohesive models have been applied to various  
41 biomechanical problems, such as simulation of cell-matrix adhesion described by [Cóndor and García-Aznar  
42 \(2017\)](#), arterial dissection presented by [Gasser and Holzapfel \(2006\)](#) and [Noble et al. \(2017\)](#), soft material  
43 tearing provided by [Bhattacharjee et al. \(2013\)](#).

44 In order to address the adhesion between the soft palate and the tongue due to the lining liquid in the numer-  
 45 ical simulation, we used the cohesive approach to simulate the surface tension, which has not been presented in  
 46 the literature. An OSA patient's CT images were used to reconstruct the soft palate geometry. Corresponding  
 47 cohesive properties was obtained based on experimental data. Moreover, investigation of different cohesive  
 48 properties of the lining liquid may provide improved understanding of how the humidity of the upper airway  
 49 influences the soft palate collapse and how to improve treatment for OSA patients.

## 50 2. Methods

### 51 2.1. Cohesive approach

52 In the FE code ABAQUS, adhesion modeling of the bonded bodies can be achieved with cohesive elements  
 53 and a traction-separation model for a specifically defined cohesive layer. The relationship between the traction  
 54 and the separation is presented as Fig. 2. Assuming that the normal direction (debonding direction) of the  
 55 cohesive layer is direction 1, and directions 2 and 3 denote the orthogonal in-plane directions, the nominal  
 56 traction stress  $\mathbf{t}$  is described as  $[t_i] = [t_1, t_2, t_3]^T$ . The initial constitutive thickness of the cohesive layer is  
 57 denoted by  $T_0$ . According to the separation displacement  $\delta$ , the nominal strain is calculated as  $\epsilon = \delta/T_0$ . In  
 58 ABAQUS, the default value for  $T_0$  is set to be 1, this ensures that the nominal strain is equal to the separation.  
 59 The detailed description is reported in ABAQUS analysis user's guide (ABAQUS, 2014). When the separation  
 60 displacement  $\delta$  increases to the damage initiation value  $\delta_i$  (see Fig. 2), the traction stress reaches the maximum  
 61 value  $t_c$ , and the degradation process begins. In the degradation process, the traction stress decreases gradually  
 62 to 0 as  $\delta$  increases towards  $\delta_f$ , which is defined as the separation at failure.

63 In this study, we mainly considered the normal direction's mechanical behavior, and a linear elastic stiffness  
 64  $E$  was used until the traction stress reaches the critical level (the initial stage). The elastic behavior is described  
 65 as

$$t_1 = E\epsilon_{11} = E\delta, \quad (1)$$

66 where  $\delta$  is the separation in the normal direction of the cohesive layer. Additionally, in the degradation  
 67 process, a scalar damage variable  $D$  is used to define the traction stress within ABAQUS (ABAQUS, 2014)

$$t_1 = (1 - D)E\delta, \quad (2)$$

68 where  $D$  has an initial value 0 and increases to 1 in the end of the degradation process (failure). In this  
 69 study, the linear degradation process was considered (see Fig. 2).

### 70 2.2. Pull-off test of upper airway lining liquid

71 In the cohesive approach, two important values are needed when investigating the adhesion of the mucosal  
 72 lining liquid between the soft palate and tongue. One is the normal direction's critical traction stress  $t_c$  defined

73 as the critical value for damage initiation, and the other one is the failure separation displacement  $\delta_f$  in the  
 74 normal direction of the cohesive layer. Based on the surface tension test of the upper airway lining liquid  
 75 provided by Kirkness et al. (2005), we can determine these two values. In their experimental test, the lining  
 76 liquid's thickness was initially equal to 21  $\mu\text{m}$ . This is close to the normal human lining liquid's thickness of  
 77 26.4  $\mu\text{m}$  (Lee et al., 2002) on the soft palate's oral surface. In addition, the pull-off force separating the adhesion  
 78 was calculated in the experimental test as

$$F = 2\pi R\gamma(2\cos\theta + \sin(\theta + \phi)), \quad (3)$$

79 where  $\gamma$  is the surface tension, which was tested to be 56.1  $\text{mNm}^{-1}$  for the lining liquid between the soft  
 80 palate and the tongue.  $R$ ,  $\theta$  and  $\phi$  are geometry constants measured in the experimental test, see Fig. 3. In their  
 81 pull-off test, two cylindrical silica discs were stacked together with the axes perpendicular to each other. The  
 82 lining liquid from the upper airway was coated on the disc's surface and generated an annulus, which can be  
 83 shown as in Fig. 3. The detailed description is provided in Kirkness et al. (2005). The values for the geometrical  
 84 constants are:  $R = 20$  mm and  $\theta = \phi = 5^\circ$ . Based on the obtained surface tension and Eq. (3), the pull-off  
 85 force can be calculated to be 15.27 mN.

