

# The Vanishing water/oil interface in the presence of antagonistic salt

Cite as: J. Chem. Phys. **152**, 124707 (2020); <https://doi.org/10.1063/1.5142811>

Submitted: 18 December 2019 . Accepted: 03 March 2020 . Published Online: 25 March 2020

Gudrun Glende,  Astrid S. de Wijn, and  Faezeh Pousaneh



View Online



Export Citation



CrossMark

## ARTICLES YOU MAY BE INTERESTED IN

[Molecular polarizabilities as fingerprints of perturbations to water by ions and confinement](#)  
The Journal of Chemical Physics **152**, 124501 (2020); <https://doi.org/10.1063/1.5143317>

[TRAVIS—A free analyzer for trajectories from molecular simulation](#)  
The Journal of Chemical Physics **152**, 164105 (2020); <https://doi.org/10.1063/5.0005078>

[Theoretical analysis of electrode-dependent interfacial structures on hydrate-melt electrolytes](#)  
The Journal of Chemical Physics **152**, 124706 (2020); <https://doi.org/10.1063/5.0003196>



**Your Qubits. Measured.**

Meet the next generation of quantum analyzers

- Readout for up to 64 qubits
- Operation at up to 8.5 GHz, mixer-calibration-free
- Signal optimization with minimal latency

[Find out more](#)



# The Vanishing water/oil interface in the presence of antagonistic salt

Cite as: *J. Chem. Phys.* **152**, 124707 (2020); doi: [10.1063/1.5142811](https://doi.org/10.1063/1.5142811)

Submitted: 18 December 2019 • Accepted: 3 March 2020 •

Published Online: 25 March 2020



View Online



Export Citation



CrossMark

Gudrun Glende,<sup>1</sup> Astrid S. de Wijn,<sup>1,2</sup>  and Faezeh Pousaneh<sup>1,a)</sup> 

## AFFILIATIONS

<sup>1</sup>Department of Mechanical and Industrial Engineering, Norwegian University of Science and Technology, 7491 Trondheim, Norway

<sup>2</sup>Department of Physics, Stockholm University, Stockholm, Sweden

<sup>a)</sup>Author to whom correspondence should be addressed: [faezeh.pousaneh@ntnu.no](mailto:faezeh.pousaneh@ntnu.no)

## ABSTRACT

Antagonistic salts are salts that consist of hydrophilic and hydrophobic ions. In a binary mixture of water and an organic solvent, these ions preferentially dissolve into different phases. We investigate the effect of an antagonistic salt, tetraphenylphosphonium chloride  $\text{PPh}_4^+\text{Cl}^-$ , in a mixture of water and 2,6-lutidine by means of Molecular Dynamics (MD) simulations. For increasing concentrations of the salt, the two-phase region is shrunk and the interfacial tension is reduced, in contrast to what happens when a normal salt is added to such a mixture. The MD simulations allow us to investigate in detail the mechanism behind the reduction of the surface tension. We obtain the ion and composition distributions around the interface and determine the hydrogen bonds in the system and conclude that the addition of salt alters the hydrogen bonding.

Published under license by AIP Publishing. <https://doi.org/10.1063/1.5142811>

## I. INTRODUCTION

Mixtures of water and oil are ubiquitous in nature and technology. Biological systems are largely water-based but also contain oily molecules. In maritime and other off-shore applications, lubricants and other oils come into contact with seawater often. What many of these mixtures have in common is the additional presence of salt ions. This addition of a small amount of salt can significantly change the properties of an oil/water mixture.

Most binary mixtures of water and organic molecules are phase-separated at ambient conditions. When a simple inorganic salt is added to the mixture, the two-phase region enlarges.<sup>1</sup> This picture changes for more complex salts with organic groups that are hydrophobic. When an antagonistic (or amphiphilic) salt is added to an oil/water mixture, this can lead to reduction in the interfacial tension between water and oil and, therefore, to a shrinkage of the two-phase region. It has been shown experimentally that the interfacial tension decreases with an increase in the amount of antagonistic salts.<sup>2</sup> If enough salt is added, it can even make the interface disappear altogether and cause the oil and water to mix.<sup>3–7</sup> In addition to changes in the interfacial tension, adding an antagonistic salt can also lead to interesting

structural changes, such as lamellar phases.<sup>3,5,6</sup> These effects have been studied and explained theoretically for a small salt concentration<sup>8</sup> and for higher salt concentrations.<sup>9–11</sup> Analytical calculations of the ion distribution have suggested that the cations and anions of an antagonistic salt tend to adsorb around a water–oil interface.<sup>12</sup>

Here, we study this problem in a different way, using detailed atomistic Molecular Dynamics (MD) simulations. Because such simulations give us the trajectories of all particles in the system, they can be used as numerical experiments. They allow us to observe in much more detail where the ions are and how they are behaving, which enables us to study and verify the mechanisms of the reduction in the interfacial tension. We investigate a mixture of water and 2,6-lutidine (2,6-dimethylpyridin) with various concentrations of the antagonistic salt tetraphenylphosphonium chloride  $\text{PPh}_4^+\text{Cl}^-$ , with hydrophilic anion  $\text{Cl}^-$  and hydrophobic cation  $\text{PPh}_4^+$ .

