

# Emission decay and energy transfer in Yb/Tm Y-codoped fibers based on nano-modified glass

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**Abstract:** We report the results of experimental investigation and theoretical analysis of luminescence decay in Yb/Tm Y-codoped fibers based on nano-modified glass. Based on the experimental results, numerical simulations allowed us to estimate the energy transfer efficiency between Yb<sup>3+</sup> and Tm<sup>3+</sup> ions. It was shown that yttria enhances the Yb/Tm energy transfer making fibers with Y-codoping promising for development of the light sources for laser applications and of up-conversion emitters for visualization applications. These fibers demonstrate energy transfer efficiency of ~50 %, which makes them attractive for diode-pumping of Yb-ions at 975 nm wavelength.

Keywords: Fiber optics, Infrared fiber lasers, Energy transfer

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## 1. Introduction

The interest to Yb/Tm-doped fibers and crystals is caused by the possibility of their direct pumping by conventional low cost telecom diode lasers operating at 975 nm [1-3] - an alternative pumping method of Tm, when simple and cost-effective low- to medium power Tm-fiber sources are sought. In this context the investigation of efficiency of the energy transfer (ET) from Yb<sup>3+</sup> to Tm<sup>3+</sup> ions becomes of great importance [4,5]. This efficiency is known to be dependent on the fiber composition. The spectroscopic investigation of Yb/Tm-doped fibers codoped with yttria was reported recently in [6]. It was assumed that yttria allows enhancing the Yb-Tm ET. The subject of this work is the investigation of the ET processes between Yb<sup>3+</sup> and Tm<sup>3+</sup> ions in relation to the new compositions of core glass of Yb/Tm-doped Y-codoped fibers.

For Yb/Tm system the efficiency of the ET from Yb<sup>3+</sup> to Tm<sup>3+</sup> ions reaches 50% [7]. The typical lifetimes for the <sup>3</sup>F<sub>4</sub> level of Tm<sup>3+</sup> vary from 200 to 650 μs, and for the <sup>2</sup>F<sub>5/2</sub> level of Yb<sup>3+</sup> from 470 to 830 μs depending on the host composition [4, 8-10]. Despite the large importance of the ET processes, there are not so much data in the literature on the efficiency of the ET from Yb<sup>3+</sup> to Tm<sup>3+</sup> ions, which makes our work particularly relevant.

In this paper we report the results of the measurements of Yb<sup>3+</sup> luminescence decay in the presence of Tm<sup>3+</sup> ions in Y-codoped silica glass fibers based on nano-modified glass, the decay time of Tm<sup>3+</sup> ions as well as the calculations of the ET efficiency in the set of Yb/Tm-doped fibers with different doping concentration levels. The technological processes as well as the effect of nano-engineering of the glass host and its influence on the ET efficiency are also being discussed in this paper. The purpose of such nano-engineering is to increase the ET efficiency from Yb<sup>3+</sup> to Tm<sup>3+</sup> ions thus making such fibers promising for development of the light sources for laser applications and of the up-conversion emitters for visualization applications [11].

## 2. Fiber fabrication and material analysis

In our experiment we have studied the set of Yb/Tm-doped silicate fibers, namely: LCT-6, LCT-8 [6], and newly produced LCT-9, LCT-10, and HTY-1 fibers with different concentrations of dopants. The waveguide parameters of the fibers are presented in Table 1, the compositions as well as the calculation and measurement results are given in Table 2. Typical luminescence and optical absorption spectra are presented in Fig.1.

**Table 1. Waveguide parameters of the fibers under investigation**

Fiber Sample	Core diameter, $\mu\text{m}$	Numerical Aperture
LCT-9	18.0	0.29
LCT-6	15.8	0.22
LCT-10	10.0	0.23
LCT-8	14.2	0.26
HTY-1	8.0	0.16