86 Based on the obtained pull-off force, we calculated the critical traction stress  $t_c$  (Fig. 2) in the normal  
 87 direction as

$$t_c = \frac{F}{A_0}, \quad (4)$$

88 where  $A_0$  is the contact area of the tested lining liquid surface and  $F$  is the obtained pull-off force. As  
 89 shown in Fig. 3, the contact area of the lining liquid in the test is calculated to be  $A_0 = \pi r^2$ . Since the radius  
 90  $R$  was set to be 20 mm and the angle  $\phi$  was measured to be  $5^\circ$  (Kirkness et al., 2005), the radius of the contact  
 91 zone  $r$  in Fig. 3 is 1.74 mm. Then, the contact area  $A_0$  in Eq. (4) was obtained as 9.51  $\text{mm}^2$ , and the critical  
 92 traction stress was calculated as 1.61 kPa.

93 On the other hand, in the experimental pull-off test, a jump apart of the two discs stacked together with the  
 94 lining liquid was observed. This jump displacement was named  $D$ . Relating the stiffness  $k$  of the spring that  
 95 was used in the experimental test to the pull-off force, the jump apart displacement was calculated as:

$$D = F/k. \quad (5)$$

96 The detailed description was reported in Kirkness et al. (2005). The jump apart displacement  $D$  was calcu-  
 97 lated to be 0.57 mm, based on Eq. (5). In this study, we assumed the failure separation  $\delta_f$  is equal to this jump  
 98 apart displacement observed in the experimental test.

99 Hence, the two key values for the cohesive approach (the critical traction stress  $t_c$  and the failure separation  
 100 displacement  $\delta_f$ ) are obtained. Since the separation evolves in the normal direction, for the cohesive approach,  
 101 the failure separation displacement  $\delta_f$  is only needed to be set in the normal direction.

102 Additionally, we assumed the human lining liquid has an isotropic cohesive property. Then, we used the  
103 same elastic stiffness  $E$  as in Eq. (1) for all three directions in the simulation. Based on the obtained critical  
104 traction stress  $t_c$  and the failure separation  $\delta_f$  in the normal direction, we only needed to determine the initial  
105 linear elastic stiffness  $E$ , which is related to the damage initial separation  $\delta_i$ .  $\delta_i$  has a limited influence on the  
106 global response of the separation of two bonded bodies, such as the soft palate and the tongue. The strain energy  
107 represented as the area underneath the curve shown in Fig. 2 will be the same for different choices of  $\delta_i$ . We set  
108  $\delta_i$  to be 0.01 mm in this study to address the adhesion simulation of the soft palate.

109 In order to ensure our numerical model gives the same pull-off force as in the experimental test, a simplified  
110 FE model was created to mimic the pull-off test of the human upper airway mucosal lining liquid. According  
111 to the numerical simulation result, the obtained pull-off pressure corresponds with the experimental test value  
112 1.61 kPa.

113 In addition, the influence of  $E$  on the calculation of the critical traction pressure was checked, showing that  
114 it was very limited. We tested two more cases:  $\delta_i = 0.1$  and  $\delta_i = 0.2$ . The obtained critical pull-off pressure is  
115 the same as for  $\delta_i = 0.01$ . Moreover, since the normal human lining liquid thickness on the oral surface of the  
116 soft palate was measured to be 26.4  $\mu\text{m}$  (Lee et al., 2002), we set the thickness of the lining liquid in the soft  
117 palate simulation to be 26.4  $\mu\text{m}$ . As the test results for different thicknesses showed no difference (Kirkness  
118 et al., 2005), the obtained cohesive properties in the description above were used in the following simulations  
119 for the soft palate.

### 120 2.3. Numerical modeling for the soft palate and tongue

#### 121 2.3.1. 3D patient-specific geometry

122 Recently, computational 3D modeling based on CT or MR images has been applied successfully for studies  
123 of the upper airway (Sera et al., 2015; Sung et al., 2006). In the current study, based on one specific patient's CT  
124 images, we obtained the 3D geometry in ABAQUS. Usage of the images was approved by the Norwegian Re-  
125 gional Committee for Medical Research Ethics (REK) and was registered in Clinicaltrials.gov. (NCT01282125).  
126 The detailed description for the geometry modeling has been presented in our previous work Liu et al. (2018).  
127 In addition, we simplified the boundary between the tongue and the soft palate as a straight line, which con-  
128 tributes to obtaining an efficient numerical cohesive simulation. Moreover, to simplify the simulation, as shown  
129 in Fig. 4, the whole tongue was not reconstructed. The cohesive part of the tongue was modeled as a 20 mm  
130 thickness brick and the length of the cohesive zone was measured to be 17 mm based on the specific patient's  
131 CT image. The thickness of the cohesive layer was set to be 26.4  $\mu\text{m}$  (Lee et al., 2002). Additionally, since the  
132 patient-specific CT images was scanned when the patient was lying down with the supine position, the influence  
133 of the gravity of the tongue was considered in the remaining numerical simulation.