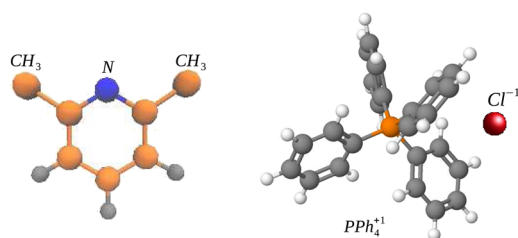
The phase diagram and mixing behavior of neat water/2,6-lutidine mixtures have gained considerable attention in the context of colloidal systems<sup>13–15</sup> and critical Casimir forces.<sup>16,17</sup> Critical Casimir forces rise between two selective (hydrophilic or hydrophobic) colloids in a suspension when the liquid base of the suspension

is near its critical point.<sup>18,19</sup> When colloids are charged, the force is tuned by electrostatic interactions.<sup>16,17,20–25</sup> Moreover, the addition of a salt to the water/2,6-lutidine mixture allows one to tune these interactions and create different structures.<sup>3,5,6,9</sup> Antagonistic salts are of particular interest in this context because the two ions have substantially different interactions with the oil and water.

The neat water/2,6-lutidine binary mixture and its phase diagram have been studied experimentally and analytically<sup>26,27</sup> and recently by molecular dynamics simulations.<sup>28</sup> Without the additional salt, this mixture has a closed-loop phase diagram with a relatively wide temperature miscibility gap, which makes it very suitable for studying (de)mixing. The lower critical point is close to room temperature, around  $T_c \approx 307.1$  K with the lutidine mole fraction  $x_{lut} \approx 6.1\%$ .<sup>26</sup>

An antagonistic salt has one hydrophilic ion and one hydrophobic ion. This makes them potentially very different from more simple hydrophilic salts, where both the cation and anion are hydrophilic. With the addition of such hydrophilic salts, the solubility of water and the organic solvent decreases, and therefore, the two-phase region is enhanced. Rather than both ions preferring to dissolve in water, in antagonistic salts, one ion will prefer to dissolve in the water, while the other prefers to dissolve in the oil. The effect of the antagonistic salt  $\text{Na}^+\text{BPh}_4^-$  on a binary mixture of  $\text{D}_2\text{O}$  and 3-methylpyridine (3MP) was recently studied experimentally by Sadakane *et al.*<sup>3,5</sup> by means of small-angle neutron scattering and optical microscopy. They show that for increasing salt concentration, the two-phase region shrinks and eventually even disappears. A similar effect has been reported with the salt  $\text{PPh}_4^+\text{Cl}^-$  [see Fig. 1 (right)] that we study in this work. The hydrophobic cation has four phenyl rings that interfere with hydrogen bonding and the hydration shell. Thus, the cations preferentially dissolve in an organic solvent (oil), whereas the anions prefer to stay within the water. Consequently, the cations and anions behave antagonistically and may distribute heterogeneously when added to a binary liquid consisting of water and an organic solvent.

In this article, we study the effect of an antagonistic salt, tetraphenylphosphonium chloride  $\text{PPh}_4^+\text{Cl}^-$ , in a mixture of water and 2,6-lutidine by means of atomistic MD simulations. We investigate the effect of the addition of salt on properties of the mixture, such as surface tension and hydrogen bonds.



**FIG. 1.** (Left) Schematic representation of the 2,6-lutidine molecule,  $\text{C}_7\text{H}_9\text{N}$ .  $\text{CH}_3$  groups are shown as single united atoms. For the charge distributions, see Ref. 28. (Right) Schematic representation of salt  $\text{PPh}_4^+\text{Cl}^-$ . The cation is tetraphenylphosphonium  $\text{PPh}_4^+$  with four phenyl rings, where the orange atom in the middle represents the phosphonium. The anion is  $\text{Cl}^-$ .

The organization of this paper is as follows. In Sec. II, we explain the MD simulations setup and models for the water and 2,6-lutidine mixture with the antagonistic salt  $\text{PPh}_4^+\text{Cl}^-$ . Section III gives and discusses the simulation results, which is divided into two subsections, concentration profiles (Sec. III A) and surface tension and hydrogen bonds (Sec. III B). We conclude the paper in Sec. IV.