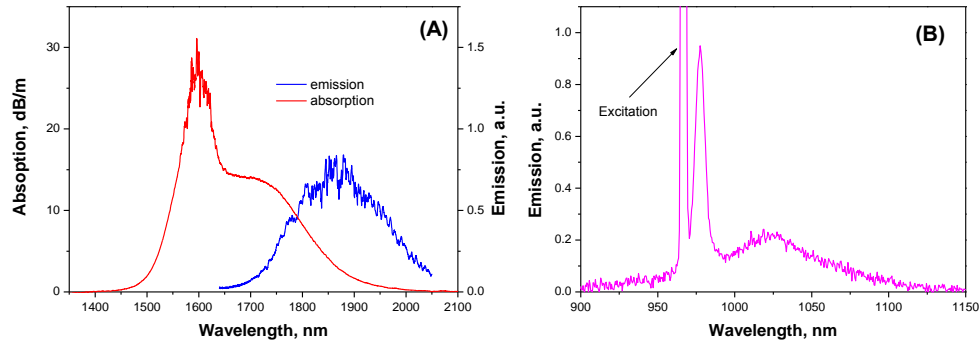


Fig. 1. Typical absorption and emission spectra of  $\text{Tm}^{3+}$  (A), and emission spectrum of  $\text{Yb}^{3+}$  (B).

Yb/Tm co-doped yttrium-alumino-silicate (YAS) fibers were fabricated by modified chemical vapor deposition (MCVD) process in conjunction with solution doping (SD) technique. Since 1987, SD technique has been used to incorporate rare-earths [12] and nanoparticles [10, 13-16] into the silica glass matrix of optical fibers.

We assume that Y-codoping allows modification of the fiber glass by creation of the nanostructures with high active ion concentration. We have studied the nature of the nanostructures in the core glass of our samples by the transmission electron microscopy (TEM). The TEM image of the core region of the LCT-9 fiber is shown in Fig. 2(A). This image suggests that the two separate phases are formed inside the core. The granules of the diameter around 0.5 - 5 nm with dark appearance are formed mainly by the atoms with high atomic mass. The glass outside the black granules is formed mainly by the atoms with lower atomic mass. For the proper analysis of the nature of these phases the Electron Diffraction (ED) pattern (Fig. 2(B)) and EDX analysis were performed. The diffraction pattern showed that the black particle-like granular features have amorphous nature. EDX spectra taken 'on' and 'outside' the granules shown in Fig. 2(C) and Fig. 2(D) showed that the black granules are formed by Y, Al, Tm, Yb and Si oxides. The other part of the matrix outside the dark granules is constructed mostly by the Si-oxide with a trace of Al. It was found from EDX that Tm and Yb are concentrated mainly within the black granules.

From this analysis we can say that the black phases correspond to YAS glass phases and form the hosts for active  $\text{Tm}^{3+}$  and  $\text{Yb}^{3+}$  ions in the core. From the TEM images it is also seen that the YAS glass phases are distributed almost homogeneously in the silicate glass in the form of nano-phases. Thus, these YAS phases help Tm and Yb to be distributed homogeneously in the core. This analysis also cleared that there is a rear-earth-rich (RE-rich) and RE-poor zones within the core. Particularly, the nano YAS phases are the RE-rich zones.

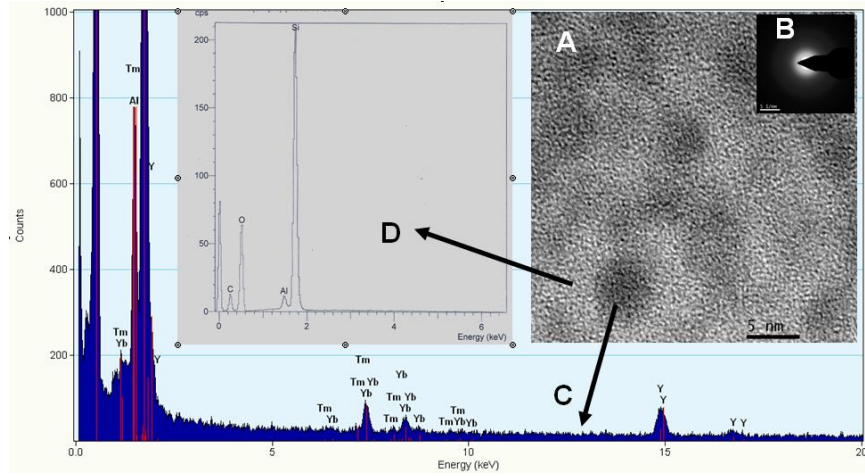


Fig. 2. (A) TEM image and (B) electron diffraction pattern of the fiber (LCT-9), (C) EDX plot taken on and (D) outside of black spots of the fiber sample.