134 *2.3.2. Material properties and boundary conditions*

135 A neo-Hookean hyperelastic model was employed. The strain-energy function reads:

$$\Psi(\bar{I}_1, J) = \underbrace{c(\bar{I}_1 - 3)}_{\text{isochoric}} + \underbrace{\frac{1}{D_1}(J - 1)^2}_{\text{volumetric}}. \quad (6)$$

136 Here,  $c$ ,  $D_1$  are material parameters derived from the Young's modulus  $E$  and Poisson's ratio  $\nu$  with the follow-  
137 ing relations (Berry et al., 1999):

$$c = \frac{E}{4(1 + \nu)}, \quad D_1 = \frac{6(1 - 2\nu)}{E}. \quad (7)$$

138 The value of the Young's modulus for the soft palate was determined according to an in vivo magnetic  
139 resonance elastography measurement of the soft palate provided by Cheng et al. (2011). Assuming the Poisson's  
140 ratio value to be 0.49, the Young's modulus was calculated to be 7.539 kPa, based on the measured shear  
141 modulus 2.53 kPa. In addition, the neo-Hookean model was also used to define the material property of the  
142 tongue. According to the in vivo measured shear modulus 2.67 kPa by Cheng et al. (2011), the Young's modulus  
143 of the tongue was calculated to be 7.957 kPa with a Poisson's ration of 0.49. Moreover, for the cohesive layer,  
144 the cohesive properties in Section 2.2 were used, including the obtained critical traction stress  $t_c = 1.61\text{kPa}$   
145 and the failure separation  $\delta_f = 0.57\text{mm}$ .

146 The boundary conditions for the soft palate were set as in Fig. 5, and the bottom of the tongue shown in  
147 Fig. 4 was also constrained. Finally, since the collapse of the soft palate in the negative pressure field is one  
148 of the main reasons to cause OSA, we applied a uniformly distributed pressure, corresponding to the pressure  
149 difference between the anterior and posterior sides of the soft palate (Fig. 4), on the surface where the soft  
150 palate contacts with the airway (Fig. 5).

$$P_{\text{negative}} = P_{\text{posterior}} - P_{\text{anterior}}. \quad (8)$$

151 When the negative pressure develops, the soft palate will have a posterior oblique deformation. If the  
152 negative pressure is large enough, the soft palate will stick to the pharynx wall. Then, the OSA occurs. We call  
153 this critical negative pressure the closing pressure. However, the adhesion force from the mucosal lining liquid  
154 between the soft palate and the tongue will act against this posterior oblique deformation of the soft palate.  
155 Therefore, the investigation of the global response of the soft palate considering the adhesion effect contributes  
156 to obtaining physiological simulation results.

157 *2.3.3. Mesh convergence study*

158 The 3D patient-specific soft palate model was meshed with four noded hybrid tetrahedral elements (C3D4H  
159 ABAQUS type). A -5 cm H2O negative pressure was applied and the neo-Hookean model with the aforemen-  
160 tioned property data of the soft palate was assigned to the 3D patient-specific model. As shown in Fig. 6, four  
161 mesh densities were tested with 139 338, 397 716, 651 742 and 852 870 elements, corresponding to Mesh 1,

162 Mesh 2, Mesh 3 and Mesh 4, respectively. The difference for the critical parameter between Mesh 3 and Mesh 4  
163 was 0.7% (Fig. 6). Therefore, considering the simulation accuracy and computational time efficiency, we used  
164 a mesh with the size of Mesh 3 in the remaining simulation of this paper. Meanwhile, the cohesive element with  
165 the same mesh size was assigned to the cohesive layer, and the tongue was also meshed with the size of Mesh 3  
166 and the eight noded hybrid elements (C3D8H ABAQUS type) were addressed.

### 167 3. Results

168 The collapse of the soft palate in the upper airway and the failure of the adhesion with respect to the cohesive  
169 layer could be observed directly with the 3D patient-specific geometry as shown in Fig. 6. Additionally, in order  
170 to present the global response of the soft palate, we used the norm of the displacement of point A to quantify the  
171 inclination displacement of the posterior surface of the soft palate, see Fig. 7. When point A makes contact with  
172 the pharynx wall, the corresponding negative pressure was set to be the closing pressure, which was regarded as  
173 a critical parameter to evaluate the global response of the soft palate. The detailed calculation results are shown  
174 as Fig. 8. Based on the simulation results, for this specific patient, we found that the influence of adhesion from  
175 the lining liquid between the soft palate and tongue on the global response of the soft palate is negligible. The  
176 closing pressure, considering the adhesion effect, was calculated to be -5.56 cm H<sub>2</sub>O. In the no adhesion case,  
177 the closing pressure was -5.54 cm H<sub>2</sub>O. This is reasonable, considering the length of the cohesive layer was  
178 measured to be only half the length of the soft palate, while point A locates at the distal part of the soft palate.