## II. SIMULATION SETUP

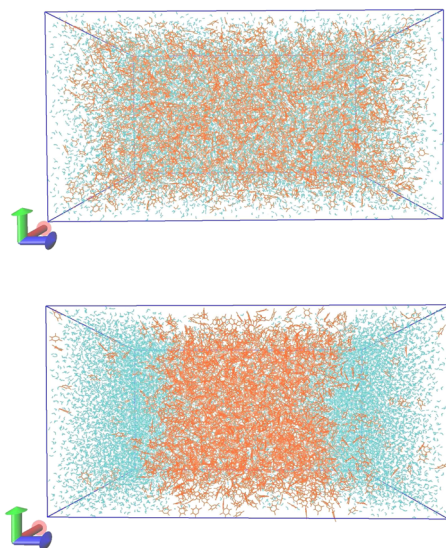
In this work, we consider the same mixture of water/2,6-lutidine provided in Ref. 28. For the organic molecule 2,6-lutidine, we use the recent parameterization<sup>28</sup> developed using the Gromacs package.<sup>29</sup> The model successfully captures the lutidine bulk properties and the lutidine/water binary mixture properties, such as the phase diagram, critical properties, and surface tension. We therefore choose the same topology and charge distributions for the lutidine molecule as in Ref. 28. The lutidine molecule  $\text{C}_7\text{H}_9\text{N}$  is presented by 11 atoms with the  $\text{CH}_3$  groups as united atoms [Fig. 1 (left)]. For the charge distributions on the lutidine, see Ref. 28.

All simulations in this work are performed using the Gromacs/2016 MD simulation package.<sup>29</sup> The Gromos54a7 force-field<sup>30</sup> is used for almost all interactions. Bond lengths are kept fixed at the Gromos54a7 equilibrium length, using the LINCS algorithm.<sup>31</sup> Water is described using the TIP4P/2005 model. For the lutidine molecules, the partial charges and dihedrals are taken from Ref. 28. The Particle Mesh Ewald (PME) approach is applied to the electrostatic interactions beyond a 1.2 (nm) cutoff. A cutoff length of 1.2 (nm) is applied to the van der Waals interactions.

Simulations are performed in the NPT ensemble. The temperature is controlled by using a V-rescale thermostat at  $T = 380$  K, and the pressure is controlled by using a Parrinello–Rahman barostat (isotropic,  $p = 1$  atm). At  $T = 380$  K, the mixture is in the phase-separated region.<sup>28</sup>

For the binary mixture, we simulate a box containing 19 000 water molecules and 3000 lutidine molecules. This corresponds to 13% mole fraction of lutidine. The initial box size is (7, 7, 19) nm [see Fig. 2 (top)]. The simulation time step is 2 fs. We run the simulation for 500 ns to equilibrate the system and allow the water and lutidine to phase separate. The equilibrated configuration is shown in Fig. 2 (bottom).

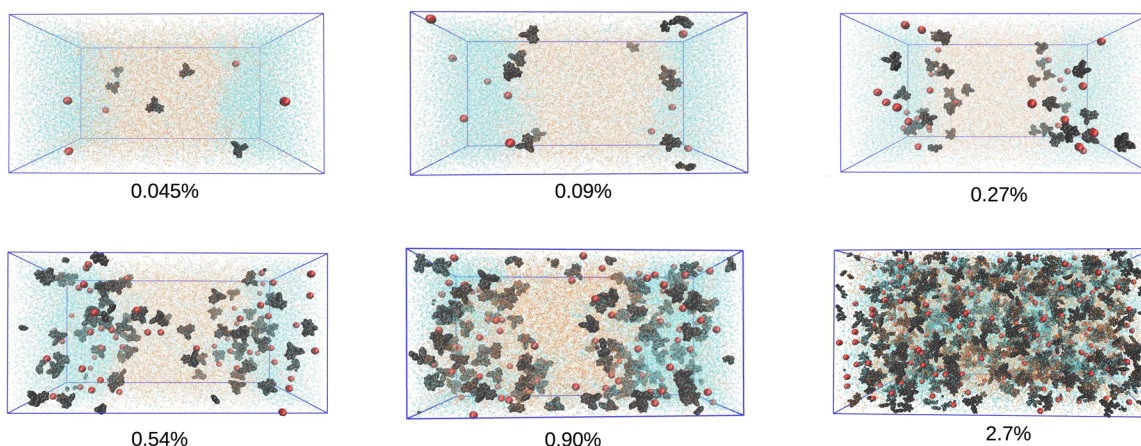
Next, we add the antagonistic salt to the equilibrated mixture. We obtain the topology for the antagonistic salt  $\text{PPh}_4^+\text{Cl}^-$  using the automatic topology builder<sup>32</sup> for the Gromos54a7 force field [see Fig. 1 (right)]. In order to investigate the effect of the salt on the interfacial tension of the water/lutidine mixture, we simulate six different systems with different salt concentrations from 0.045% to 2.70% mole fractions of salt. For every ion we add to the box, we remove one water molecule. We then run the simulations for up to 600 ns depending on the system. Equilibrium typically is reached around 200–300 ns. We verify that the system has reached equilibrium not only by checking density, pressure, and temperature but also the structural properties such as partial densities, radial distribution functions, and hydrogen bonds.