Typical microscopic view of the fiber LCT-9 sample is shown in Fig. 3(A). The refractive index profile of the fiber preform LCT-8 sample measured by the preform analyzer is shown in Fig. 3(B). The dopants' concentrations along the core diameter were measured by the electron probe micro-analysis (EPMA). The dopants' distribution of the fiber LCT-10 sample is shown in Fig. 3(C), and for the other samples it is given in Table 2.

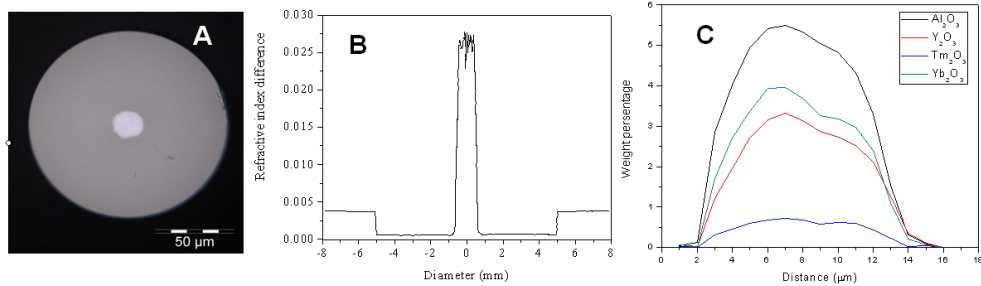


Fig. 3. (A) The microscopic cross-sectional view of the fiber sample (LCT-9), (B) Refractive index profile of fiber preform sample (LCT-8) and (C) EPMA curve for distribution of different dopants along the diameter of fiber sample (LCT-10).

Due to the small size (0.5 – 5.0 nm) of nano YAS-RE phases, fibers exhibit negligible Rayleigh scattering losses [16]. To estimate the possible scattering losses caused by the nano-phases we have performed the loss measurements of the fiber samples. The spectral loss curve of the LCT-8 fiber is shown in Fig. 4. The absorption spectrum of the fiber sample was measured at the wavelength range of 350 – 1700 nm. The absorption peak of  $\text{Yb}^{3+}$  obtained at 975 nm can be associated with a hump at 920 nm due to the transitions from  ${}^2\text{F}_{7/2}$  to  ${}^2\text{F}_{5/2}$  levels. In case of  $\text{Tm}^{3+}$  ions, there are three electronic transitions involved:  ${}^3\text{H}_6 \rightarrow {}^3\text{H}_4$ ,  ${}^3\text{H}_6 \rightarrow {}^3\text{F}_{2,3}$ ,  ${}^3\text{H}_6 \rightarrow {}^1\text{G}_4$ , which correspond to the absorption peaks at 789 nm, 678 nm and 465 nm respectively [17-18]. Another absorption peak of  $\text{Tm}^{3+}$  is also observed at 1205 nm in the attenuation spectra, which represents the electronic transition from  ${}^3\text{H}_6$  to  ${}^3\text{H}_5$  level.

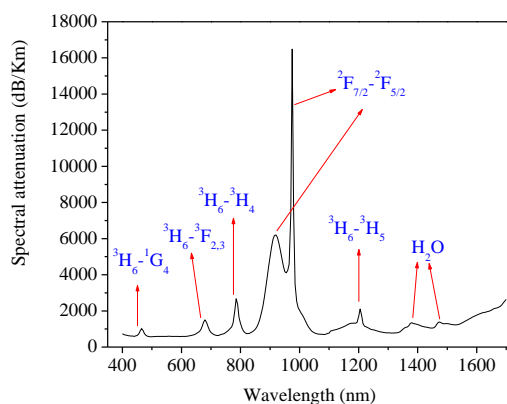


Fig. 4: Spectral attenuation curve of the Tm/Yb codoped fiber (LCT-8).

### 3. Experimental

To find the efficiency of the ET between  $\text{Yb}^{3+}$  and  $\text{Tm}^{3+}$  ions and the dopants' effective lifetimes, we examined luminescence decay kinetics of  $\text{Yb}^{3+} \ ^2\text{F}_{5/2}$  and  $\text{Tm}^{3+} \ ^3\text{F}_4$  levels after pulsed excitation by two diode lasers: one with the emission wavelength of 980 nm for the ET efficiency measurement and the other with the wavelength of 1550 nm for the effective lifetime measurements of  $\text{Tm}^{3+}$  ions.