179 In order to have a more comprehensive study of the adhesion effect, we investigated 18 patients' CT images  
180 and found different contact types. We divided them into 4 adhesion types (Table 1), based on the ratio between  
181 the adhesion length and the soft palate length, including below 50%, 50%-70%, above 70% and tip adhesion  
182 types. For simplicity, we did not create all these patients' 3D geometry model. Instead, we extended the length  
183 of the cohesive layer (shown in Fig. 4) in the created patient-specific model. The length of the tongue model  
184 was also extended following the extension of the cohesive layer. In addition, for the tip adhesion case, we  
185 only modeled the tip part's contact, and the anterior part's contact between the soft palate and the tongue was  
186 neglected, because for a short contact length the adhesion influence is very limited, according to Fig. 8 (A).  
187 The global responses of soft palate versus the negative pressure for each contact type are shown in Fig. 8.  
188 The closing pressure for each case are reported in Table 1. According to the simulation results, the adhesion  
189 effect is observed to be strengthened following the increase of the adhesion length. Failure of the adhesion (see  
190 Fig. 7) occurs when the adhesion length is larger than 70 percent of the soft palate. Additionally, when the  
191 adhesion failure happens, a step-increase of the norm of displacement of point A is presented, corresponding  
192 to the decrease of the traction stress during the degradation process (Fig. 2). Moreover, the closing pressure is  
193 also observed to become more negative following the increase of the adhesion length. We also used the created  
194 patient-specific 3D geometry for the simulation of the tip adhesion case.

195 We used the critical traction stress 1.61 kPa in the above simulations, based on the pull-off experiment test.  
196 The critical traction stress relates directly to the adhesion strength. This motivates us to investigate another



197 case for evaluating the influence of changing the critical traction stress on the global response of the soft palate.  
198 A smaller critical traction stress of magnitude 0.5 kPa was considered. The adhesion length was set to be 70  
199 percent of the soft palate. Using the same initial separation displacement for damage ( $\delta_i = 0.01$ ), the initial  
200 linear elastic stiffness was calculated to be 50 kPa. The comparison between these two kinds of critical traction  
201 stresses is shown in Fig. 9. According to the calculation results, different global responses of soft palate evolve,  
202 and the closing pressure for the smaller critical traction stress case is less negative than the larger one. Moreover,  
203 it is obvious that the adhesion failure starts earlier for the smaller critical traction stress case. This shows that  
204 the critical traction stress has a significant influence on the adhesive behavior.

#### 205 4. Discussion and conclusion

206 The adhesion from the mucosal lining liquid on the soft tissue's surface has an influence on the biomechanical  
207 behavior of the corresponding soft tissue. This is presented in the above simulations for the deformation of  
208 the soft palate, accounting for the adhesion due to the lining liquid between the soft palate and tongue. Based  
209 on the experimental pull-off test of the human lining liquid from the oral surface of soft palate (Kirkness et al.,  
210 2005), using the traction-separation cohesive approach, a 3D patient-specific numerical simulation considering  
211 the adhesion effect was achieved. The simulation results show that the adhesion strengthening effect depends  
212 on the contact area between the soft palate and the tongue. When the contact length is smaller than half of  
213 the soft palate length, the adhesion's influence on the global response of the soft palate was negligible. For  
214 a contact length larger than 70 percent of the length of the soft palate, the adhesion leads to a more negative  
215 closing pressure for the soft palate, but this improvement rate is moderate. Kirkness et al. (2003) showed that  
216 decreasing the surface tension of the upper airway lining liquid contributes to reducing the upper airway col-  
217 lapsibility. Their research focused on the collapse of the whole upper airway, which can be treated as a tube.  
218 Our study investigated the adhesion between the soft palate and the tongue, and the surface tension of the lining  
219 liquid on the posterior surface of the soft palate was not accounted in the numerical simulation. Therefore,  
220 we conclude that the adhesion between the soft palate and the tongue prevents the soft palate from collapse.  
221 A further investigation of the pharyngeal lining liquid's surface tension working on the soft palate's posterior  
222 surface remains to be addressed.

223 The strength of the adhesion depends mainly on the critical traction stress in the normal direction. We  
224 obtained this value from the pull-off experiment test. Some patients may suffer from dry throat (lining liquid  
225 humidity change), and the adhesion from the tongue may be weaker than for healthy people. We investigated  
226 a weaker adhesion case where the critical traction stress was set to be 0.5 kPa, which is lower than the healthy  
227 people's 1.61 kPa. The comparison result showed that the closing pressure is less negative for the weaker  
228 adhesion case. Whether the humidity degree of the upper airway has a direct influence on the critical traction  
229 stress of the lining liquid remains to be validated in the further work. Our simulation results point out that  
230 increasing the adhesion force from tongue can make the closing pressure of soft palate more negative and may  
231 improve the obstructive sleeping condition for OSA patients. Moreover, we used the failure separation with

232 0.57 mm. In the real case, small variation of this value may be found. We also investigated another case with  
233 failure separation of 0.5 mm. The results showed a small difference between these two cases' calculations. This  
234 indicates that the variation of the failure separation  $\delta_f$  in a small range has a limited influence on the global  
235 response of the soft palate.