**FIG. 2.** (Top) The initial configuration of water/2,6-lutidine box and (bottom) simulation result after 500 ns at  $T = 380$  K. Blue molecules indicate water molecules, and the orange-colored ring molecules indicate lutidine molecules.

### III. RESULTS

Examples of equilibrated water/2,6-lutidine/salt configurations are shown in Fig. 3. From the snapshots, it can be seen that the salt tends to locate around the interface. As the concentration increases, the interface becomes saturated and the salt ions are present deeper into the bulk of the water and lutidine. Finally, at the extremely high concentration of 2.70% mole fraction, the interface disappears completely and the water and lutidine are mixed.



**FIG. 3.** Equilibrated configurations of the water/2,6-lutidine systems with different salt concentrations. Red spherical particles indicate the  $\text{Cl}^-$  anion, and black bigger particles indicate  $\text{PPh}^+$  cations. The salt concentrations are indicated below each figure. At lower concentrations, the ions predominantly sit around the interface, but at the highest concentration, the interface disappears and the system is mixed.

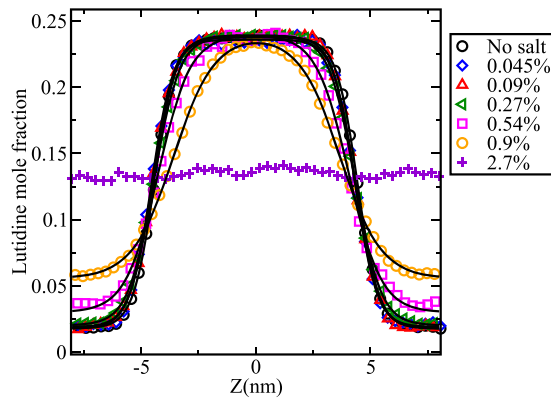
### A. Concentration profiles

We have investigated the concentration profiles in more detail. We first focus on lutidine mole fractions in both phases, lutidine-rich and lutidine-poor. Figure 4 shows the lutidine mole fractions as a function of the position in the box, obtained from simulations. One can see that for higher salt concentrations, the interface between two phases becomes softer. At the highest concentration, the system becomes homogeneous with water and lutidine mixed. The concentration values in the two phases, water-rich and lutidine-rich, can be obtained from the classical theories for interfaces between coexisting phases, such as Cahn and Hilliard or Landau–Ginzburg theory, which both predict a hyperbolic-tangent shaped interfacial density profile,<sup>28,33–35</sup>

$$w_{\text{lut}}(z) = w_{\text{lut}}^{\text{p}} + \frac{w_{\text{lut}}^{\text{r}} - w_{\text{lut}}^{\text{p}}}{2} \left[ \tanh\left(\frac{z - z_0 + c}{\lambda}\right) - \tanh\left(\frac{z - z_0 - c}{\lambda}\right) \right], \quad (1)$$

with  $w_{\text{lut}}^{\text{r}}$  and  $w_{\text{lut}}^{\text{p}}$  as the fractions of 2,6-lutidine in the lutidine-rich and lutidine-poor phases, respectively. The width of the interface is given by  $\lambda$ .  $c$  is the half-width of the lutidine-rich region, and  $z_0$  is the center of the lutidine-rich phase. We fit the concentration profiles from the simulations to this function. These fits are indicated by solid lines in Fig. 4. The fit parameters obtained for the lutidine mole fractions in the rich and poor phases are shown in Fig. 5 and given in the Appendix. The values indicate that with an increase in the salt concentration, the lutidine mole fraction in the lutidine-rich phase decreases, while it increases in the lutidine-poor phase.