Three photodetectors were used for the kinetics detection: Ge photodetector operating in 0.8 – 1.8  $\mu\text{m}$  range (ThorLabs PDA50B-EC), Si photodetector (0.4 – 1.1  $\mu\text{m}$ ) (ThorLabs FDS1010) and extended InGaAs photodetector (1.2 – 2.55  $\mu\text{m}$ ) (Hamamatsu G5853-23). Ge and Si photodetectors allowed measuring the kinetics of  $\text{Yb}^{3+}$  ions. Due to fast drop of its sensitivity at the longer wavelength edge the possible impact of Tm emission was negligible. On the other hand, InGaAs photodetector allowed us measuring the kinetics of  $^3\text{F}_4$  level of  $\text{Tm}^{3+}$  ions. That choice of the detectors ensured the spectral separation of our measurements.

Both sources were supplied with electrical pulses of 1 ms duration, at the frequency of 89 Hz, with fall-of time of 1  $\mu\text{s}$ . The optical pulses with average excitation power of around 10 mW were produced. The fall-off times measured with Ge and Si detectors were 15  $\mu\text{s}$  and the fall-off time measured with InGaAs detector was 40  $\mu\text{s}$ , respectively. This elongation is mainly attributed to the detectors' response. The decays were detected from the outer fiber surface and recorded with a digital Instek GDS-3354 oscilloscope.

### 3. Results

The decays of  $^3\text{F}_4$  level of  $\text{Tm}^{3+}$  ions obtained under 1550 nm excitation and collected by InGaAs were exponential, while the decays of  $\text{Yb}^{3+}$  ions obtained under 980 nm excitation were non-exponential. The kinetics for the fibers (LCT-9 and LCT-6) are shown in Fig. 5. and Fig. 6.

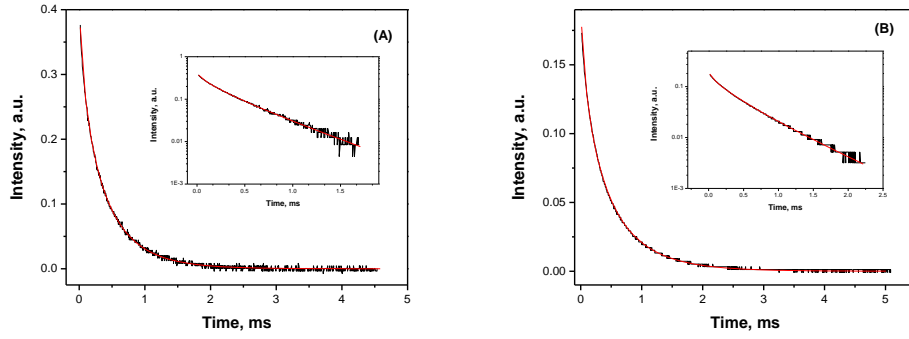


Fig. 5. Typical non-exponential decays with excitation at 980 nm wavelength corresponding to the energy transfer between the ions, measured with Ge photodetector, the black line corresponds to the experimental data and the red one is the approximated using the eq.(2). The inset is in logarithmic scale. (A) LCT-9 and (B) LCT-6 fibers.

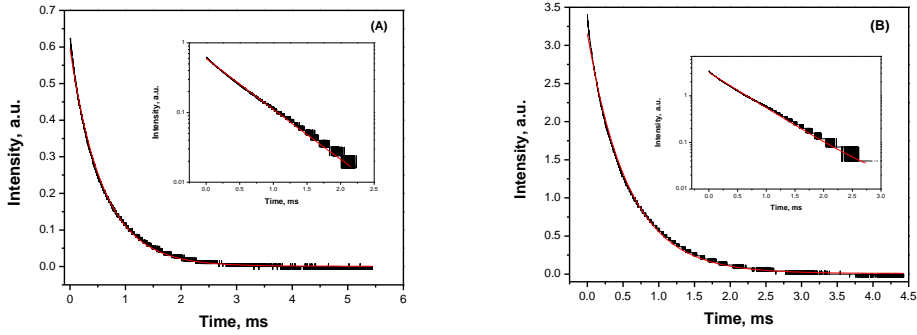


Fig. 6.  $Tm^{3+}$  kinetics with the excitation wavelength of 1550 nm, measured by InGaAs photodetector, the black line corresponds to the experimental data and the red one is the exponential fitting of the curve. The inset is in logarithmic scale. (A) LCT-9 and (B) LCT-6 fibers.