236 In our simulation, only the passive condition was considered, and the neuromuscular activation effect from  
237 palatal muscle fibers was neglected, considering the patients' activation level is defective (McGinley et al.,  
238 2008). Moreover, in order to simplify the calculation, we applied the negative pressure as a uniformly distributed  
239 load. However, this is not the case in reality. Therefore, fluid-structure interaction analysis may be employed  
240 in order to predict a more realistic pressure distribution when the large deformation problem is considered.  
241 This will be a task in the further studies of the soft palate. In addition, we defined the closing pressure as the  
242 negative pressure that make the point A contacts the pharynx wall. This is an approximate definition since the  
243 soft palate is not fully collapsed. We used this approximate method to get a consistence result when comparing  
244 different simulation cases. In addition, we simplified the geometry between the soft palate and the tongue as a  
245 plane, which contributes to efficient cohesive simulation. Moreover, we simplified the geometry of the tongue  
246 to be a block body. In the further study, more physiological geometry model is needed to achieve more realistic  
247 simulation results. Additionally, for different cohesive length cases' simulation, we used only one patient-  
248 specific geometry. The patient-specific model for each case will be created in our next step work that achieves  
249 more physiological simulation results.

250 Based on the pull-off test of the human lining liquid, using the traction-separation cohesive approach, we  
251 provide a method for numerical modeling of the soft palate, considering the adhesion effect from the mucosal  
252 lining liquid. The adhesion effect is evaluated to have an influence on the global response of soft palate, but  
253 this influence in terms of the change of closing pressure is moderate. The cohesive simulation gives input to  
254 the clinical research on how the adhesion effect of the upper airway lining liquid influences the biomechanical  
255 behavior of the soft palate. For example, keeping the optimal humidity in the the OSA patients' throat may be  
256 considered as a method to reduce the soft palate collapse tendency and improve the patients' sleep condition.

### **Conflict of interest statement**

The authors declare that they have no conflict of interest.

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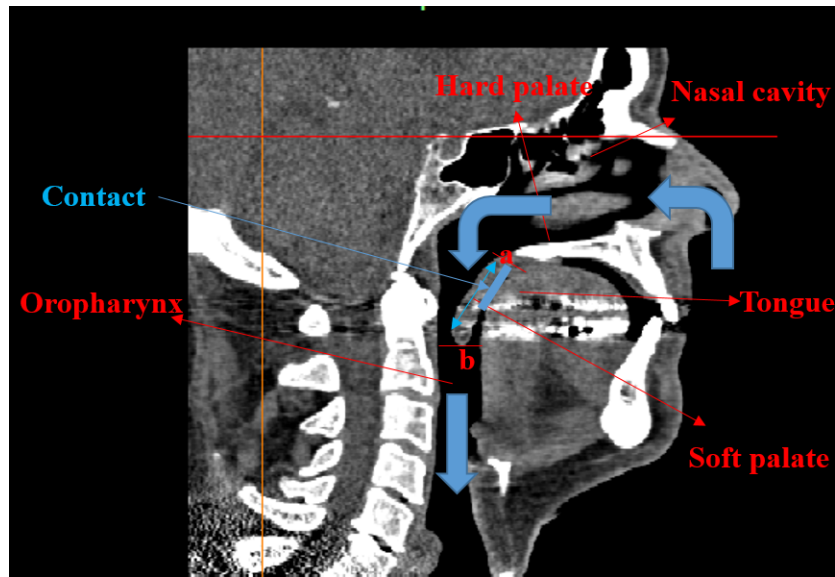
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## Table and figure captions

**Table 1.** The closing pressure for different length ratio cases and the patients' distribution for these cases.

Length ratio	No adhesion	Below 50%	50% – 70%	Above 70%	Tip adhesion
Closing pressure ( $-cm H_2O$ )	5.54	5.54 – 5.56	5.56 – 6.16	6.16 – 7.44	5.96
Number of patients	0	4	4	6	4



**Fig. 1.** CT image of the upper airway in the sagittal midsection. In the breathing process with mouth closed, the air flow comes into the lung through the nasal cavity, pharynx. The soft palate locates in the top part of the pharynx, connecting to the hard palate. We use “a” and “b” markers to have a clear description of the soft palate’s position in the following. The “a” section means the side between the soft palate and the hard palate, the “b” section means the bottom edge of the soft palate tip. In addition, the contact between the soft palate and tongue is marked (blue short line).

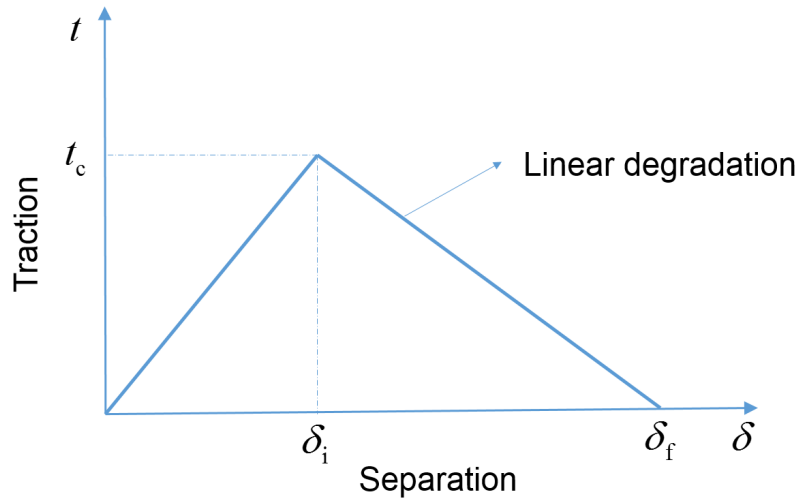


Fig. 2. Traction-separation model for the cohesive approach, the damage initial separation and failure separation are presented.