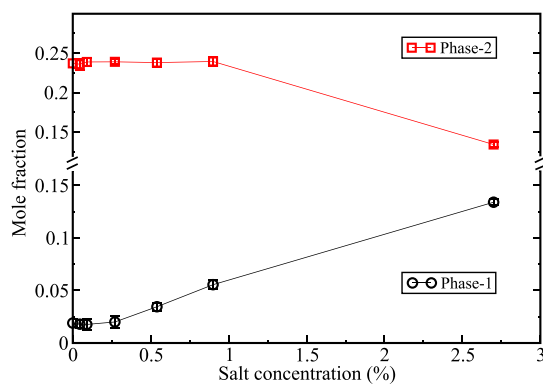
Next, we study the ion concentrations in the two phases and near the interfaces. Figure 6 shows ion concentrations obtained from simulation: the top figure shows the  $\text{PPh}_4^+$  mole fractions and the bottom one shows the  $\text{Cl}^-$  mole fractions. There are clear concentration peaks around the interface between water and lutidine, rather



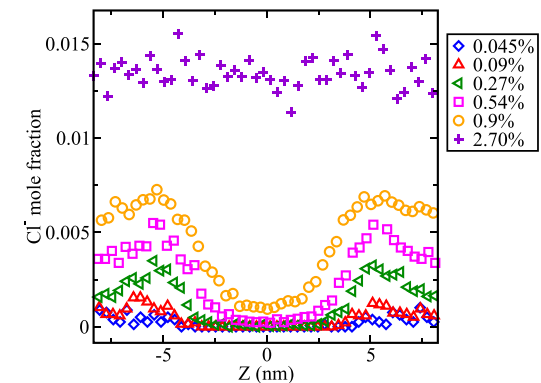
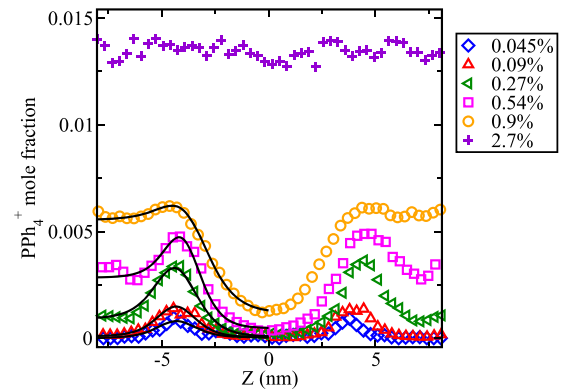
**FIG. 4.** Mole fraction of lutidine at different salt concentrations as a function of position in the simulation box along the axis perpendicular to the interfaces.

than each ion staying at its preferred phase (hydrophilic anion into water and hydrophobic cations into lutidine). At extremely high salt concentration, when the water and lutidine are mixed, the salt concentration becomes homogeneous. In order to obtain the concentration values of  $\text{PPh}_4^+$  in each phase, we fit the concentration profiles to the modified hyperbolic-tangent function from Cahn and Hilliard theory. The modified function corresponds to an asymmetric interface where the three concentrations, at the interface,  $w^m$ , and in the two phases,  $w^r$ ,  $w^l$ , next to the interface, are different. The function is

$$w(z) = \frac{(w^l + w^r)}{2} - \left[ \frac{\tanh(c/\lambda)[w^l - w^r] + [w^l + w^r - 2w^m]}{4 \tanh(c/\lambda)} \right] \times \tanh\left(\frac{z - z_0 + c}{\lambda}\right) - \left[ \frac{\tanh(c/\lambda)[w^l - w^r] - [w^l + w^r - 2w^m]}{4 \tanh(c/\lambda)} \right] \times \tanh\left(\frac{z - z_0 - c}{\lambda}\right), \quad (2)$$

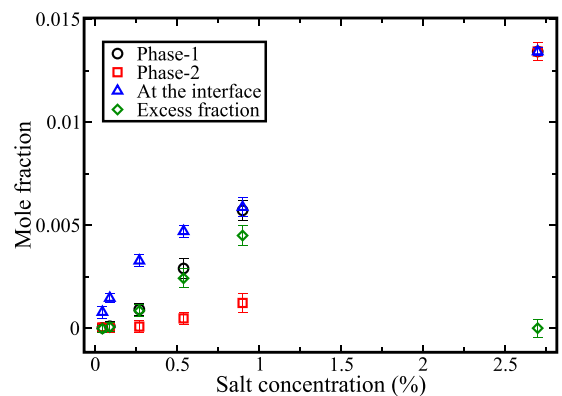


**FIG. 5.** Lutidine mole fraction in the two phases, phase-1 (lutidine-poor) and phase-2 (lutidine-rich) obtained from fitting to Eq. (2).