Since the decay kinetics of the  $Yb^{3+}$  ions were not exponential, the effective decay constant  $\tau_{eff}$ , of  $Yb$  kinetics was calculated using eq.(1). These kinetics were also approximated using the eq.(2) (Inokuti-Hirayama formula [19]), which describes the donor fluorescence kinetics under the assumption of dipole-dipole interaction between donor and acceptor for migration-assisted ET [20]. The migration-assisted and direct donor-acceptor energy-transfer rates were also obtained from the fitting of the kinetics using eq.(2). The ET-efficiency  $\eta$  was calculated using the eq.(3) [21] and reached 56% for LCT-9 fiber, which is slightly higher than it was previously reported. For the calculations we used the value of  $Yb^{3+}$  lifetime ( $\tau_{Yb}$ ) equal to 830  $\mu s$  that was measured for the  $Tm^{3+}$ -free fibers [4]. The comparable value of 810  $\mu s$  of  $Yb^{3+}$  lifetime for  $Tm^{3+}$ -free fibers produced by the same technology with similar concentration was obtained in [10].

$$\tau_{eff} = \frac{\int I(t)dt}{\int I(t)dt} \quad (1),$$

$$I(t) = A_0 \exp\left(-\frac{t}{\tau_{Yb}} - \gamma t^{1/2} - Kt\right) \quad (2),$$

$$\eta = 1 - \frac{\int I(t)dt}{A_0\tau_{yb}} \quad (3),$$

where  $I(t)$  is a decay kinetic,  $A_0$  is the amplitude of the decay,  $K$  is the migration-assisted energy-transfer rate, and  $\gamma$  is the direct donor–acceptor ET rate.

The analysis of  $\text{Tm}^{3+}$  kinetics allowed us to determine the lifetime of  ${}^3\text{F}_4$  level for all samples (Table 2). These values are comparable to the previously reported lifetimes (200-650  $\mu\text{s}$  depending on the host [4, 8, 9]). During the fitting procedures the reduced chi-squared distribution of the fitting varied from  $6.5 \times 10^{-8}$  to  $5.5 \times 10^{-7}$  depending on the fiber, and the coefficient of determination ( $R^2$ ) for all the samples was around 0.995. We estimated the standard deviation error in ET-efficiency to be around 2%.

**Table 2. Summary of the parameters obtained in this work.**

Sample	$\tau_{eff}$ , $\mu\text{s}$	$\tau_{\text{Tm}}$ , $\mu\text{s}$	$\text{Yb}_2\text{O}_3$ , wt%	$\text{Y}_2\text{O}_3$ , wt%	$\text{Al}_2\text{O}_3$ , wt%	$\text{Tm}_2\text{O}_3$ , wt%	$\gamma$	$K$ , $\text{s}^{-1}$	$\eta$ , %
LCT-9	450	590	4.5	4.0	6.0	1	39	226	56
LCT-6	520	570	2	1.9	2.0	0.8	33	105	49
LCT-10	540	510	4	3.3	5.5	0.7	22	547	46
LCT-8	570	570	2	3.0	1.0	0.5	23	279	42
HTY-1	720	400	0.9	2.5	4.5	0.75	20	466	21

The luminescence measurements of Tm-free Yb-doped fibers of similar composition produced with the same technology with the concentration of about 1.6 wt.% of Yb [10] showed that  $\text{Yb}^{3+}$  experiences single-exponential decay with a lifetime of about 800  $\mu\text{s}$ . This value significantly exceeded the Yb decay constant  $\tau_{eff}$  of the fibers with a similar concentration, LCT-6, LCT-8, see Table 2. Thus, the non-exponential decay with the lower constant as compared to the Yb lifetime in Tm-free samples is attributed to the ET between  $\text{Yb}^{3+}$  and  $\text{Tm}^{3+}$ , but not to the concentration quenching. We also conclude that the possible concentration quenching is not the critical issue for these fibers.

As it can be seen from Table 2, there is no influence of the Al concentration on the ET efficiency. But at the same time, one can see the direct proportionality between the Y concentration and ET efficiency. The double change in concentration of Y leads to the double change in the efficiency.