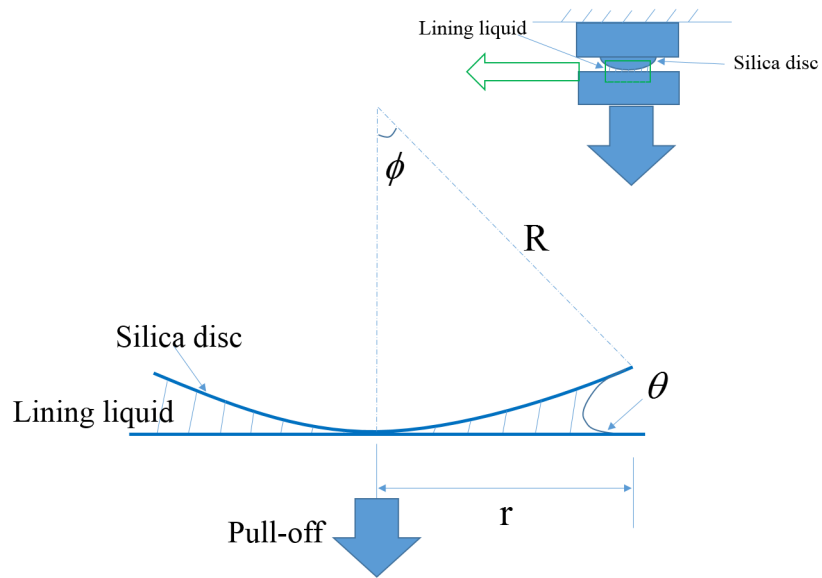
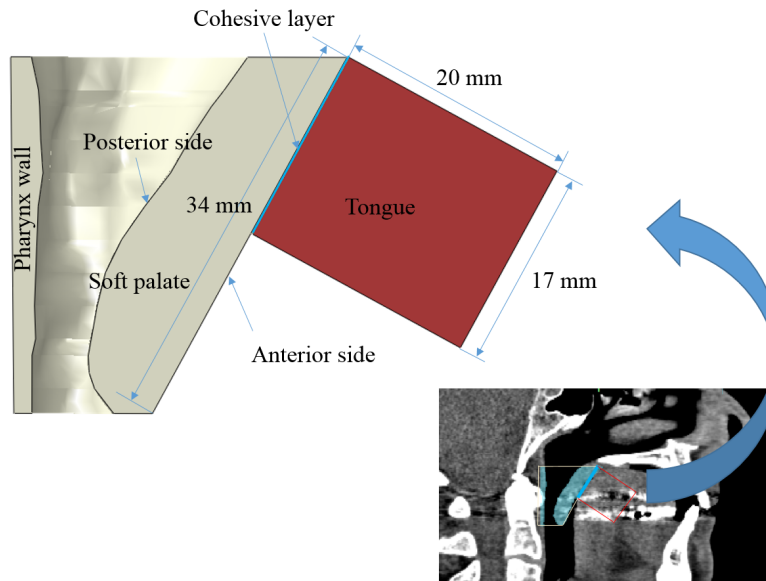
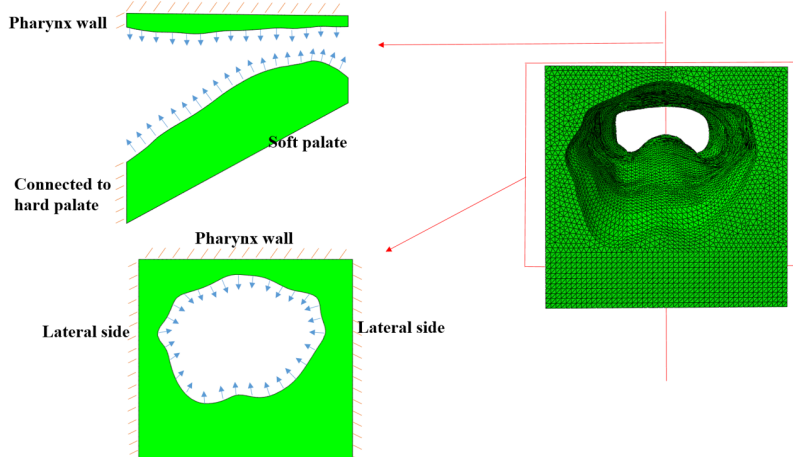


Fig. 3. The schematic of the pull-off test of human upper airway lining liquid according to Kirkness et al. (2005). The top disc is constrained and the lower disc is pulled down with a spring that can measure the pull-off force. The angle  $\theta$  represents the contact angle between the liquid and the lower disc and angle  $\phi$  is an angular measurement of the size of the liquid annulus, which can be used to calculate the contact area between the disc and lining liquid.