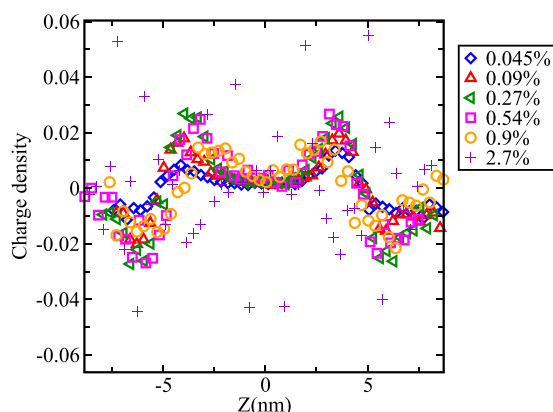


**FIG. 6.** The mole fractions of  $\text{PPh}_4^+$  (top) and  $\text{Cl}^-$  (bottom) for different salt concentrations as a function of position in the simulation box along the axis perpendicular to the interfaces.

where  $w^r$ ,  $w^l$ , and  $w^m$  are the three concentrations. The solid lines in Fig. 6 (top) show the fittings to the last equation. They indicate that with an increase in the salt concentration, the amount of ions around the interfaces increases. This can be better seen in Fig. 7. The fit parameters are also given in the Appendix.



**FIG. 7.**  $\text{PPh}_4^+$  mole fraction in different phases, phase-1 (lutidine-poor) and phase-2 (lutidine-rich), and at the interface obtained from fitting to Eq. (2). The excess mole fraction of the cation between two phases is also shown as green diamonds.



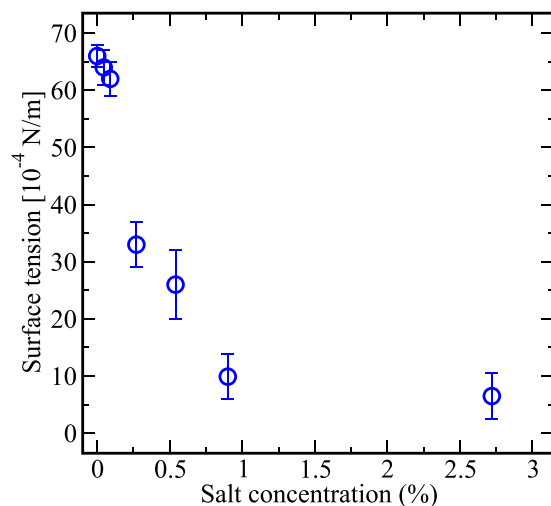
**FIG. 8.** Charge number density for different salt concentrations as a function of position in the simulation box along the axis perpendicular to the interfaces.

We furthermore show the charge density profile in Fig. 8. The figure illustrates that the positive and negative charges pile up when the salt concentration is increased up to a point. However, at large salt concentrations, they become uniformly distributed due to mixing in the system.

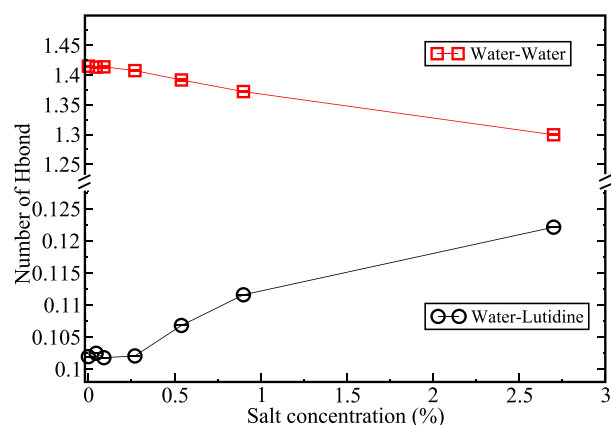
## B. Interfacial tension

In order to obtain the interfacial tension, we continue the simulations in the NPT ensemble using a Berendsen barostat for 200 ns. In Gromacs, the average interfacial tension  $\gamma$  can be calculated directly from the difference between the normal and the lateral pressure,

$$\gamma(t) = \frac{1}{n} \int_0^{L_z} \left[ P_{zz}(z, t) - \frac{P_{xx}(z, t) + P_{yy}(z, t)}{2} \right] dz, \quad (3)$$



**FIG. 9.** Interfacial tension in systems with different salt concentrations. Interfacial tension decreases by increasing the salt concentration in the system.



**FIG. 10.** Number of hydrogen bonds between water-water and water-lutidine molecules (per water molecule) in systems with different salt concentrations.

where  $L_z$  is the height of the box and  $n$  is the number of surfaces. The results are plotted in Fig. 9. The error estimates for the averages are obtained based on block averages. We can see a clear reduction of interfacial tension with an increase in the salt concentration.