#### 4. Discussion

The energy level diagram of the Tm/Yb system (Fig. 7. [4]) shows the processes of ET between the energy levels of  $\text{Tm}^{3+}$  and  $\text{Yb}^{3+}$  ions in the aluminosilicate host.

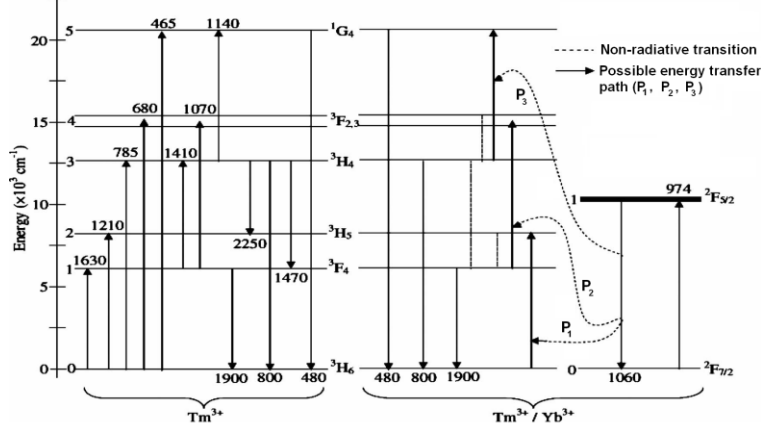
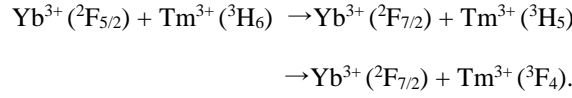
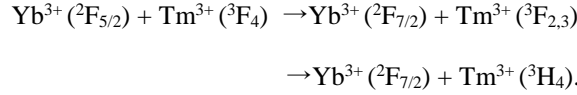


Fig. 7. Energy level diagram and energy-transfer processes in Tm<sup>3+</sup>/Yb<sup>3+</sup> co-doped system [4].

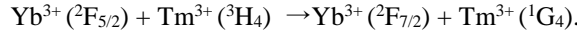
According to the diagram, three transitions are possible in the Tm/Yb system: P<sub>1</sub>, P<sub>2</sub> and P<sub>3</sub>. The first transition results in the 2 μm emission and can be described as



The P<sub>2</sub> transition corresponds to the following relaxation/excitation scheme:



The third one results in population of the <sup>1</sup>G<sub>4</sub> level, which is responsible for the 480 nm wavelength emission:



We applied the commonly used Förster–Dexter [22] approach to the analysis of ET in Yb/Tm system, under the assumption of dipole-dipole interaction between donor and acceptor. As a consequence from this theory [19], the dependency of the direct donor–acceptor ET rate ( $\gamma$ ) on the concentration of the acceptors ( $N_A$ ) can be expressed by:

$$\gamma = \frac{4}{3} \pi^2 N_A \sqrt{C_{DA}} \quad (4)$$

where,  $C_{DA}$  is the donor-acceptor energy transfer microparameter and  $N_A$  is the concentration of the acceptors.

According to (4), there should be a linear dependence between  $\gamma$  and  $N_A$ . **Such a linear behavior, Fig. 8(A).** It means, that energy transfer processes in our Tm/Yb system can be described within the commonly applied Förster–Dexter theory under the assumption of dipole-dipole interaction.

At the same time, we have found that the efficiency of the ET linearly increases with the increase of the Thulium <sup>3</sup>F<sub>4</sub> level lifetime, Fig. 8(B). To explain this phenomenon, we suggest the following mechanisms of the observed correlation. First, this dependence can be explained according to the diagram shown in Fig. 7. The increase of the <sup>3</sup>F<sub>4</sub> level lifetime favors the up-conversion processes. Such processes terminating in <sup>1</sup>G<sub>4</sub> energy level also increase the visible 480 nm wavelength emission intensity which can be interesting for certain applications [11].