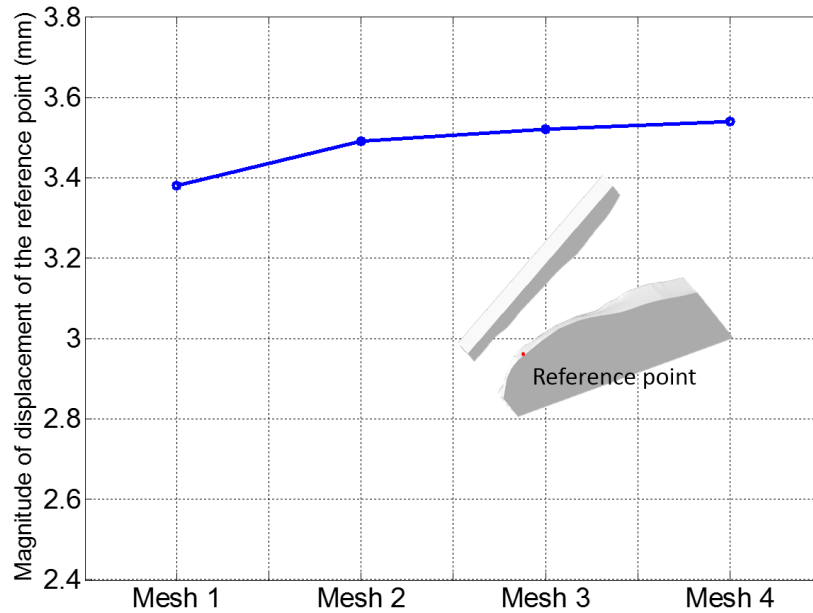


**Fig. 4.** Obtained patient-specific geometry model that was used to investigate the influence of the adhesion from the lining liquid between the soft palate and tongue on the global response of the soft palate in the sagittal midsection view. The length of the soft palate is 34 mm and the contact length between the soft palate and tongue is 17 mm.

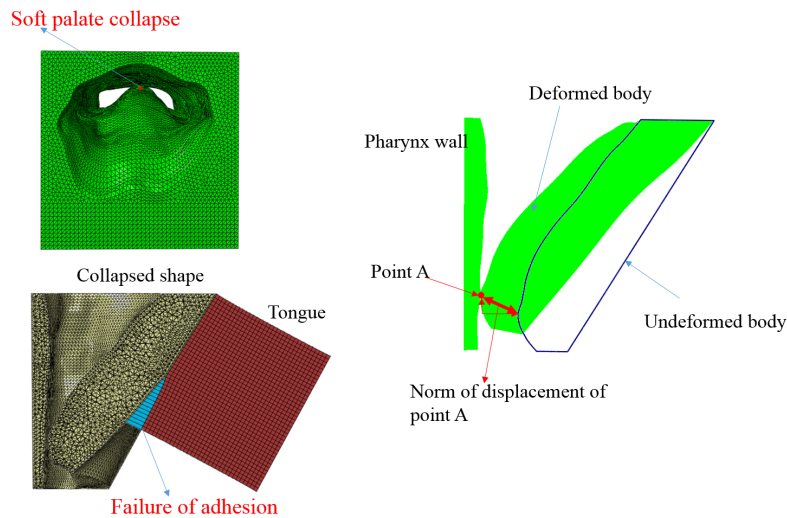


**Fig. 5.** Boundary conditions for the 3D patient-specific model: the external surface of the pharynx wall, the lateral sides and the side connected to the hard palate were constrained in all directions. The left-top figure is the in the view from right to left and the left-bottom figure is in the view looking inside.