To further investigate the effect of antagonistic salt in the mixture, we now focus on the hydrogen bonds between lutidine and water. We obtain the number of hydrogen bonds between water-lutidine molecules and between water-water molecules using Gromacs. Hydrogen bonds are determined based on cutoffs for the angle hydrogen-donor-acceptor and the distance donor-acceptor. O and N are acceptors here.<sup>29</sup> The results are plotted in Fig. 10 for different salt concentrations. The number of hydrogen bonds between water and water decreases with an increase in the salt concentration, while the number of hydrogen bonds between water and lutidine increases. This is a clear indication of the mixing of the water and lutidine. The error estimate of the averages is obtained based on block averages.

## IV. CONCLUSION

The addition of antagonistic salt tetraphenylphosphonium chloride  $\text{PPh}_4^+\text{Cl}^-$  in water/2,6-lutidine has been studied by using Molecular Dynamics simulations. The salt has a hydrophilic anion and hydrophobic cation, therefore different affinity with the water and lutidine phases. We have examined how the salt affects the mixture properties. We study the concentration of all four components, and we determine that the ions tend to stay around the interface of the two phases. However, the ion concentration profiles in both phases become the same at a high value of salt concentration. Thus, our simulations reveal that with increasing salt concentration, the two-phase region starts to shrink and finally at a high salt concentration, water and lutidine mix.

We obtain the surface tension between water and lutidine phases and show that it decreases with increasing the salt concentration. We further investigate the hydrogen bonds between water and lutidine molecules. We report that with an increase in the salt concentration, the number of hydrogen bonds between water-water molecules decreases, while the number of hydrogen bonds between

water–lutidine increases. This is one of the indications for the water and lutidine mixing.

It would be interesting to investigate the possibility of mesophases before and after mixing as well as the structure of the mixed phase.<sup>36,37</sup> However, such a study would necessitate a much bigger system than the one we have simulated here. We are limited in the system size by the computational power required for the detailed atomistic description of the molecules.

## ACKNOWLEDGMENTS

This work was supported by the National Infrastructure for Computational Science in Norway (UNINETT Sigma2) with computer time for the Center for High Performance Computing (Grant Nos. NN9573K and NN9572K). The authors acknowledge the Research Council of Norway for NFR Project No. 275507 for financial support.

## APPENDIX: TABLES OF FITTING PARAMETERS

The fit parameters obtained for the lutidine and PPh<sub>4</sub><sup>+</sup> mole fractions in the rich and poor phases are given in Tables I and II, respectively.

**TABLE I.** Lutidine mole fractions at two phases, lutidine-poor (phase-1) and lutidine-rich (phase-2), for systems with different salt concentrations.

Salt (%)	Phase-1	Phase-2
No salt	0.019 ± 0.004	0.237 ± 0.008
0.045	0.0179 ± 0.004	0.236 ± 0.007
0.09	0.0177 ± 0.004	0.238 ± 0.005
0.27	0.020 ± 0.006	0.239 ± 0.003
0.54	0.0345 ± 0.004	0.238 ± 0.005
0.90	0.055 ± 0.004	0.239 ± 0.006
2.70	0.134 ± 0.003	0.134 ± 0.003

**TABLE II.** PPh<sub>4</sub><sup>+</sup> mole fractions at two phases, lutidine-poor (phase-1) and lutidine-rich (phase-2), and at the interface for systems with different salt concentrations. The values are given in unit of 10<sup>-3</sup>.

Salt (%)	Phase-1	Phase-2	At the interface
0.045	0.0014 ± 0.0004	0.0037 ± 0.0002	0.077 ± 0.003
0.09	0.009 ± 0.0002	0.004 ± 0.0002	0.14 ± 0.01
0.27	0.091 ± 0.003	0.0076 ± 0.0003	0.32 ± 0.03
0.54	0.28 ± 0.06	0.047 ± 0.003	0.46 ± 0.06
0.90	0.57 ± 0.08	0.122 ± 0.04	0.58 ± 0.04
2.70	1.34 ± 0.2	1.34 ± 0.2	1.34 ± 0.2