Second, there is an obvious dependence of the ET-efficiency on the Y concentration that allows us to conclude that Y favors the energy transfer. Moreover, yttria allows enhancing the lifetime of  $\text{Tm}^{3+}$  in comparison to the Y-free silica hosts. We attribute it to the structural changes of the glass due to the addition of Y to the host. Indeed, the  $\text{Y}_2\text{O}_3$  helps to create nano-structured YAS glass phase in the host and acts as a glass network generator [23-25]. The  $\text{Al}_2\text{O}_3$  increases the refractive index (RI) of silica by  $2.3 \times 10^{-3}$  per mole and is used as intermediate core glass network modifier [26]. It helps to distribute  $\text{Tm}^{3+}$  and  $\text{Yb}^{3+}$  ions homogeneously to decrease phonon energy of silica glass [26]. The modified silica glass host is chosen due to its optical transparency from 200 nm to beyond 3000 nm wavelength region [27]. Furthermore, due to the lower phonon energy of YAS, the magnitude of radiative emission probability of metastable states of  $\text{Tm}^{3+}$  increases and the non-radiative multi-phonon decay decreases. This leads to the enhancement of the ET process from  $\text{Yb}^{3+}$  to  $\text{Tm}^{3+}$  ions. In this system the 975 nm excitation wavelength can be absorbed only by  $\text{Yb}^{3+}$ . Thus, the pump source excites the  $\text{Yb}^{3+}$  first and then the energy is transferred to nearby  $\text{Tm}^{3+}$  ion to excite it in an efficient manner.

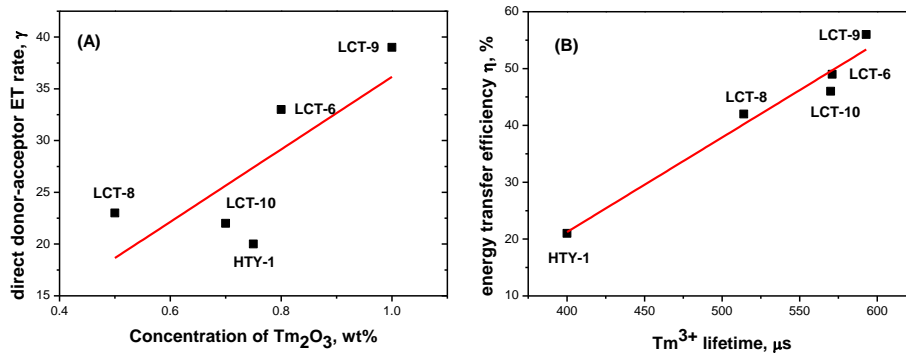


Fig. 8. (A) The dependency of direct donor-acceptor rate on the concentration of  $\text{Tm}_2\text{O}_3$ .  
 (B) Correlation between the transfer efficiency  $\eta$  and  $\text{Tm}^{3+}$  lifetime  
 (linear approximation by the straight line is shown as a guide-to the eye).

## 5. Conclusions

We measured the luminescence decay kinetics of the set of Yb/Tm nano-particles-doped silicate fibers codoped with yttria under excitation with the two pulsed diode lasers: one with the wavelength of 980 nm for the energy transfer analysis and the other with the wavelength of 1550 nm for the calculation of the  $\text{Tm}^{3+} \ ^3\text{F}_4$  level lifetime. We obtained from our experiment the average decay constants and the lifetimes of  $\text{Tm}^{3+} \ ^3\text{F}_4$  level, which varied from 490 to 890  $\mu\text{s}$  and from 180 to 590  $\mu\text{s}$ , respectively, depending on a fiber, as well as the energy transfer efficiency from  $\text{Yb}^{3+}$  to  $\text{Tm}^{3+}$  ions. The latter reached 56%. The correlation between the transfer efficiency and the  $^3\text{F}_4$  level lifetime of  $\text{Tm}^{3+}$  was found and explained. The nano-engineering of the host by the  $\text{Y}_2\text{O}_3$  doping allows to increase the efficiency of ET by creation of the nano-structures with the high doping levels of active  $\text{Tm}^{3+}$  and  $\text{Yb}^{3+}$  ions, which are organized in nano-phases. The  $\text{Y}_2\text{O}_3$  effectively lowers the phonon energy and, as a consequence, increases the ET rates - the origin of the high up-conversion efficiency. The ET efficiency increases with the increase of the  $\text{Y}_2\text{O}_3$  doping levels.

Up-conversion effects make these fibers attractive for development of cost-efficient up-conversion blue light emitters for visualization applications, as well as the effective energy transfer makes these fibers promising for development of Tm-doped fiber lasers pumped with the low-cost diodes operating at 975 nm wavelength.

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