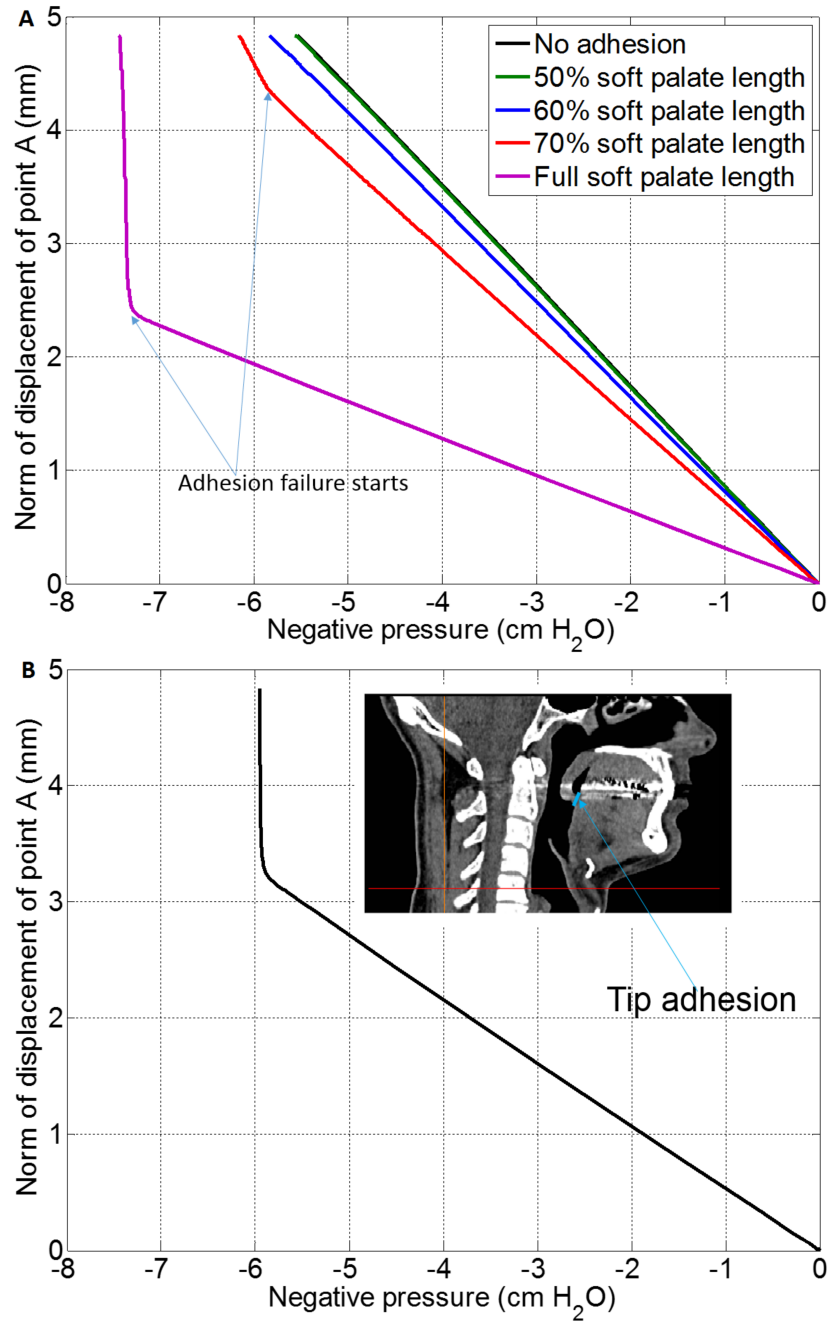




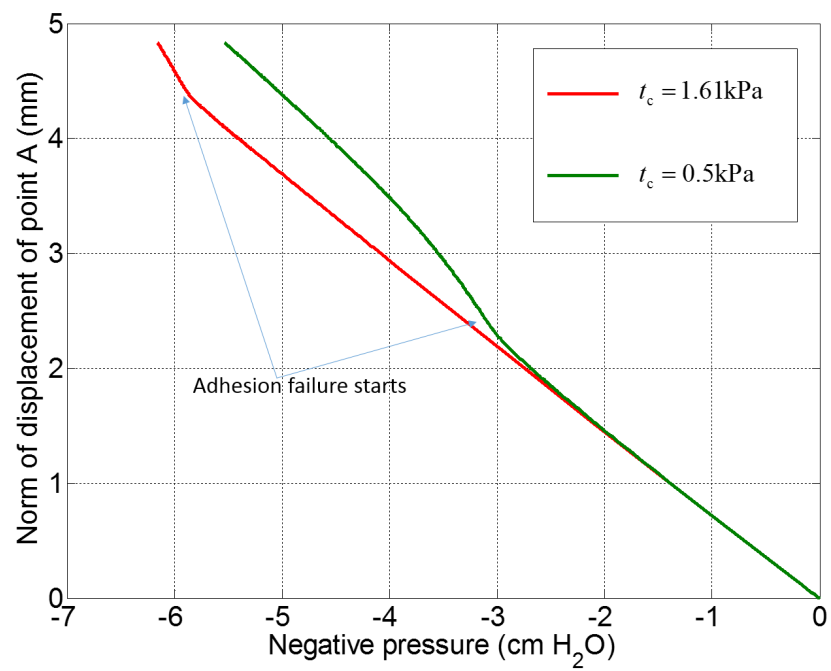
**Fig. 6.** Mesh convergence analysis result. The displacement of the reference point in the midsection of the soft palate's posterior surface was chosen as a critical parameter and compared for different mesh densities.



**Fig. 7.** Collapsed deformation of the 3D patient-specific soft palate model (view from the nasopharynx cavity), the failure of the adhesion for the cohesive layer (70% soft palate length for the cohesive layer) in the sagittal midsection view and the definition of point A and of the norm of its displacement for the 3D patient-specific model in the sagittal midsection plane. **Point A was defined to be the first point of the soft palate posterior surface to be in contact with the pharynx wall in the sagittal midsection plane.** The displacement of point A was used to represent the inclination displacement of soft palate tip's posterior surface.



**Fig. 8.** (A) Negative pressure against the norm of displacement of point A for the 3D patient-specific adhesion model, including different contact length cases except the tip adhesion case. (B) Negative pressure against the norm of displacement of point A for the tip adhesion case. A step-increasing of the displacement of point C was observed, and the closing pressure was confirmed to be  $-5.96$  cm H<sub>2</sub>O. Note that this CT image for the tip adhesion case is from another patient.



**Fig. 9.** Negative pressure against the norm of displacement of point A for the comparison of different critical traction stress cases. The length of the cohesive zone was set to be 70% soft palate length.