## REFERENCES

- V. Balevicius and H. Fuess, *Phys. Chem. Chem. Phys.* **1**, 1507–1510 (1999).
- G. Luo, S. Malkova, J. Yoon, D. G. Schultz, B. Lin, M. Meron, I. Benjamin, P. Vanýsek, and M. L. Schlossman, *Science* **311**, 216–218 (2006).
- K. Sadakane, M. Nagao, H. Endo, and H. Seto, *Chem. Lett.* **139**, 234905 (2013).
- K. Sadakane, A. Onuki, K. Nishida, S. Koizumi, and H. Seto, *Soft Matter* **7**, 1334 (2011).
- K. Sadakane, A. Onuki, K. Nishida, S. Koizumi, and H. Seto, *Phys. Rev. Lett.* **103**, 167803 (2009).
- J. Leys, D. Subramanian, E. Rodezno, B. Hammouda, and M. A. Anisimov, *Soft Matter* **9**, 9326 (2013).
- K. Sadakane, Y. Horikawa, M. Nagao, and H. Seto, *Chem. Lett.* **41**, 1075–1077 (2012).
- A. Onuki and H. Kitamura, *J. Chem. Phys.* **121**, 3143 (2004).
- F. Pousaneh and A. Ciach, *Soft Matter* **10**, 8188–8201 (2014).
- A. Onuki, T. Araki, and R. Okamoto, *Phys. Rev. Lett.* **23**, 284113 (2011).
- A. Onuki and R. Okamoto, *Curr. Opin. Colloid Interface Sci.* **16**, 525 (2011).
- A. Onuki, *Phys. Rev. E* **73**, 021506 (2006).
- M. L. Broide, Y. Garrabos, and D. Beysens, *J. Chem. Phys.* **47**, 3768 (1993).
- B. J. Frisken, F. Ferri, and D. S. Cannel, *Phys. Rev. Lett.* **66**, 2754 (1991).
- P. D. Gallagher and J. M. Maher, *Phys. Rev. A* **46**, 2012 (1992).
- A. Gambassi, A. Maciolek, C. Hertlein, U. Nellen, L. Helden, C. Bechinger, and S. Dietrich, *Phys. Rev. E* **80**, 061143 (2009).
- C. Hertlein, L. Helden, A. Gambassi, S. Dietrich, and C. Bechinger, *Nature* **451**, 172 (2008).
- M. E. Fisher and P. G. de Gennes, *C. R. Acad. Sci. Paris B* **287**, 207 (1978).
- M. Krech, *J. Phys.: Condens. Matter* **11**, R391 (1999).
- U. Nellen, L. Helden, and C. Bechinger, *Europhys. Lett.* **88**, 26001 (2009).
- U. Nellen, J. Dietrich, L. Helden, S. Chodankar, K. Nygard, J. F. van der Veen, and C. Bechinger, *Soft Matter* **7**, 5360 (2011).
- M. Bier, A. Gambassi, M. Oettel, and S. Dietrich, *Europhys. Lett.* **95**, 60001 (2011).
- F. Pousaneh and A. Ciach, *J. Phys.: Condens. Matter* **23**, 412101 (2011).
- F. Pousaneh, A. Ciach, and A. Maciolek, *Soft Matter* **8**, 7567–7581 (2012).
- F. Pousaneh, A. Ciach, and A. Maciolek, *Soft Matter* **10**, 470–483 (2014).
- R. J. L. Andon and J. D. Cox, *J. Chem. Soc.* **1**, 4601–4606 (1952).
- Y. Jayalaxshmi, J. S. V. Duijneveldt, and D. Beysens, *J. Chem. Phys.* **100**, 604 (1993).
- F. Pousaneh, O. Edholm, and A. Maciolek, *J. Chem. Phys.* **145**, 014501 (2016).
- D. van der Spoel, E. Lindahl, B. Hess, G. Groenhof, A. E. Mark, and H. J. Berendsen, *J. Comput. Chem.* **26**, 1701–1718 (2005).
- N. Schmid, A. P. Eichenberger, A. Choutko, S. Riniker, M. Winger, A. E. Mark, and W. F. van Gunsteren, *Eur. Biophys. J.* **40**, 843–856 (2011).
- B. Hess, H. Bekker, H. J. C. Berendsen, and J. G. E. M. Fraaije, *J. Comput. Chem.* **18**, 1463–1472 (1997).
- A. K. Malde, L. Zuo, M. Breeze, M. Stroet, D. Poger, P. C. Nair, C. Oostenbrink, and A. E. Mark, *J. Chem. Theory Comput.* **7**, 4026–4037 (2011).
- J. W. Cahn and J. E. Hilliard, *J. Chem. Phys.* **28**, 258 (1958).
- G. A. Chapela, G. Saville, S. M. Thompson, and J. S. Rowlinson, *J. Chem. Soc., Faraday Trans. 2* **73**, 1133 (1977).
- K. Binder, M. Bowker, J. E. Inglesfield, and P. J. Rous, *Cohesion and Structure of Surfaces* (Elsevier Science, The Netherlands, 1995).
- N. Tasios, S. Samin, R. van Roij, and M. Dijkstra, *Phys. Rev. Lett.* **119**, 218001 (2017).
- T. Araki and A. Onuki, *J. Phys.: Condens. Matter* **21**, 424116 (2009).